

The U-Spin-CP Anomaly in charm

based on 2210.16330

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Outline

- ① Introduction
- ② Model
- ③ Phenomenology & Constraints
- ④ Conclusion

Introduction

CP Violation in charm: Experimental Status

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- ▶ LHCb'19: First observation of CP violation in charm sector [1903.08726]
⇒ Difference of CP asymmetries in $D^0 \rightarrow K^+ K^-, \pi^+ \pi^-$ decays

$$\Delta A_{CP} = A_{CP}(K^+ K^-) - A_{CP}(\pi^+ \pi^-) = (-15.4 \pm 2.9) \cdot 10^{-4}$$

- ▶ LHCb'22: Direct measurement of $A_{CP}(K^+ K^-)$ [2209.03179]

$$A_{CP}(K^+ K^-) = (6.8 \pm 5.4 \pm 1.6) \cdot 10^{-4}$$

- ▶ LHCb'22: Fit of both direct CP asymmetries [2209.03179]

$$a_{K^+ K^-}^d = (7.7 \pm 5.7) \cdot 10^{-4}$$

$$a_{\pi^+ \pi^-}^d = (23.2 \pm 6.1) \cdot 10^{-4}$$

⇒ 3.8σ evidence for direct CP violation in $D^0 \rightarrow \pi^+ \pi^-$!

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- ▶ $a_{\pi^-\pi^+}^d > |\Delta A_{CP}|$ is quite large!
- ⇒ SM interpretation: $a_{\pi^-\pi^+}^{d\text{SM}} \sim 2 \cdot \text{Im}(V_{cb}^* V_{ub} / (V_{cd}^* V_{ud})) h/t$
- ⇒ Would require enhanced higher order contributions h over tree-level ones t by $h/t \sim 2$ to compensate CKM suppression!

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- ▶ The SM has an approximate U-spin symmetry ($d \leftrightarrow s$)
 $\Rightarrow$ U-Spin limit $a_{K^+K^-}^d = -a_{\pi^+\pi^-}^d$ broken by $2.7 \sigma!$ [2207.08539]
- ▶ U-spin breaking in SM: order $\frac{m_s - m_d}{\Lambda_{QCD}} \sim 30\%$
 $\Rightarrow$ sufficient to explain $\frac{\mathcal{B}(D \rightarrow \pi^+\pi^-)}{\mathcal{B}(D \rightarrow K^+K^-)} \sim 2.8$
 $\Rightarrow$ insufficient to explain $a_{\pi^-\pi^+}^{d,\text{exp}}, a_{K^-K^+}^{d,\text{exp}}$ by a factor 4-5!

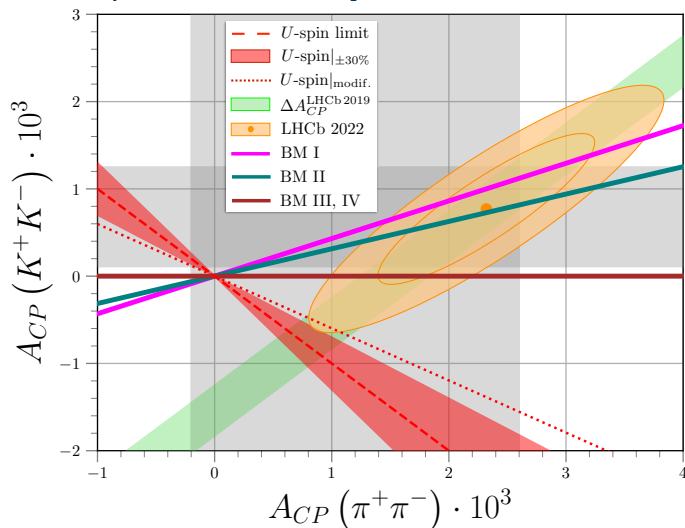
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$\Rightarrow$ **Hint for CP- and U-Spin violating NP!**

The U-Spin-CP Anomaly



Model

How to adress the anomaly?

Idea: Z' with generation dependent charges

$$D^0 \left\{ \begin{array}{l} c \rightarrow u \\ \bar{u} \end{array} \right\} \left\{ \begin{array}{l} \pi^+ \\ \pi^- \end{array} \right\}$$

$$D^0 \left\{ \begin{array}{l} c \rightarrow u \\ \bar{u} \end{array} \right\} \left\{ \begin{array}{l} K^+ \\ K^- \end{array} \right\}$$

$$a_{\pi^-\pi^+}^d = \frac{g_4^2}{M_{Z'}^2} \Delta \tilde{F}_R [c_\pi F_{Q1} + d_\pi F_{d1}], \quad a_{K^-K^+}^d = \frac{g_4^2}{M_{Z'}^2} \Delta \tilde{F}_R [c_K F_{Q2} + d_K F_{d2}]$$

- ▶ $g_4, M_{Z'}$ are the $U(1)'$ coupling and Z' -mass
- ▶ F_{ψ_i} are SM fermion $U(1)'$ charges
- ▶ $c_{\pi,K}, d_{\pi,K}$ are hadronic parameters (contain HMEs & RG effects)

Benchmarks & Ingredients

- ▶ $\Delta \tilde{F}_R = \sin \theta_u \cos \theta_u (F_{u_2} - F_{u_1}) \neq 0$ to induce ucZ' -vertex
 $\Rightarrow$ Right-handed c - u -mixing by angle $\theta_u \neq 0$ (adjust)
 $\Rightarrow$ different $U(1)'$ charges $F_{u_2} \neq F_{u_1}$
- ▶ $F_{Q_{1,2}} = 0$ to avoid Kaon constraints
- ▶ $\frac{F_{d_2}}{F_{d_1}} = \frac{d_\pi a_{K^-K^+}^d}{d_K a_{\pi^-\pi^+}^d}$ to generate $a_{K^+K^-, \pi^+\pi^-}^{d,\text{exp}} \Rightarrow |F_{d_2}| \ll |F_{d_1}|$
- ▶ six anomaly cancellation conditions on F_{ψ_i}
 $\Rightarrow$ add ν_R with $F_{\nu_i} \neq 0$

Model	F_{Q_i}			F_{u_i}			F_{d_i}			F_{L_i}			F_{e_i}			F_{ν_i}		
BM III	0	0	0	G	$-F$	0	F	$-G$	0	0	0	0	$-G$	F	0	G	$-F$	
BM IV	0	0	0	$-F_u$	F_u	0	F_d	0	$-F_d$	0	0	0	F_e	0	$-F_e$	F_ν	$-F_\nu$	0

Phenomenology & Constraints

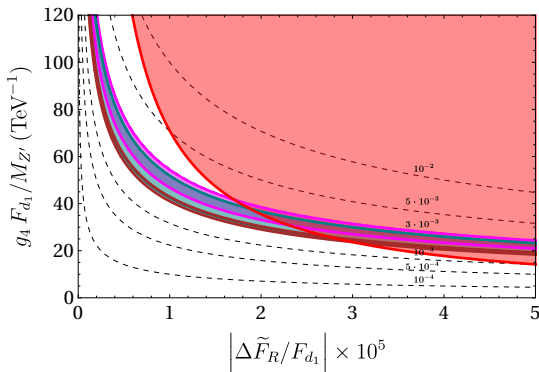
Predicting $A_{CP}(\pi^0\pi^0)$ and $A_{CP}(\pi^0\pi^+)$

- ▶ $F_{d_1} \neq F_{u_1}$ violates Isospin and induces $A_{CP}(\pi^0\pi^{0/+}) \neq 0$
- ▶ Similar size: $\frac{A_{CP}(\pi^0\pi^0)}{A_{CP}(\pi^+\pi^0)} = \frac{d_{\pi^0}}{d_{\pi'}} \simeq 1.08 \pm 0.10$ [2004.01206]
- ▶ For $|F_{d_2}| \ll |F_{d_1}|$ all CP asymmetries involving pions are correlated!
 $A_{CP}(\pi^+\pi^0) \simeq \frac{d_{\pi'}}{d_{\pi^0}} A_{CP}(\pi^0\pi^0) \simeq -\frac{d_{\pi'}}{d_{\pi}} \left(1 - \frac{F_{u_1}}{F_{d_1}}\right) \Delta A_{CP}$
- ▶ $d_{\pi} \sim d_{\pi'} \sim d_{\pi^0}$ and in our BMs $|F_{u_1}| < |F_{d_1}|$
 $\Rightarrow A_{CP}(\pi^0\pi^{0/+})$ are positive
 $\Rightarrow A_{CP}(\pi^0\pi^{0/+})$ are of similar magnitude as $\Delta A_{CP} \sim 10^{-3}$

$$A_{CP}^{\text{BM III}}(\pi^0\pi^{0/+}) \simeq -1 \cdot \Delta A_{CP}$$

$$A_{CP}^{\text{BM IV}}(\pi^0\pi^{0/+}) \simeq -\left(1 \pm \frac{1}{\sqrt{2}}\right) \cdot \Delta A_{CP}$$

D-Mixing



$$\frac{g_4}{M_{Z'}} \Delta \tilde{F}_R < 7.1 \cdot 10^{-4} \text{ TeV}^{-1} \text{ from } D\text{-mixing (e.g. LHCb'21 [2106.03744])}$$

$$\left| \frac{\Delta \tilde{F}_R}{F_{d_1}} \right| \lesssim \mathcal{O}(10^{-5}) \quad \Rightarrow \quad \boxed{\theta_u \ll 1}$$

$$\frac{g_4 F_{d_1}}{M_{Z'}} \gtrsim \frac{1}{25 \text{ GeV}} \cdot \frac{|a_{\pi^+ \pi^-}^d|}{0.002} \quad \Rightarrow \quad \boxed{\text{light } Z' \text{ (or large } g_4 F_{d_1})!}$$

Charm dilepton & invisibles data

Lepton couplings of the Z' are constrained by various processes

- ▶ BR's of (semi-)muonic D -decays (LHCb talk @ ICHEP'22)
 $\Rightarrow D^0 \rightarrow \mu^+ \mu^-$, $D^0 \rightarrow \pi^0 \mu^+ \mu^-$, $D^+ \rightarrow \pi^+ \mu^+ \mu^-$
- ▶ Drell-Yan data for $\ell = e, \tau$ [2003.12421]

$\Rightarrow$ Lepton charges should not be excessive compared to quarks

$\Rightarrow$ very roughly $|F_{L_i}, F_{e_i}| \lesssim |F_{d_1}|$

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Z' couples to ν_R (+ potentially SM singlet vector-like fermions χ)

$\Rightarrow$ Couplings constrained by charm BR's with invisibles in final state!

- ▶ $\mathcal{B}(D^0 \rightarrow \pi^0 \text{ inv.}) < 2.1 \cdot 10^{-4}$ (BESIII'21 [2112.14236]) constrains $|F_{\nu, \chi}|$
- ▶ $\mathcal{B}(D^0 \rightarrow \text{inv.}) < 9.4 \cdot 10^{-5}$ (BELLE'16 [1611.09455]) constrains $|F_\nu|$

$$\Rightarrow |F_{\nu, \chi}| \lesssim 110 |F_{d_1}|$$

A flavorful Z' of $\mathcal{O}(10 \text{ GeV})$?

Mass window severely constrained!

- Dijets + ISR (CMS'19 [1905.10331])

$$\Rightarrow g_4 F_{d_1} \lesssim 0.5 \quad \text{implying} \quad 10 \text{ GeV} \lesssim M_{Z'} \lesssim 20 \text{ GeV}$$

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Z' - γ kinetic mixing induces Z' -fermion-couplings via $\mathcal{L}_\epsilon = -\epsilon e J^\mu Z'_\mu$

- Limits from dark photon searches! $\Rightarrow \epsilon(M_{Z'}) \lesssim 10^{-3}$

$$\Rightarrow Z' \rightarrow \mu^+ \mu^- \text{ (LHCb'19 [1910.06926])}, Z' \rightarrow e^+ e^- \text{ [1801.04847]}$$

- Translate to bounds on lepton charges

$\Rightarrow$ **The Z' has to be (quasi-)leptophobic!**

$$\frac{F_{L_{1,2}, e_{1,2}}}{F_{d_1}} \lesssim \frac{1}{750}$$

$\Rightarrow$ much stronger than constraints from DY & semileptonic decays

Conclusion

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- ▶ $a_{\pi^+\pi^-}^d$ factor ~ 2 enhanced over SM tree-level (LHCb'22)
- ▶ U-Spin breaking ($\sim 4 - 5$) times larger than in SM (LHCb'22)
- ▶ Explanation by light, flavorful Z' (very predictive!)
 $\Rightarrow M_{Z'} = 10\text{-}20$ GeV, μ -,e-phobic, coupling only to RH fermions
 $\Rightarrow$ Up to now the only available explanation for the anomaly!

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- ▶ $a_{\pi^+\pi^-}^d$ factor ~ 2 enhanced over SM tree-level (LHCb'22)
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What to do next?

- ▶ Many search channels & signatures (also for LHCb)
 $\Rightarrow A_{CP}(\pi^0\pi^{0/+})$, dijets, invisible D -decays, τ Drell-Yan
- ▶ Relation to U-Spin puzzle in B -Decays? [2211.06994]
- ▶ Reduce large hadronic uncertainties (rescattering) [2203.04056], [PVS@ICHEP'22]

$\Rightarrow$ Joint experimental & theoretical effort to disentangle SM & NP

BACKUP

Anomaly Cancellation Conditions

Chiral fermion contributions to gauge anomalies have to cancel!

$$SU(3)_C^2 \times U(1)': 2\langle\mathcal{F}_Q\rangle - \langle\mathcal{F}_u\rangle - \langle\mathcal{F}_d\rangle = 0,$$

$$SU(2)_L^2 \times U(1)': 3\langle\mathcal{F}_Q\rangle + \langle\mathcal{F}_L\rangle = 0,$$

$$U(1)_Y^2 \times U(1)': \langle\mathcal{F}_Q\rangle + 3\langle\mathcal{F}_L\rangle - 8\langle\mathcal{F}_u\rangle - 2\langle\mathcal{F}_d\rangle - 6\langle\mathcal{F}_e\rangle = 0,$$

$$\text{gauge-gravity} : 6\langle\mathcal{F}_Q\rangle + 2\langle\mathcal{F}_L\rangle - 3\langle\mathcal{F}_u\rangle - 3\langle\mathcal{F}_d\rangle - \langle\mathcal{F}_e\rangle - \langle\mathcal{F}_\nu\rangle = 0,$$

$$U(1)_Y \times U(1)'^2 : \langle\mathcal{F}_Q^2\rangle - \langle\mathcal{F}_L^2\rangle - 2\langle\mathcal{F}_u^2\rangle + \langle\mathcal{F}_d^2\rangle + \langle\mathcal{F}_e^2\rangle = 0,$$

$$U(1)'^3 : 6\langle\mathcal{F}_Q^3\rangle + 2\langle\mathcal{F}_L^3\rangle - 3\langle\mathcal{F}_u^3\rangle - 3\langle\mathcal{F}_d^3\rangle - \langle\mathcal{F}_e^3\rangle - \langle\mathcal{F}_\nu^3\rangle = 0$$

$$\text{where } \langle\mathcal{F}_\psi^{(n)}\rangle = \sum_{i=1}^3 F_{\psi_i}^{(n)}$$

Additionally to avoid kinetic mixing at 1-loop:

$$\langle\mathcal{F}_Q\rangle - \langle\mathcal{F}_L\rangle + 2\langle\mathcal{F}_u\rangle - \langle\mathcal{F}_d\rangle - \langle\mathcal{F}_e\rangle = 0$$

Constraints on lepton charges

Lepton couplings & charges are constrained by

- ▶ BR's of rare (semi-)leptonic charm decays [2011.09478]
 $\mathcal{B}(D^0 \rightarrow \mu^+ \mu^-) < 2.9 \cdot 10^{-9}$ @90% c.l. (LHCb talk ICHEP'22)
- ▶ Drell-Yan searches [2003.12421]
 $\Rightarrow$ Combination with $a_{\pi^+ \pi^-}^d$ yields as most stringent bounds:

$$\sqrt{F_{L_1}^2 + F_{e_1}^2} \lesssim 2.3 |F_{d_1}| \quad \text{Drell-Yan } e$$

$$|F_{L_2} - F_{e_2}|, \sqrt{F_{L_2}^2 + F_{e_2}^2} \lesssim 0.8 |F_{d_1}| \quad \text{(semi-)muonic } D^0 \text{ decays}$$

$$\sqrt{F_{L_3}^2 + F_{e_3}^2} \lesssim 4.7 |F_{d_1}| \quad \text{Drell-Yan } \tau$$

Kinetic Mixing

Z' - γ kinetic mixing via $\mathcal{L}_\eta = -\frac{\eta}{2} F^{\mu\nu} Z'_{\mu\nu}$

► From ρ -parameter: $\boxed{\varepsilon(M_Z) \lesssim 3 \cdot 10^{-1}}$ where $\varepsilon = \frac{-\eta}{\sqrt{1-\eta^2}}$

► From induced Z' -lepton-couplings via $\mathcal{L}_\varepsilon = -\varepsilon e J^\mu Z'_\mu$:

$$\boxed{\varepsilon(M_{Z'}) \lesssim 10^{-3}} \text{ and } \boxed{F_{L_{1,2}, e_{1,2}} \lesssim 10^{-3} \cdot F_{d_1}}$$

$$\Rightarrow \text{BM III: } |G| \ll |F| \quad \& \quad \text{BM IV: } |F_e| \ll |F_{u,d}|$$

► *Problem:* $\varepsilon(M_Z)$ is in general RG induced even if $\varepsilon(M_{Z'}) = 0$

$$\Rightarrow \text{BM IV: } \varepsilon \text{ is natural } \frac{d\varepsilon}{d \ln \mu} \propto \varepsilon \text{ and controlled by small } \varepsilon(M_{Z'})$$

$$\Rightarrow \text{BM III: RGEs induce } |\varepsilon(M_Z) - \varepsilon(M_{Z'})| \gtrsim 10^{-2}$$

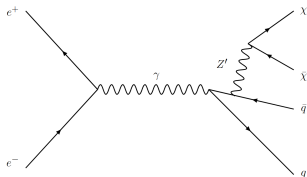
$$\Rightarrow \text{Require moderate tuning of } \varepsilon(M_{Z'}) \text{ order } 10\%$$

$\Rightarrow$ kinetic mixing constraints can be avoided

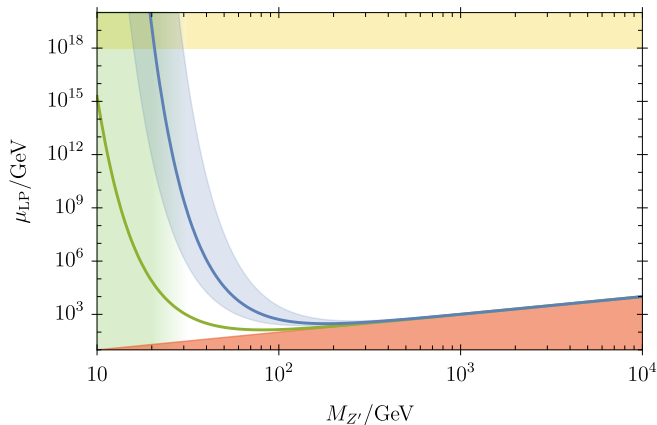
Z' Branching ratios

Model	jets	b	c	e	μ	τ	$\nu_{e,\mu,\tau}$
BM III	38%	0%	37%	0%	0%	12%	13%
BM IV $ _{M_{Z'}=10 \text{ GeV}}$	59%	22%	18%	0%	0%	0%	0%
BM IV $ _{M_{Z'}=20 \text{ GeV}}$	52%	31%	17%	0%	0%	0%	0%

- Decays via kinetic mixing suppressed by $\lesssim \mathcal{O}(10^{-7})$
 $\Rightarrow$ negligible
- Branching ratios can be suppressed by adding $|F_{\nu_i, \chi}| \gtrsim |F_{\psi_i}|$
 $\Rightarrow e^+e^- \rightarrow \text{hadrons} + E/\tau$
 $\Rightarrow |F_{\nu_i, \chi}|$ limited by Landau poles



High Energy Behaviour



$M_{Z'} \lesssim \text{few} \times 10 \text{ GeV}$ to avoid sub-TeV-ish Landau pole