

Standard Model: Flavor mixing and CP violation

Flavor Mixing and CP Violation

1. CKM Matrix
2. Mixing of neutral mesons
3. CP violation in the B^0 system

1. CKM Matrix

$$\begin{pmatrix} d' \\ s' \\ b' \end{pmatrix} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \cdot \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$

weak eigenstates

CKM matrix

mass eigenstates

Unitarity

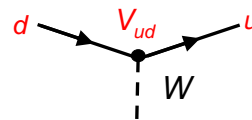
$$V_{CKM} V_{CKM}^+ = 1$$

Charged currents:

$$J_\mu^+ \propto (\bar{u}, \bar{c}, \bar{t}) \gamma_\mu (1 - \gamma_5) \begin{pmatrix} d' \\ s' \\ b' \end{pmatrix} = (\bar{u}, \bar{c}, \bar{t}) \gamma_\mu (1 - \gamma_5) V_{CKM} \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$

weak

mass/
flavor



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1.1 Parameters of CKM matrix

Number of independent parameters:

18 parameter (9 complex elements)
 -5 relative quark phases (unobservable)
 -9 unitarity conditions
 =4 independent parameters: **3 angles + 1 phase**

PDG parametrization

3 Euler angles

$$\theta_{23}, \theta_{13}, \theta_{12}$$

1 Phase

$$\delta$$

$$\begin{pmatrix} d' \\ s' \\ b' \end{pmatrix} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \begin{pmatrix} c_{13} & 0 & s_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta} & 0 & c_{13} \end{pmatrix} \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$

$$\begin{pmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta} \\ -s_{12}c_{23} - c_{12}s_{23}s_{13}e^{i\delta} & c_{12}c_{23} - s_{12}s_{23}s_{13}e^{i\delta} & s_{23}c_{13} \\ s_{12}s_{23} - c_{12}c_{23}s_{13}e^{i\delta} & -c_{12}s_{23} - s_{12}c_{23}s_{13}e^{i\delta} & c_{23}c_{13} \end{pmatrix}$$

where $c_{ij} = \cos\theta_{ij}$, $s_{ij} = \sin\theta_{ij}$

Remark: Unobservable relative quark phases

Phases of left-handed fields in J^{cc} are unobservable: possible redefinition

$$u_L \rightarrow e^{i\phi(u)}u_L \quad c_L \rightarrow e^{i\phi(c)}c_L \quad t_L \rightarrow e^{i\phi(t)}t_L$$

$$d_L \rightarrow e^{i\phi(d)}d_L \quad s_L \rightarrow e^{i\phi(s)}s_L \quad b_L \rightarrow e^{i\phi(b)}b_L$$

Real numbers

Under phase transformation:

$$V \rightarrow \begin{pmatrix} e^{-i\phi(u)} & 0 & 0 \\ 0 & e^{-i\phi(c)} & 0 \\ 0 & 0 & e^{-i\phi(t)} \end{pmatrix} \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \begin{pmatrix} e^{i\phi(d)} & 0 & 0 \\ 0 & e^{i\phi(s)} & 0 \\ 0 & 0 & e^{i\phi(b)} \end{pmatrix}$$

$$V_{\alpha j} \rightarrow \exp[i(\phi(j) - \phi(\alpha))]V_{\alpha j}$$

5 unobservable phase differences !

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Magnitude of elements

$$\begin{pmatrix} d' \\ s' \\ b' \end{pmatrix} = \begin{pmatrix} u & d & s & b \\ c & & & \\ t & & & \end{pmatrix} \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$

komplex in $O(\lambda^3)$

Wolfenstein Parametrization $\lambda, A, \rho, \eta, \lambda = 0.22$

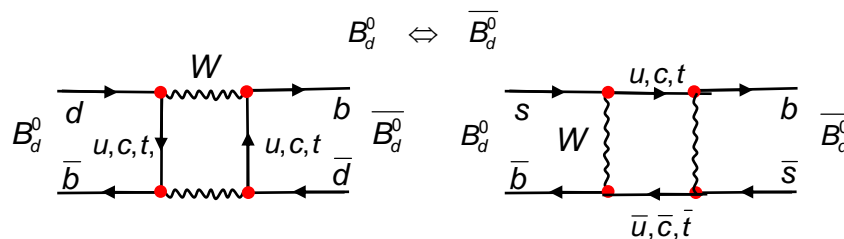
$$V_{CKM} = \begin{pmatrix} 1 - \frac{\lambda^2}{2} & \lambda & A\lambda^3(\rho - i\eta) \\ -\lambda & 1 - \frac{\lambda^2}{2} & A\lambda^2 \\ A\lambda^3(1 - \rho - i\eta) & -A\lambda^2 & 1 \end{pmatrix} + O(\lambda^4)$$

$V_{ub} = |V_{ub}| e^{-i\gamma}$
 $V_{td} = |V_{td}| e^{-i\beta}$

Reflects hierarchy of elements in $O(\lambda)$

2. Mixing of neutral mesons

As result of the quark mixing the Standard Model predicts oscillations of neutral mesons:



Neutral mesons:

	$ P^0\rangle: K^0 = d\bar{s}\rangle$	$D^0 = \bar{u}c\rangle$	$B_d^0 = d\bar{b}\rangle$	$B_s^0 = s\bar{b}\rangle$
	$ \bar{P}^0\rangle: \bar{K}^0 = \bar{d}s\rangle$	$\bar{D}^0 = \bar{u}c\rangle$	$\bar{B}_d^0 = \bar{d}b\rangle$	$\bar{B}_s^0 = \bar{s}b\rangle$
discovery of mixing	1960	2007	1987	2006

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2.1 Mixing phenomenology

Consider time dependent Schrödinger eq. for 2 component wave function $\begin{pmatrix} B^0 \\ \bar{B}^0 \end{pmatrix}$:

$$i \frac{d}{dt} \begin{pmatrix} B^0 \\ \bar{B}^0 \end{pmatrix} = \mathbf{H} \begin{pmatrix} B^0 \\ \bar{B}^0 \end{pmatrix} = \left(\mathbf{M} - \frac{i}{2} \mathbf{\Gamma} \right) \begin{pmatrix} B^0 \\ \bar{B}^0 \end{pmatrix} = \begin{pmatrix} m_{11} - \frac{i}{2} \Gamma_{11} & m_{12}^* - \frac{i}{2} \Gamma_{12} \\ m_{12} - \frac{i}{2} \Gamma_{12} & m_{22} - \frac{i}{2} \Gamma_{22} \end{pmatrix} \begin{pmatrix} B^0 \\ \bar{B}^0 \end{pmatrix}$$

Dispersive & absorptive

As the matrix H is not diagonal B^0 and \bar{B}^0 are not mass eigenstates.

Diagonalizing H of finds the mass eigenstates:

$$\begin{aligned} |B_L\rangle &= p |B^0\rangle + q |\bar{B}^0\rangle & \text{with } m_L, \Gamma_L & & |p|^2 + |q|^2 = 1 \\ |B_H\rangle &= p |B^0\rangle - q |\bar{B}^0\rangle & \text{with } m_H, \Gamma_H & & \text{complex coefficients} \end{aligned}$$

Free particle wave funct. $|B_{H,L}(t)\rangle = |B_{H,L}(0)\rangle \cdot \underbrace{e^{-im_{H,L}t} \cdot e^{-\frac{1}{2}\Gamma_{H,L}t}}_{b_{H,L}(t)}$ $\Delta m = m_H - m_L$
 $\Delta \Gamma = \Gamma_H - \Gamma_L$

Time development of B^0 and \bar{B}^0

$$|B^0\rangle = \frac{1}{2p} (|B_L\rangle + |B_H\rangle) \quad |\bar{B}^0\rangle = \frac{1}{2q} (|B_L\rangle - |B_H\rangle)$$

$$\begin{aligned} |\psi_{B^0}(t)\rangle &= \frac{|B_L(t)\rangle + |B_H(t)\rangle}{2p} = \frac{1}{2p} \left(b_L(t) \cdot (p |B^0\rangle + q |\bar{B}^0\rangle) + b_H(t) \cdot (p |B^0\rangle - q |\bar{B}^0\rangle) \right) \\ &= f_+(t) \cdot |B^0\rangle + \frac{q}{p} f_-(t) \cdot |\bar{B}^0\rangle \quad f_{\pm}(t) = \frac{1}{2} \cdot \left[e^{-im_+t} e^{-\Gamma_+t/2} \pm e^{-im_-t} e^{-\Gamma_-t/2} \right] \end{aligned}$$

$$|\psi_{\bar{B}^0}(t)\rangle = f_+(t) \cdot |\bar{B}^0\rangle + \frac{p}{q} f_-(t) \cdot |B^0\rangle$$

$$\begin{aligned} P(B^0(0) \rightarrow B^0(t)) &= |f_+(t)|^2 \\ P(B^0(0) \rightarrow \bar{B}^0(t)) &= \left| \frac{q}{p} \right|^2 |f_-(t)|^2 \end{aligned}$$

$$\begin{aligned} P(\bar{B}^0(0) \rightarrow \bar{B}^0(t)) &= |f_+(t)|^2 \\ P(\bar{B}^0(0) \rightarrow B^0(t)) &= \left| \frac{p}{q} \right|^2 |f_-(t)|^2 \end{aligned}$$

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Oscillation frequency

$$P(B^0 \rightarrow B^0) = P(\bar{B}^0 \rightarrow \bar{B}^0) = \frac{1}{4} \left[e^{-\Gamma_L t} + e^{-\Gamma_H t} + 2e^{-(\Gamma_L + \Gamma_H)t/2} \cos \Delta m t \right]$$

CPT

$$P(B^0 \rightarrow \bar{B}^0) = \frac{1}{4} \left| \frac{q}{p} \right|^2 \left[e^{-\Gamma_L t} + e^{-\Gamma_H t} - 2e^{-(\Gamma_L + \Gamma_H)t/2} \cos \Delta m t \right]$$

$$P(\bar{B}^0 \rightarrow B^0) = \frac{1}{4} \left| \frac{p}{q} \right|^2 \left[e^{-\Gamma_L t} + e^{-\Gamma_H t} - 2e^{-(\Gamma_L + \Gamma_H)t/2} \cos \Delta m t \right]$$

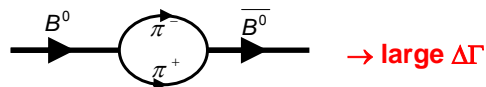
CP- violation in mixing: $P(B^0 \rightarrow \bar{B}^0) \neq P(\bar{B}^0 \rightarrow B^0) \Rightarrow \left| \frac{q}{p} \right| \neq 1$

2.2 Standard Model prediction for B^0 mixing

Mixing mechanisms:

- Mixing through decay:

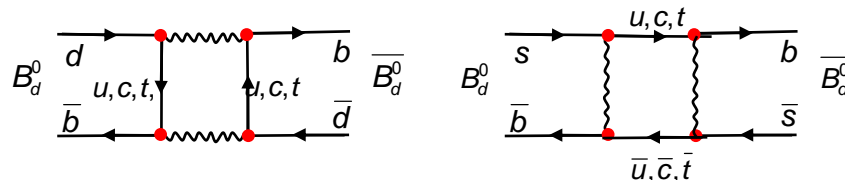
„long distant, on-shell states“



For B mesons there are many possible hadronic decays $\rightarrow \Gamma$ is large in addition decays like $B \rightarrow \pi\pi$ are suppressed

$$\Rightarrow y = \frac{\Delta\Gamma}{2\Gamma} \text{ is small} = \begin{cases} \approx 0 \text{ for } B_d^0 \\ \approx O(0.1) \text{ for } B_s^0 \end{cases} \Rightarrow \text{don't expect mixing via decay}$$

- Mixing through oscillation \rightarrow large Δm



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- Standard Model result** → Significant contribution only from top loop

B_d^0 \bar{B}_d^0

B_s^0 \bar{B}_s^0

$$\Delta m \sim m_t^2 |V_{tb} V_{td}|^2 \sim m_t^2 \cdot O(\lambda^6)$$

$$\Delta m \sim m_t^2 |V_{tb} V_{ts}|^2 \sim m_t^2 \cdot O(\lambda^4)$$

Large $\Delta m_{s,d}$: $\Delta m_s \sim 1/\lambda^2 \Delta m_d \rightarrow B_s$ osc. is about 35 times faster than B_d osc.

2.3 Experimental observation of B meson mixing

for $\Gamma_H \approx \Gamma_L \approx \Gamma$

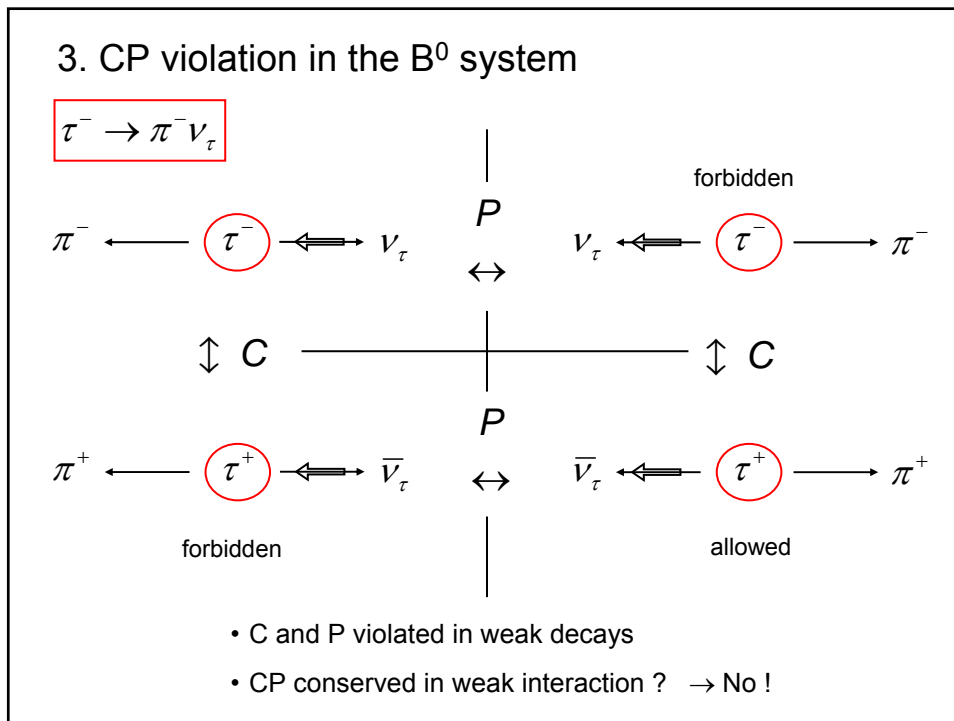
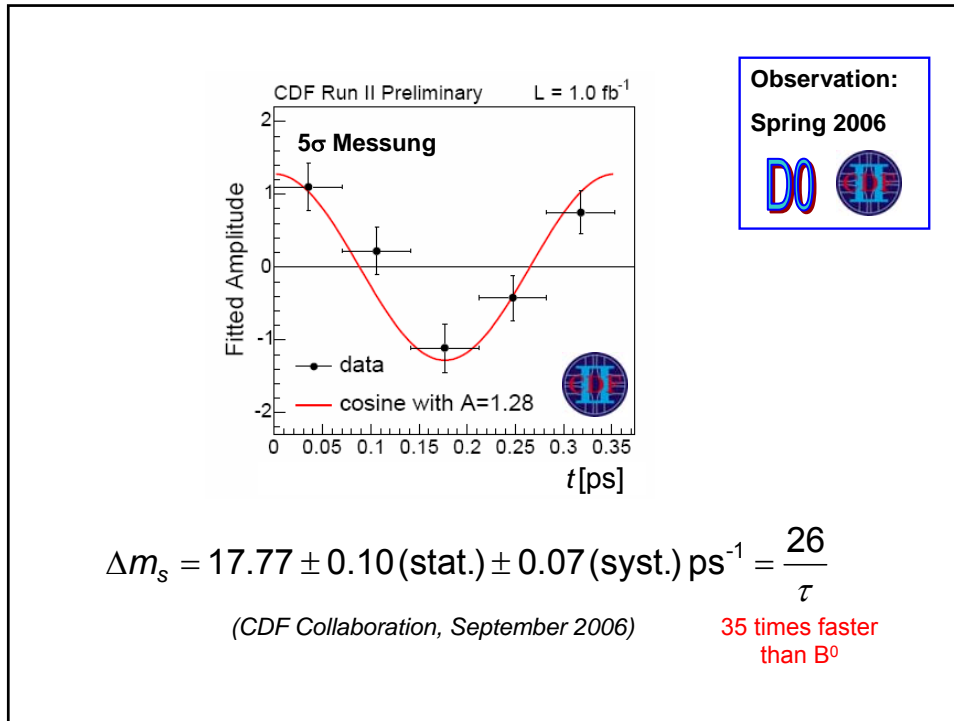
$P(B^0 \rightarrow B^0) = \frac{1}{2} \Gamma e^{-\Gamma t} (1 + \cos \Delta m t)$
 $P(B^0 \rightarrow \bar{B}^0) = \frac{1}{2} \Gamma e^{-\Gamma t} (1 - \cos \Delta m t)$

$\frac{P(B^0 \rightarrow B^0) - P(B^0 \rightarrow \bar{B}^0)}{P(B^0 \rightarrow B^0) + P(B^0 \rightarrow \bar{B}^0)}$

$$A = \frac{\text{unmixed} - \text{mixed}}{\text{unmixed} + \text{mixed}}$$

$\Delta m_d = 0.506 \pm 0.006 \pm 0.004 \text{ ps}^{-1}$
 $\approx \frac{0.774}{\tau_B}$

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3.1 CP Violation in Standard Model: complex CKM elements

CP (T) violation $\Leftrightarrow V_{ji} \neq V_{ji}^*$
i.e. Complex elements

Remark: For 2 quark generations the mixing is described by the **real 2x2** Cabbibo matrix \rightarrow **no CP violation !!**. To explain **CPV** in the SM Kobayashi and Maskawa have predicted a **third quark generation**.
Moreover, as can be shown, CPV requires that all u-type and all d-type quarks have different masses.

3.2 Unitarity Triangle

Unitary CKM matrix: $\mathbf{V}\mathbf{V}^\dagger = \mathbf{1}$ \rightarrow 6 "triangle" relations in complex plane:

$$V = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix}$$

$$V_{ud}V_{ub}^* + V_{cd}V_{cb}^* + V_{td}V_{tb}^* = 0$$

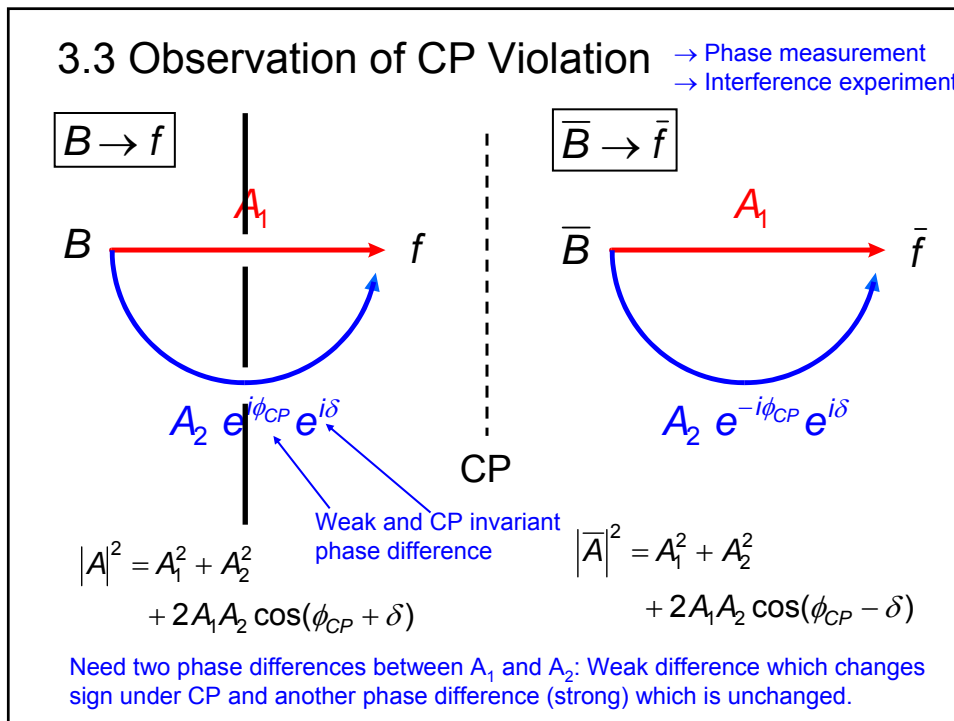
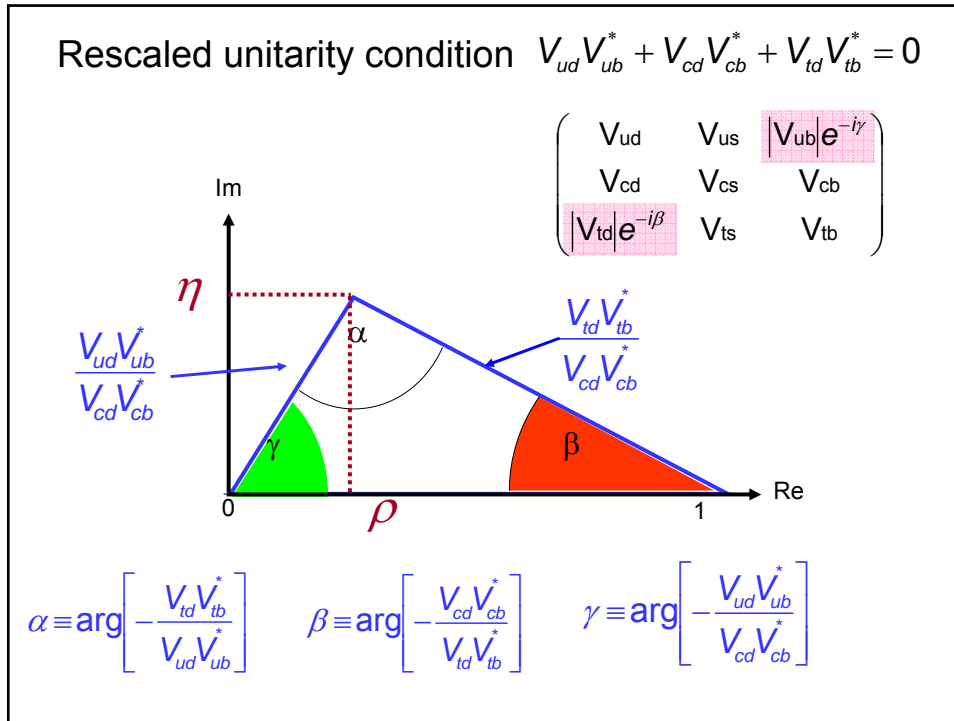
$$V_{td}V_{ud}^* + V_{ts}V_{us}^* + V_{tb}V_{ub}^* = 0$$

Important for \mathbf{B}_d and \mathbf{B}_s decays

Strength of CPV: Characterized by Jarlskog invariant $J = \text{Im} (V_{ij} V_{kl} V_{il}^* V_{kj}^*)$

In SM: $J = \text{Im}[V_{us} V_{cb} V_{ub}^* V_{cs}^*] = A^2 \lambda^6 \eta (1 - \lambda^2/2) + O(\lambda^{10}) \sim 10^{-5}$

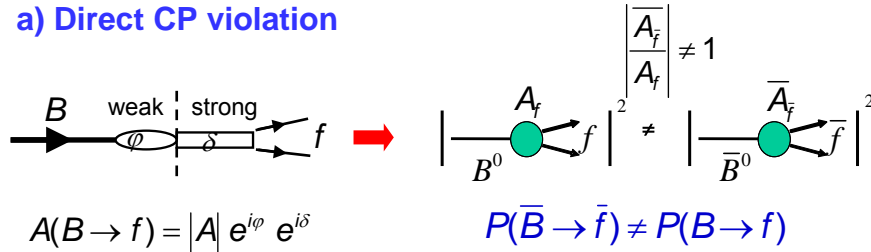
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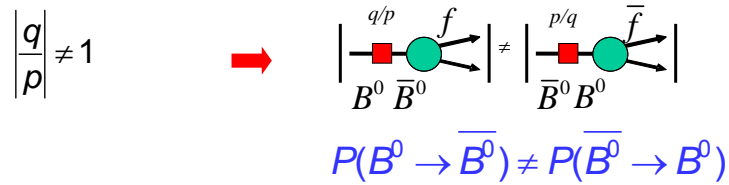
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“3 Ways” of CP violation in meson decays

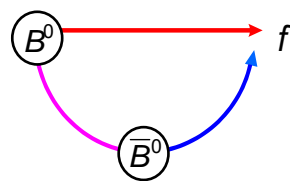
a) Direct CP violation



b) CP violation in mixing



c) CP violation through interference of mixed and unmixed amplitudes

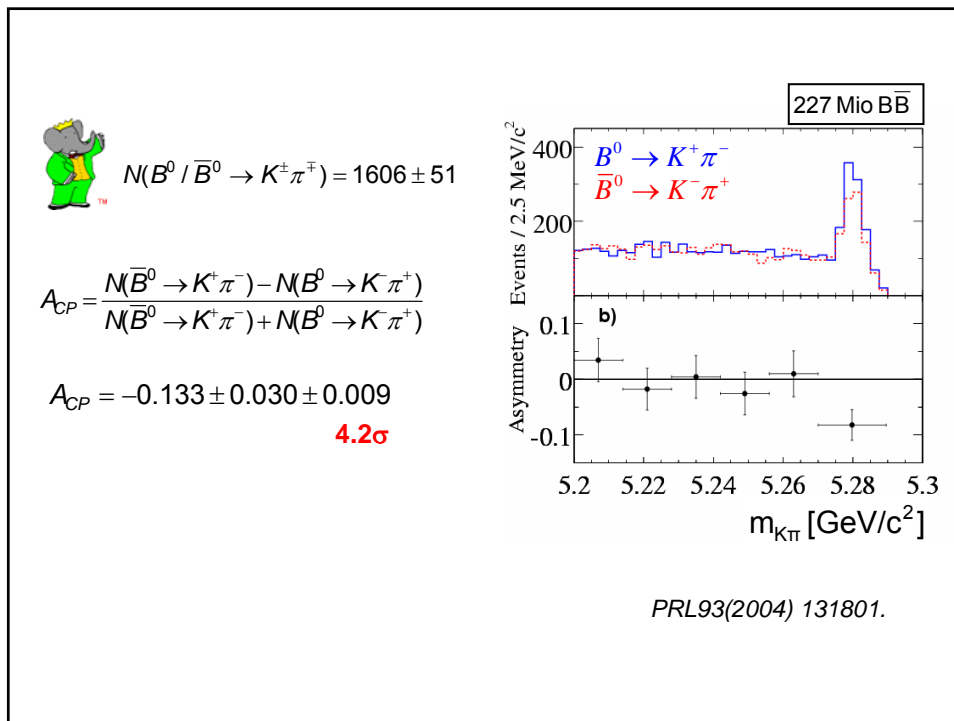
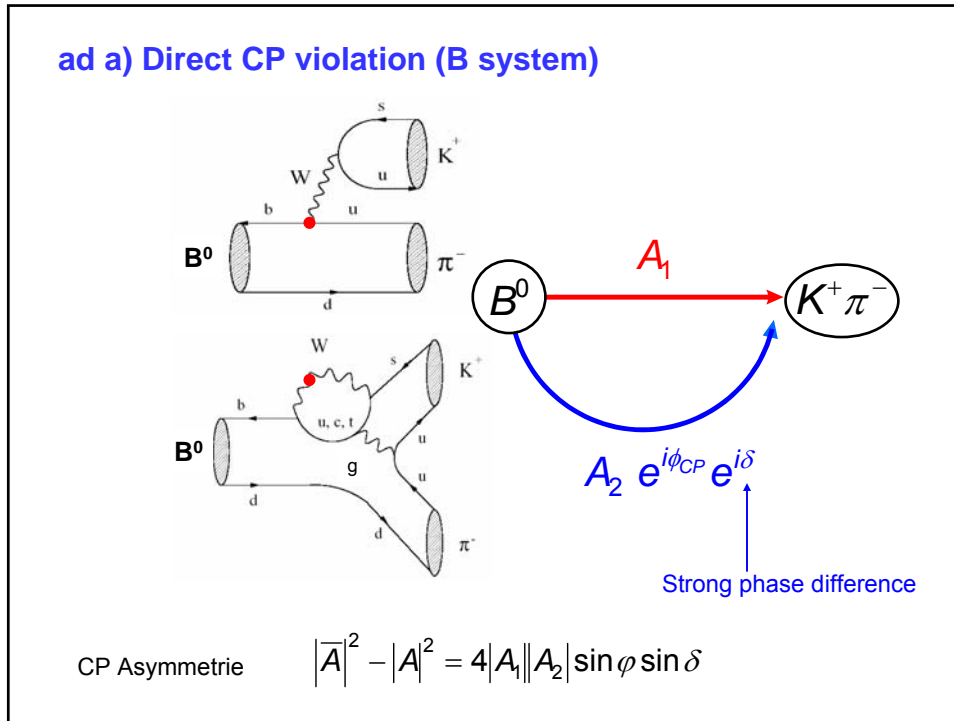


$$\Gamma(B_{t=0}^0 \rightarrow f)(t) \neq \Gamma(\bar{B}_{t=0}^0 \rightarrow f)(t)$$

Asymmetrie modulated by $\sim \sin \Delta m t$

Combinations of the 3 ways are possible!

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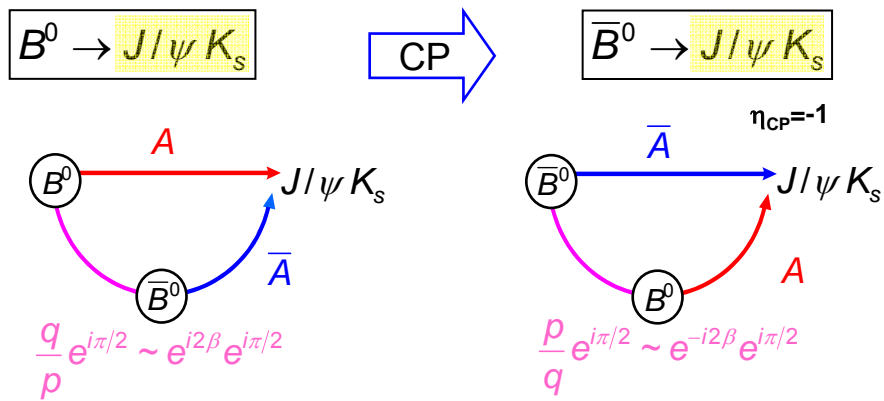
b) CP (T) violation in mixing

$$\left| \frac{q}{p} \right| \neq 1 \quad P(B^0 \rightarrow \bar{B}^0) \neq P(\bar{B}^0 \rightarrow B^0)$$

T violation

Skipped.

c) CP violation in interference between mixing and decay



$$|B^0 \rightarrow J/\psi K_s\rangle = A(f_+(t) + \lambda_{CP} f_-(t)) \quad | \bar{B}^0 \rightarrow J/\psi K_s \rangle = \bar{A} \left(f_+(t) + \frac{1}{\lambda_{CP}} f_-(t) \right)$$

$$\lambda_{CP} \equiv \frac{q}{p} \cdot \frac{\bar{A}}{A}$$

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SM prediction of λ_{CP} for $B^0 \rightarrow J/\psi K_s$ $\eta_{CP} = -1$

$$\lambda_{CP} = \frac{q \bar{A}}{p A} = \frac{V_{tb}^* V_{td} V_{cb} V_{cs}^* V_{cs} V_{cd}^*}{V_{tb} V_{td}^* V_{cb}^* V_{cs} V_{cs}^* V_{cd}} = \frac{V_{tb}^* V_{td} V_{cb} V_{cd}^*}{V_{tb} V_{td}^* V_{cb}^* V_{cd}} = e^{-2i\beta}$$

Beside V_{td} all other CKM elements are real

$$V_{td} \approx |V_{td}| e^{-i\beta} \Rightarrow \begin{cases} |\lambda_{CP}| = 1 \\ \text{Im}(\lambda_{CP}) = \sin(2\beta) \end{cases} \quad \text{no direct CPV, no CPV in mixing}$$

Calculation of the time-dependent CP asymmetry

$$\Gamma(B^0 \rightarrow f_{CP})(t) \propto \frac{e^{-\Delta t/\tau_{B^0}}}{(1+|\lambda_{CP}|^2)} \times \left[\frac{1+|\lambda_{CP}|^2}{2} - \text{Im}(\lambda_{CP}) \sin(\Delta m_d t) + \frac{1-|\lambda_{CP}|^2}{2} \cos(\Delta m_d t) \right]$$

$$\Gamma(\bar{B}^0 \rightarrow f_{CP})(t) \propto \frac{e^{-\Delta t/\tau_{B^0}}}{(1+|\lambda_{CP}|^2)} \times \left[\frac{1+|\lambda_{CP}|^2}{2} + \text{Im}(\lambda_{CP}) \sin(\Delta m_d t) - \frac{1-|\lambda_{CP}|^2}{2} \cos(\Delta m_d t) \right]$$

$$A_{CP}(t) = \frac{\Gamma(\bar{B}^0(t) \rightarrow f_{CP}) - \Gamma(B^0(t) \rightarrow f_{CP})}{\Gamma(\bar{B}^0(t) \rightarrow f_{CP}) + \Gamma(B^0(t) \rightarrow f_{CP})} = [S_f \sin(\Delta m_d t) - C_f \cos(\Delta m_d t)]$$

Time resolved

$$S_f = \frac{2 \text{Im} \lambda_{CP}}{1 + |\lambda_{CP}|^2} \quad C_f = \frac{1 - |\lambda_{CP}|^2}{1 + |\lambda_{CP}|^2}$$

Interference = $\sin 2\beta$ for $B^0 \rightarrow J/\psi K_S$

indicates direct CP violation if $|q/p| \neq 1$

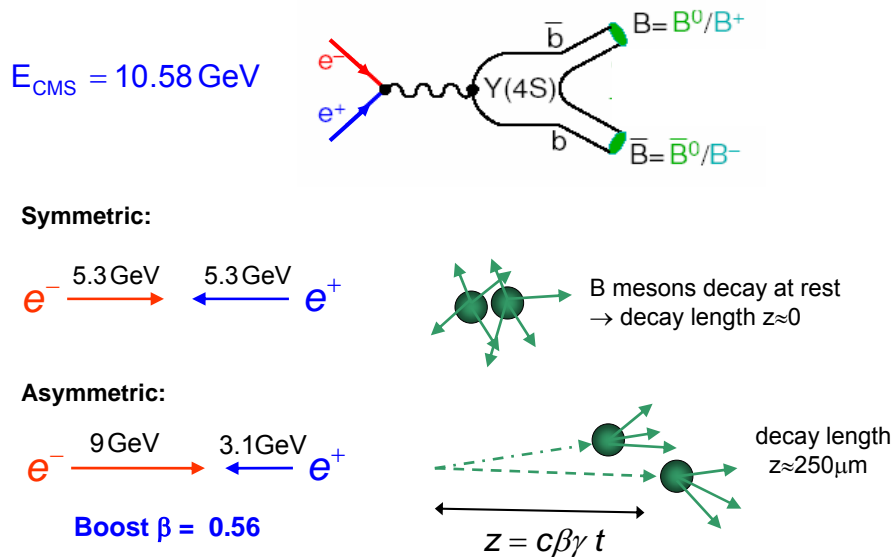
negligible

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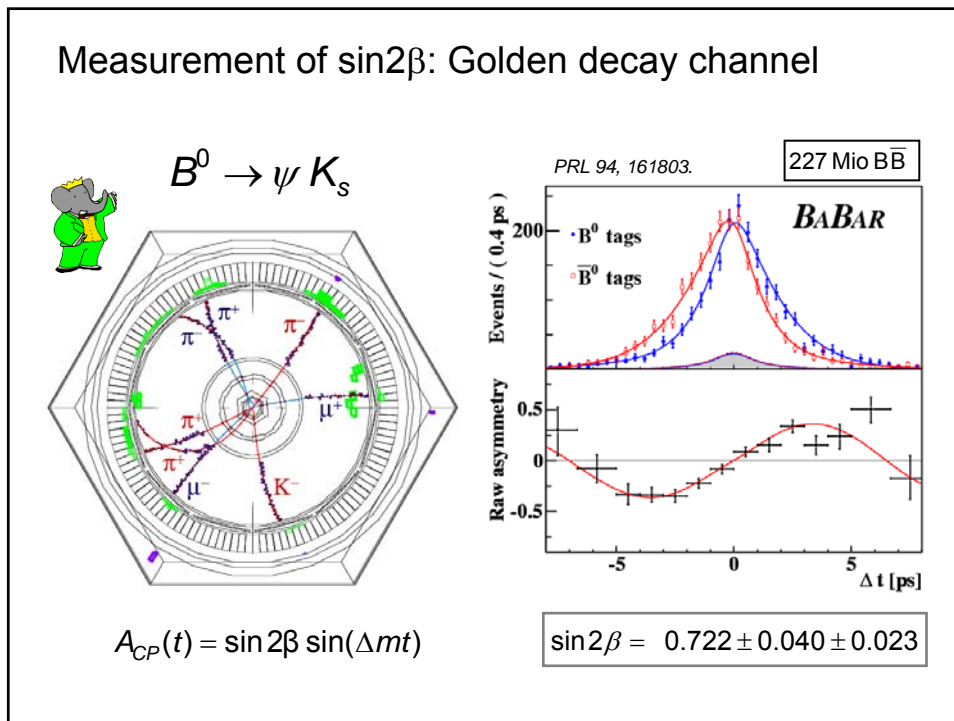
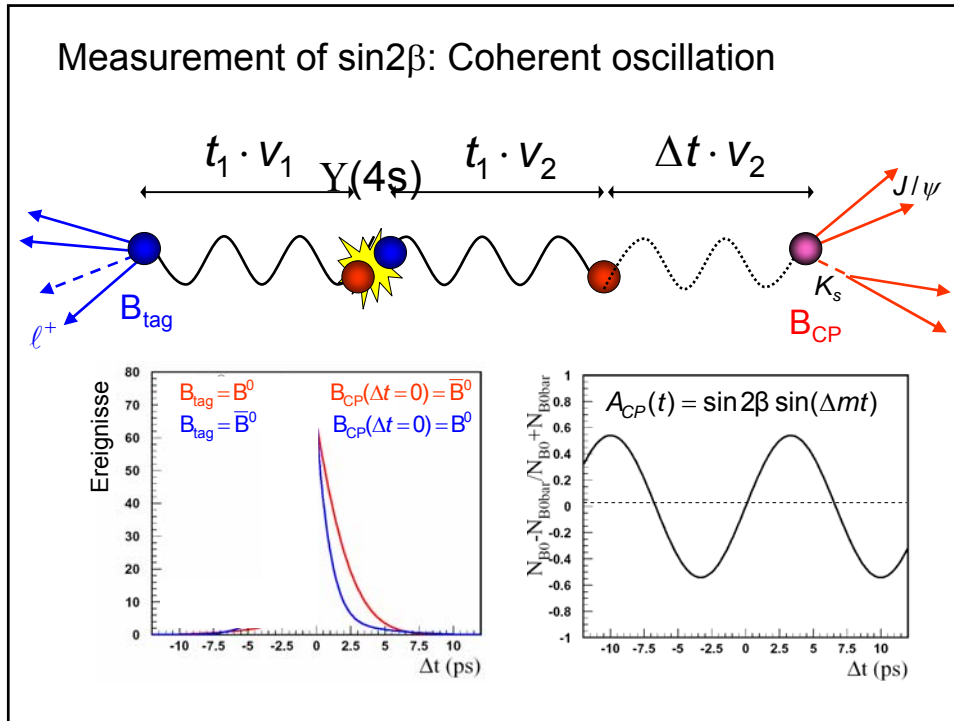
To measure CP violation in B_d system:

- Need many B (several 100×10^9)
- Need to know the flavor of the B at $t=0$
- Need to reconstruct the decay length to measure t

3.4 Measurement of $\sin 2\beta$: Asymmetric $e^+ e^-$ B factory

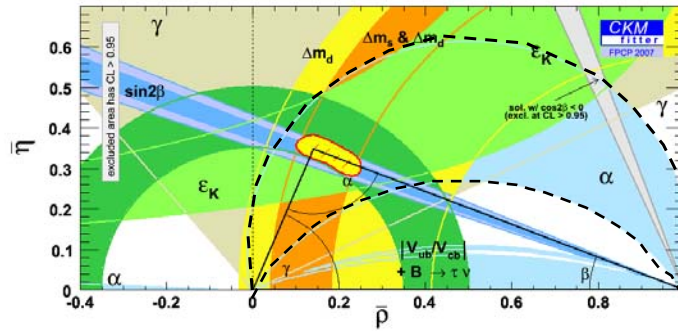


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3.5 Experimental status of the Unitarity Triangle



Standard Model CKM mechanism confirmed

1. Large CP Violation in B decays
 2. Large direct CP violation observed
 3. CPV parameter related to magnitude of non-CP observables
- A triple triumph**