Could the Dark Matter be the QCD Axion?

Guy D. Moore, TU Darmstadt

- Dark matter: how are we sure it's there?
- **T** symmetry in QCD: a mystery.
- The Axion: a possible explanation
- Cosmology of the axion: a dark matter candidate
- Relating the Dark Matter abundance to the Axion particle's mass
- The search for the axion: how knowing the mass helps

Dark Matter: a Cosmic Mystery



Atoms: Standard Model.
Dark Energy: Cosmological Constant.
Strange value, but possible
Dark Matter: MYSTERY! NOT SM!

We only know 3 things about dark matter:

- It's Matter: gravitationally clumps.
- It's **Dark**: negligible electric charge, interactions too feeble to be detected except by gravity
- It's Cold: negligible pressure by redshift z = 3000

Heidelberg, 6 Dez. 2019: Folie 2 von 36

Dark matter: are you sure?

Yes we are sure! Many independent lines of evidence!

- Microwave sky: pattern of hot/cold spots requires DM to explain (models without DM fail at the $> 100 \sigma$ level)
- X-ray appearance of rich galaxy clusters: 5× more total mass than "ordinary" (baryonic) mass
- large scale structure in distribution of galaxies
- Many other lines of evidence . . .

I will explain the galaxy cluster evidence as an example

Heidelberg, 6 Dez. 2019: Folie 3 von 36

Rich galaxy clusters



100's to 1000's of galaxies Gravitationally bound Full of plasma Plasma: hydrostatic equilibrium Glow in X-rays

Consider X-ray signal more closely

Heidelberg, 6 Dez. 2019: Folie 4 von 36

X-rays from galaxy clusters

Mohr et al ApJ 517, 627; Vikhlinin et al astro-ph/0507092; Allen et al astro-ph/0405340...

X rays emitted when e^- , p in plasma scatter inelastically



2-body scattering: Rate $\propto \rho^2$ density². Spectrum set by e^- velocity Hence by the temperature

Intensity: determines $M_{\rm plasma} \gg M_{\rm galaxies}$. Spectrum + Hydrostatic equil.: determines $M_{\rm grav} = M_{\rm tot}$. Total (grav.) mass $\simeq 5 \times$ plasma mass \gg mass of galaxies

Heidelberg, 6 Dez. 2019: Folie 5 von 36

Really sure? Look at cluster 2 other ways



Left: grav. lensing gives independent measure of M_{grav} Right: Sunyaev-Zeldovich effect measures M_{plasma}

So Dark Matter really is there!

Heidelberg, 6 Dez. 2019: Folie 6 von 36

Changing the subject: T symmetry

Obviously wrong! Time reversed! "Entropic Arrow of Time" Statistical Physics

Microphysics is OK: E&M, Strong force T invariant

Violated by tiny weak-force effects, orders of magnitude smaller than other weak physics

Heidelberg, 6 Dez. 2019: Folie 7 von 36

T in E&M

How do E, B fields change when you run movie backwards?



Q's unchanged, but J's flip. E same, but B flips.

Heidelberg, 6 Dez. 2019: Folie 8 von 36

T in E&M or QCD

How a particle physicist thinks about E&M / QCD:

- What are the *Degrees of Freedom*? Electrons/quarks $\psi_e/\psi_{f,a}$, photons/gluons A^{μ}_A
- What are the *Symmetries*? Lorentz invariance, gauge invariance
- What is the *Action*?

$$S = \int d^4x \ \mathcal{L}[A^{\mu}_{A}, \psi_{e}, \psi_{f,a}]$$

• What is the most general \mathcal{L} , given fields+symmetries?

Heidelberg, 6 Dez. 2019: Folie 9 von 36

T in E&M, QCD Continued

Lorentz + gauge symmetry: \mathcal{L} is gauge-invariant scalar. Without prejudice, write the most general answer you can!

$$\mathcal{L} = \frac{1}{4e^2 \ 4g^2} \left(\vec{E_a}^2 - \vec{B_a}^2 \right) + \frac{\Theta}{8\pi^2} \left(\vec{E_a} \cdot \vec{B_a} \right) + \mathcal{L}_{\psi} \,.$$

Here $\vec{E}_i^a = -\partial_t A_i^a - \partial_i \Phi^a + g f_{abc} A_i^b \Phi^c$ and \vec{B} is ...

Both respect gauge, Lorentz invariance \rightarrow should be OK! First term: **T** respecting. Second term: **T** violating.

Heidelberg, 6 Dez. 2019: Folie 10 von 36

Time in E&M and QCD

In E&M, $\vec{E} \cdot \vec{B}$ turns out not to affect equations of motion – in fact, has no effect at all!

But QCD has these nonlinearities $E_i^a = \dots + g f_{abc} A_i^b \Phi^c$, same argument does not apply.

QCD with $\Theta \neq 0$ is *not* equivalent to $\Theta = 0$ They have different *vacuum energy* The $\Theta \neq 0$ version is **T** violating!

So, what do we know *experimentally*?

Heidelberg, 6 Dez. 2019: Folie 11 von 36

Neutron Electric Dipole Moment

A Test of whether **QCD** obeys T symmetry:



Put neutron in \vec{B} field – spin lines up with \vec{B} . Does it have an electric dipole moment aligned with spin?

If so: physics when you run movie backwards is different! \mathbf{T} violating!

Heidelberg, 6 Dez. 2019: Folie 12 von 36

Neutron Electric Dipole Moment

Theory: Neutron electric dipole moment should exist,

$$d_n \simeq -2 \times 10^{-16} \, e \, \mathrm{cm} \times \Theta$$

So long as Θ is not zero! See arXiv:1904.00323, assumes Θ , modulo 2π , is small

Experiment: Consistent with zero! Baker et al (Grenoble), arXiv:hep-ex/0602020

 $|d_n| < 2.9 \times 10^{-26} \ e \ \mathrm{cm}$

Either $|\Theta| < 10^{-10}$ by (coincidence? accident?) or there is something deep going on here.

Heidelberg, 6 Dez. 2019: Folie 13 von 36

Axion: an explanation for $\Theta = 0$

Hypothesize an extra *complex scalar* field $\varphi = \phi e^{i\theta_A}$:

- Field: takes a value at each point in space Think \vec{E}
- Scalar: value is a number, without direction but with units ...
- Complex scalar: value is a real and imaginary part

Assume a symmetry: $\varphi \rightarrow e^{i\theta}\varphi$

Energy should be

Energy =
$$\int d^4x \ |\dot{\varphi}|^2 + |\vec{\nabla}\varphi|^2 + V(|\varphi|)$$

Heidelberg, 6 Dez. 2019: Folie 14 von 36

Axion: spontaneous symmetry breaking

Potential function $V(|\varphi|)$ can look like this:



Lowest energy state has $\varphi = f_a \neq 0$. Minimum not unique; all values of θ_A equally good. *Or are they?*?

Indirect Interactions with QCD

Peccei, Quinn; Kim, Shifman Vainshtein Zakharov Carefully chosen high-scale physics couples the phase of this field to $\vec{E}_a \cdot \vec{B}_a$:

$$\mathcal{L} = \frac{1}{4g^2} \left(\vec{E}_a^2 - \vec{B}_a^2 \right) + \frac{\Theta}{8\pi^2} \left(\vec{E}_a \cdot \vec{B}_a \right) + \frac{\theta_A}{8\pi^2} \left(\vec{E}_a \cdot \vec{B}_a \right)$$

T violation determined by $\Theta_{\text{eff}} = \Theta + \theta_A \pmod{2\pi}$ Now φ can change values, changing Θ_{eff} . Vacuum energy lower where $\Theta_{\text{eff}} = 0$ Dynamics choose T respecting state! Cosmological evolution of φ

Suppose φ starts with some value cosmologically:

Sinks to \mathbf{T} respecting minimum, but (Hubble damped) oscillations around that minimum. Oscillations act like matter – dark matter!

Heidelberg, 6 Dez. 2019: Folie 17 von 36

Different values at different space points?

What if axion field starts at a (random) different value at different points in space? actually likely

- Energy cost to vary through space: $\int d^3x |\nabla \varphi|^2/2$
- Locally aligns to point in one direction
- Causality: cannot align globally
- Phase-ordering dynamics

Predicting the axion mass

Suppose I could solve these *phase-ordering dynamics* Count axions at the end: *predict* Dark Matter density? No, depends on φ vacuum-value f_a :

 $\rho_{\rm DM} \propto f_a$

However the unknown axion mass also depends on f_a :

$$m_a^2 = \frac{\chi_{\text{QCD}}}{f_a^2}, \qquad \chi_{\text{QCD}} \simeq (76 \pm 1 \text{ MeV})^4$$

If the DM is axions, computing axion production efficiency predicts the axion mass.

Heidelberg, 6 Dez. 2019: Folie 19 von 36

Needs a numerical solution



 $\chi(T)$ from Borsanyi et al 1606.07494

Heidelberg, 6 Dez. 2019: Folie 20 von 36

Axions and Topology

 φ is a complex number – plot as a 2D arrow. Axion field: a field of arrows. 2D slice for instance:



Field generically has vortices Davis, PLB180 225 (1986)

Heidelberg, 6 Dez. 2019: Folie 21 von 36

Domain walls

2D slice of evolution, When the potential tilts:

Heidelberg, 6 Dez. 2019: Folie 22 von 36

Network evolution

When potential stays untilted:

When potential tilts:

Heidelberg, 6 Dez. 2019: Folie 23 von 36

Layers of String Energy

$$E_{\rm str} = \int dz \int d\phi \int r \, dr \left(\nabla \phi^* \nabla \phi \simeq f_a^2 / 2r^2 \right) \simeq \pi \ell f_a^2 \int_{\sim f_a^{-1}}^{\sim H^{-1}} \frac{r \, dr}{r^2}$$



Series of "sheaths" around string: equal energy in each $\times 2$ scale, 10^{30} scale range! $\ln(10^{30})\simeq 70$. Log-large string tension $T_{\rm str}=\pi f_a^2\ln(10^{30})\equiv \pi f_a^2\kappa$

Not reproduced by numerics (separation/core \sim 400)

Heidelberg, 6 Dez. 2019: Folie 24 von 36

Another model: Abelian Higgs model

Complex scalar plus copy of electromagnetism

$$\mathcal{L}(\varphi, A_{\mu}) = \frac{1}{4} (\partial_{\mu} A_{\nu} - \partial_{\nu} A_{\mu})^2 + (D_{\mu} \varphi)^* (D^{\mu} \varphi) + \frac{\lambda}{8} \left(2\varphi^* \varphi - f_a^2 \right)^2$$

with $D_{\mu} = \partial_{\mu} - ieA_{\mu}$ covariant derivative Relativistic version of \vec{B} fields in type-II superconductor Tension-only (magnetic-flux) strings Only massive fields outside of the string cores

Hybrid: global strings, local cores

Theory with one A_{μ} and two scalars

$$\mathcal{L}(\varphi_{1},\varphi_{2},A_{\mu}) = \frac{1}{4} (\partial_{\mu}A_{\nu} - \partial_{\nu}A_{\mu})^{2} + \frac{\lambda}{8} \left[(2\varphi_{1}^{*}\varphi_{1} - f^{2})^{2} + (2\varphi_{2}^{*}\varphi_{2} - f^{2})^{2} \right] + |(\partial_{\mu} - iq_{1}eA_{\mu})\varphi_{1}|^{2} + |(\partial_{\mu} - iq_{2}eA_{\mu})\varphi_{2}|^{2}$$

Pick $q_1 \neq q_2$, say, $q_1 = 4$, $q_2 = 3$.

One *global* [axions], one *local* rotation symmetry *"glues"* Abelian-Higgs string onto axion string core.

Extra tension represents effects of short-distance scales which give extra tension

Increases string tension by a completely tuneable amount!

Heidelberg, 6 Dez. 2019: Folie 26 von 36

Simulations with 4 different added tensions

Higher tension = higher initial density, longer lasting, hardier loops

Heidelberg, 6 Dez. 2019: Folie 27 von 36

Results



Axions produced vary mildly with increasing string tension

Heidelberg, 6 Dez. 2019: Folie 28 von 36

Put it all together

$$\begin{array}{lll} \mbox{Axion production:} & n_{\rm ax}(T=T_*)=(13\pm2)H(T_*)f_a^2\\ & \mbox{Hubble law:} & H^2=\frac{8\pi\varepsilon}{3m_{\rm pl}^2}\,,\\ \mbox{Equation of state:} & \varepsilon=\frac{\pi^2T^4g_*}{30}\,, \quad s=\frac{4\varepsilon}{3T}\,, \quad g_*(1{\rm GeV})\simeq73\\ & \mbox{Susceptibility:} & \chi(T)\simeq \left(\frac{1\ {\rm GeV}}{T}\right)^{7.6}\,(1.02(35)\times10^{-11}{\rm GeV^4})\\ & \mbox{Dark matter:} & \frac{\rho}{s}=0.39\ {\rm eV} \end{array}$$

One finds $T_* = 1.54\,{
m GeV}$ and $m_a = 26.2\pm 3.4\,\mu{
m eV}$

Heidelberg, 6 Dez. 2019: Folie 29 von 36

Summary so far

- If the QCD axion exists (solving $\Theta \mathbf{T}$ problem)
- If the axion is the Dark Matter
- If it starts with random values at different space-points

Then it has $m_A = 26.2 \pm 3.4 \ \mu \text{eV} (\simeq 6 \text{ Ghz})$

So, what are constraints and detection prospects?

Heidelberg, 6 Dez. 2019: Folie 30 von 36

Looking for DM axions today

Axion field still fluctuating today. Energy density is:

$$\varepsilon_{\rm DM} = \frac{\chi}{2} \theta_{\rm max}^2 \Rightarrow \frac{(76 \,\,{\rm MeV})^4}{2} \theta_{\rm max}^2 = \frac{0.3 \,\,{\rm GeV}}{{\rm cm}^3}$$

 $\theta_{\rm max} = 4 \times 10^{-19}$

Expected $\theta_A G_{\mu\nu} \tilde{G}^{\mu\nu}$ effect: neutron EDM oscillates by $10^{-34} ecm$ at 6 GhzGood luck finding that! Need to consider other couplings

Axions and Electromagnetism

Axions generically also couple to E&M:

$$\mathcal{L} = \ldots + \frac{\theta_A}{8\pi^2} \left(\vec{E_a} \cdot \vec{B_a} \right) + \frac{C\alpha_{\rm EM}}{2\pi} \left(\vec{E_{\rm EM}} \cdot \vec{B_{\rm EM}} \right)$$

with C model dependent (but never 0!). Changes EM behavior: Ampere's law becomes

$$\vec{\nabla} \times \vec{B} = \vec{J} + \kappa \partial_0 \vec{E} - \frac{C\alpha_{\rm EM}}{2\pi} (\partial_0 \theta_A) \vec{B}$$

Turns \vec{B} field into effective oscillating current! Nice review article: Redondo et al arXiv:1801.08127

Heidelberg, 6 Dez. 2019: Folie 32 von 36

How to turn axions into microwaves

Consider dielectric-vacuum interface, B tangent to surface

$$\dot{E}_{\parallel} = \frac{-g_{a\gamma}\dot{\theta}_{A}B_{\text{static}} + \nabla \times B}{\kappa} \qquad \dot{E}_{\parallel} = -\gamma_{a\gamma}\dot{\theta}_{A}B_{\text{static}} + \nabla \times B$$

As usual, E_{\parallel} must match at boundary Impossible without $\nabla \times B$ traveling wave component! Interface, bathed in *B*-field, emits microwaves with $\omega = m_a$

Heidelberg, 6 Dez. 2019: Folie 33 von 36

Resonant cavity detection



Conductor: $\kappa \simeq i\infty$ Separate by $\lambda/2$: resonant growth of wave

Power produced enhanced by Q. Noise reduced 1/QBandwidth $\propto 1/Q$: need Q "experiments" Ability to search enhanced by one factor of Q

Heidelberg, 6 Dez. 2019: Folie 34 von 36

Alternative: MADMAX

Parallel dielectric plates B field: microwaves.. Emission from each plate adds coherently if optical dist. between plates = λ



Plate separations can be actuated to tune sensitive λ Loses Q but gains volume $\sim m^3$ $_{\rm arXiv:1901.07401}$

Heidelberg, 6 Dez. 2019: Folie 35 von 36

Conclusions

- Two mysteries: Dark Matter and \mathbf{T} symmetry
- Axion could explain T invariance of QCD
- Axion could also be dark matter
- Nontrivial, but now solvable, early-Universe dynamics
- Prediction: $m_a = 26 \pm 3 \,\mu \text{eV}$
- The search is on. Mass window helps design experiments and narrow the search.

What about Anthropic Principle?

Trendy Explanation for "coincidences" or "tunings"

Why is Cosmological Constant so small? If it were 100 times bigger, matter would fly apart or collapse before life could evolve. Nature plays dice, universes with all values occur, but only universes with life get observed.

Why does QCD respect T symmetry? If QCD violated T, something would go wrong with nuclear physics, which would make life impossible. Nature plays dice, only universes where life is possible get observed. Except that life is fine in a world where $\Theta = 10^{-2}$!

Heidelberg, 6 Dez. 2019: Folie 37 von 36

Why should the axion-EM coupling be nonzero? Dynamics of QCD at the QCD scale are quite nontrivial. The π^0 also has a coupling

$$\mathcal{L} = \ldots + \frac{\alpha_{\rm EM} N_c}{3\pi F_{\pi}} \pi^0 \vec{E} \cdot \vec{B}$$

The axion then *mixes* with π^0 in a way dependent on the m_u/m_d mass ratio. The above induces an axion-photon coupling

$$\frac{C\alpha_{\rm EM}}{2\pi}\theta_A \vec{E} \cdot \vec{B} \qquad \text{with} \qquad C = 1.92$$

For the axion-photon coupling to vanish, the "fundamental" contribution would have to be exactly -1.92.

Heidelberg, 6 Dez. 2019: Folie 38 von 36

Why random initial conditions for θ_A ?

If Universe ever reached $T > f_a \sim 10^{11} \text{ GeV}$: axion field got "symmetry restored," then randomly broke in different directions in different places.

If not: starting θ_A value "came out of inflation"

- Inflation stretches quantum fluctuations to classical ones: $\Delta \varphi \sim H_{\text{infl.}}$.
- If $N_{\rm efolds}H^2 > f_a^2$, scambles field.
- If not: need $H < 10^{-5} f_a$ to avoid excess "isocurvature" fluctuations in axion field