

B Physics Workshop Neckarzimmern 22-24 March, 2017

Direct probes of neutrino mass Tritium β-decay and EC of ¹⁶³Ho

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4 Nobel Prizes for Neutrino Physics







Martin L. Perl Prize share: 1/2

© University of California Regents Frederick Reines Prize share: 1/2

The Nobel Prize in Physics 1995 was awarded "for pioneering experimental contributions to lepton physics" jointly with one half to Martin L. Perl "for the discovery of the tau lepton" and with one half to Frederick Reines "for the detection of the neutrino".





"Herr Auge"

1988







Leon M. Lederman Prize share: 1/3

an Melvin Schwartz Prize share: 1/3

Jack Steinberger Prize share: 1/3

The Nobel Prize in Physics 1988 was awarded jointly to Leon M. Lederman, Melvin Schwartz and Jack Steinberger *"for the neutrino beam method and the demonstration of the doublet structure of the leptons through the discovery of the muon neutrino"*.



4 Nobel Prizes for Neutrino Physics







2002



Raymond Davis Jr. Prize share: 1/4

The Nobel Prize in Physics 2002 was divided, one half jointly to Raymond Davis Jr. and Masatoshi Koshiba "for pioneering contributions to astrophysics, in particular for the detection of cosmic neutrinos" and the other half to Riccardo Giacconi "for pioneering contributions to astrophysics, which have led to the discovery of cosmic X-ray sources".

Masatoshi Koshiba Riccardo Giacconi Prize share: 1/4 Prize share: 1/2



Photo © Takaaki Kajita Takaaki Kajita Prize share: 1/2

Photo: K. MacFarlane Queen's University /SNOLAB Arthur B. McDonald

2015

Prize share: 1/2

The Nobel Prize in Physics 2015 was awarded jointly to Takaaki Kajita and Arthur B. McDonald "for the discovery of neutrino oscillations, which shows that neutrinos have mass"





Homestake





Super-KamiokaNDE



The role of massive neutrinos





What's on the menu today?

- Introduction:
 - Neutrino masses & mixing
 - Methods to measure neutrino masses
- Neutrino mass searches using kinematics:
 - Beta decay of Tritium (³H)
 - Electron capture in Holmium (¹⁶³Ho)
- Summary & outlook





Three-flavour neutrino mixing



3 x 3 unitary mixing matrix analogous to CKM:

"Pontecorvo Maki Nakagawa Sakata" (PMNS)



- 3 mixing angles: θ_{12} , θ_{23} , θ_{13} ,
- 1 CP-violating phase: δ
- 2 independent Δm^2 scales:

$$\Delta m_{13}^2 = \Delta m_{12}^2 + \Delta m_{23}^2$$

$$U = \begin{pmatrix} 1 & 0 & 0 \\ 0 & \cos\theta_{23} & \sin\theta_{23} \\ 0 & -\sin\theta_{23} & \cos\theta_{23} \end{pmatrix} \cdot \begin{pmatrix} \cos\theta_{13} & 0 & \sin\theta_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -\sin\theta_{13}e^{-i\delta} & 0 & \cos\theta_{13} \end{pmatrix} \cdot \begin{pmatrix} \cos\theta_{12} & \sin\theta_{12} & 0 \\ -\sin\theta_{12} & \cos\theta_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix}$$



→ Structure of leptonic mixing matrix very different from CKM matrix:

	$(0.800 \rightarrow 0.844)$	0.515 ightarrow 0.581	$0.139 \rightarrow 0.155$
$U_{PMNS} =$	$0.229 \rightarrow 0.516$	0.438 ightarrow 0.699	$0.614 \rightarrow 0.790$
	$0.249 \rightarrow 0.528$	$0.462 \rightarrow 0.715$	$0.595 \rightarrow 0.776$

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How do we know all this?

→ collected "world data" from many different experiments

1. & 3. generation	1. & 2. generation	
$\Delta m_{13}^2 = 2.5 \times 10^{-3} \text{ eV}^2$	$\Delta m_{12}^{2} = 7.6 \times 10^{-5} \text{ eV}^{2}$	
θ ₁₃ ≈ 8.5° (small)	θ ₂₃ ≈ 34° (large)	
reactor & long-baseline accelerator exp.	solar & reactor exp.	
MeV, $\overline{\mathbf{v}}_{e}$ GeV, \mathbf{v}_{μ} ($\overline{\mathbf{v}}_{\mu}$)	MeV, v_e (\overline{v}_e)	
	<image/>	
Double Chooz	SNO KamLAND	
	1. & 3. generation $\Delta m_{13}{}^2 = 2.5 \times 10^{-3} \text{ eV}^2$ $\theta_{13} \approx 8.5^\circ \text{ (small)}$ reactor & long-baseline accelerator exp. $MeV, \overline{v}e$ $GeV, v_{\mu} (\overline{v}_{\mu})$ $WeV = 0$	

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Results of v-oscillation experiments

V_T

Large neutrino mixing and tiny neutrino masses $m(v_i) \neq 0$ established

- Which mass ordering (normal, inverted)? ullet
- What is the absolute v mass scale? •

 $\Delta m_{\rm atm}^2$ m_{1}^{2} m_{3}^{2} 0 inverted normal

So far: only **upper** (< 2 eV) and **lower bounds** (>0.01) >0.05 eV) resp.

Well, then, how to determine the neutrino mass scale?

Astrophysicist's answer: Use astrophysical data!

today

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- Estimate v mass from energy-dependent time of flight. Current limit: $m(v_e) < 5.7 eV$.
- ➡ Purely kinematical. Main difficulties: (a) rare! (b) emission model! [Loredo et al., 2005]

24. Feb. 1987

Core-collapse SN 1987a in LMC

Well, then, how to determine the neutrino mass scale?

Cosmologist's answer: Use cosmological data!

[T. Haugbølle, Univ. of Aarhus]

Well, then, how to determine the neutrino mass scale?

Experimentalist's answer: 3 complementary paths!

ΤοοΙ	Cosmology CMB + LSS +	Neutrinoless double β-decay	β-decay endpoint and EC
Observable	$\sum m_{\nu} = \sum_{i=1}^{3} m_i$	$\langle m_{\beta\beta} \rangle = \left \sum_{j=1}^{3} U_{ej} ^2 m_j e^{i\alpha_j} \right $	$m_{\beta}^2 = \sum_{i=1}^3 U_{ei}^2 m_i^2$
Present upper limit	0.15 – 1 eV	0.2 – 0.4 eV	2 eV
Potential	20 – 50 meV	20 – 50 meV	200 meV
Model dependence	Multi-parameter cosmological model	 Majorana vs. Dirac phase cancellations possible nucl. matrix elements 	Direct , only kinematics; no cancellations in incoherent sum

Kinematics of weak decays: "direct" v mass search

β decay and neutrino mass

postulation of new particle: neutral, spin 1/2, weak interaction

1930: Wolfgang Pauli's "desperate remedy" to solve the problem of apparent violation of energy & momentum conservation in ß decay

4 December 1930 Gloriastr. Zürich

Physical Institute of the Federal Institute of Technology (ETH) Zürich

Dear radioactive ladies and gentlemen,

As the bearer of these lines, to whom I ask you to listen graciously, will explain more exactly, considering the 'false' statistics of N-14 and Li-6 nuclei, as well as the continuous β -spectrum, I have hit upon a desperate remedy to save the "exchange theorem" of statistics and the energy theorem. Namely [there is] the possibility that there could exist in the nuclei electrically neutral particles that I wish to call neutrons, ** which have spin 1/2 and obey the exclusion principle, and additionally differ from light quanta in that they do not travel with the velocity of light: The mass of the neutron must be of the same order of magnitude as the electron mass and, in any case, not larger than 0.01 proton mass. The continuous β -spectrum would then become understandable by the assumption that in β decay a neutron is emitted together with the electron, in such a way that the sum of the energies of neutron and electron is constant.

β decay and neutrino mass

E. Fermi, 1934

Versuch einer Theorie der β -Strahlen. I¹). Von E. Fermi in Rom.

Mit 3 Abbildungen. (Eingegangen am 16. Januar 1934.)

Eine quantitative Theorie des β -Zerfalls wird vorgeschlagen, in welcher man die Existenz des Neutrinos annimmt, und die Emission der Elektronen und Neutrinos aus einem Kern beim β -Zerfall mit einer ähnlichen Methode behandelt, wie die Emission eines Lichtquants aus einem angeregten Atom in der Strahlungstheorie. Formeln für die Lebensdauer und für die Form des emittierten kontinuierlichen β -Strahlenspektrums werden abgeleitet und mit der Erfahrung verglichen.

1. Grundannahmen der Theorie.

Bei dem Versuch, eine Theorie der Kernelektronen sowie der β -Emission aufzubauen, begegnet man bekanntlich zwei Schwierigkeiten. Die erste ist durch das kontinuierliche β -Strahlenspektrum bedingt. Falls der Erhaltungssatz der Energie gültig bleiben soll, muß man annehmen, daß ein Bruchteil der beim β -Zerfall frei werdenden Energie unseren bisherigen Beobachtungsmöglichkeiten entgeht. Nach dem Vorschlag von W. Pauli kann man z. B. annehmen, daß beim β -Zerfall nicht nur ein Elektron, sondern auch ein neues Teilchen, das sogenannte "Neutrino" (Masse von der Größenordnung oder kleiner als die Elektronenmasse; keine elektrische

1934: Fermi 4-point int.

1938 (Yukawa, Klein et al.): boson-mediated weak int.

- Model independent neutrino mass measurement: based solely on kinematic variables (E₀, E_e, p_e, m_e) and energy conservation
- Elusive neutrino escapes undetected, only electron energy is measured

Direct kinematic determination of m(v_e)

$$\frac{\mathrm{d}\Gamma}{\mathrm{d}E} = C F(Z, E) p \left(E + m_{\mathrm{e}}\right) \left(E_0 - E\right) \sum_i |U_{\mathrm{e}i}|^2 \sqrt{(E_0 - E)^2 - m^2(\nu_i)}$$

Key requirements:

- high-activity source
- low-endpoint β emitter (³H) or EC isotope (¹⁶³Ho)
- excellent energy resolution (MAC-E filter or calorimeter)

kinematic measurement can probe for **heavier neutrino states** → eV-scale and keV-scale sterile v spectral distortion measures "effective" mass square:

 $m^2(\nu_{\rm e}) := \sum_i |U_{\rm ei}|^2 m_i^2$

Moore's Law of direct neutrino mass searches

Magnetic vs. electrostatic spectrometers

Status of the KATRIN source

Gaseous molecular tritium source of

- high activity (~170 GBq)
- high isotopic purity ($\epsilon_T > 95\%$)
- high stability (0.1%)

Extensive control of systematics Gas column: 40% no-loss electrons

- Mechanical and cryoinfrastructure complete
- ✓ Test of 800 sensors & valves

Source: Beam-tube cooling performance

Novel 2-phase neon beam tube cooling system successfully tested (Nov. 2016)

Stability of beam tube temperature surpassing specifications

Axial temperature profile along 10 m beam tube

Status of transport & pumping sections

- fully adiabatic, **lossless electron transport** in 5.6 T magnetic field
- reduction of T₂ flow rate towards spectrometers by factor >10¹⁴ by magnetic chicane with differential and cryo-pumping
- blocking of ion flux by electrostatic barrier

passive cryo-pumping with Ar frost at 3-4 K status: cool-down to 3.5 K successful, s.c. magnets operational; *next:* Ar operation tests

active differential pumping with TMPs

next: instrumentation (FT-ICR)

status: magnets, beam tubes installed & tested,

KATRIN spectrometer status

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LFCS low-field fine-tuning

EMCS earth field compensation

Ø = 12.7 m

2011: fully commissioned large Helmholtz coil system

KATRIN spectrometer: filter characteristics

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Background mitigation in MAC-E filters

- 8 sources of background investigated and understood
- 7 out of 8 avoided or actively eliminated by
 - fine-shaping of special electrodes
 - symmetric magnetic fields
 - LN₂-cooled baffles (cold traps)
 - wire electrode grids

• 1 out of 8 remaining:

caused by ²¹⁰Pb on spectrometer walls (neutral H* atoms ionised by black-body radiation in spectrometer)

KATRIN: v-mass sensitivity

- Relative shape measurement of integrated β spectrum
- 4 fit parameters: m_{v}^{2} , E_{0} , A_{S} , R_{Bg}
- After **3 yrs** of data (~5 cal. yrs): balance of statistics and systematics

KATRIN milestone: "First Light"

Neutrinos auf der Waage

Am 14. Oktober durchflogen erstmals Elektronen das Experiment KATRIN am Karlsruher Institut für Technologie

Neutrinos durchdringen uns jede Sekunde milliardenfach, ohne dass wir das Geringste davon bemerken würden. Lange Zeit galten die mysteriösen Teilchen daher als masselos. Seit dem Nachweis von Neutrino-Oszillationen, der im vergangenen Jahr mit dem Physik Nobelpreis ausgezeichner varder v aber klar, dass Neutrinos für ringe Masse besitzen phö wie groß sie genau jst

nicht. Diese Frage

technical inauguration

of KATRIN

with photoelectrons

October 14, 2016

KATRIN: Next steps

Installation of precision electron source ("e-gun")

Major importance for systematics:

- precision determination of overall transmission function (requirements: ~few degree angular selectivity and ~100 meV energy resolution)
- determination of energy loss function
- in-line column density monitoring

[K. Valerius et al., JINST 6 (2011) P01002; J. Behrens et al., arXiv:1703.05272]

KATRIN: Next steps

Connection & testing of tritium gas circuits integration with existing infrastructure from TLK fusion research **Transport section** CMS Tritium source (WGTS) <10⁻¹⁴ mbarℓ/s 1,8 mbarℓ/s <10⁻² mbarℓ/s <10⁻⁷ mbarℓ/s WGTS DPS1-R DPS1-F CPS CMS DPS2 tube 6x (0) Inner Loop 8x(0) (0) (0)0 (0) @ WGTS (0)Buffer p- & Tvessel Buffer controlled vessel buffer Laser vessel Raman Argon Helium Inner Loop @ ISS (stabilised tritium Outer Loop Outer Loop Outer Loop injection) @ DPS @ CPS @ CMS **Central Tritium Tritium recovery & isotope separation Retention System**

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KATRIN: Next steps

Intermediate summary

 Discovery of non-zero neutrino masses → incompleteness of the Standard Model

- Mass scale & pattern remain to be uncovered!
- Increasingly stringent bounds from cosmology → ∧CDM paradigm

 Direct neutrino mass determination using kinematics (³H β decay)

- KATRIN experiment in final commissioning phase
 - → start data-taking in spring 2018

How to further improve v-mass sensitivity?

Problems:

New developments

Several avenues towards improvement:

1st avenue: exploit differential β spectrum

Idea: Upgrade to MAC-E-TOF spectrometer

Spectrometer as 24 m long "delay line"
→ very sensitive to small differences in surplus energy

TOF spectrum records full β spectrum \Rightarrow save meas. time by using only few voltage settings of MAC-E filter

Coincidence requirement → add. background suppression

Technical realization?

(a) pre-spectrometer as gated filter(b) radio frequency tagger

2nd avenue: alternative spectroscopic technique

Idea: Cyclotron Radiation Emission Spectroscopy (CRES)

Non-destructive measurement of electron energy via cyclotron frequency:

$$\omega(\gamma) = rac{\omega_{
m c}}{\gamma} = rac{eB}{E_{
m kin} + m_{
m e}}$$

uniform B-field, magnetic trap low-pressure gas cell

antenna array

UW Seattle, MIT, UCSB, Pacific NW, CfA, Yale, Livermore, KIT, U Mainz

Phase I system

→ Proof of principle of CRES technique

Project 8 – phase I results

First observation of cyclotron radiation from single electrons

Calorimeters to measure ¹⁶³Dy* atomic de-excitation

MMC: metallic magnetic calorimeters with paramagnetic sensor Au:Er

 δT in absorber from EC-decay \Rightarrow change in magnetization M of sensor

signal:
$$\delta \Phi_s \sim \frac{\partial M}{\partial T} \cdot \Delta T \sim \frac{\partial M}{\partial T} \cdot \frac{1}{C_{tot}} \cdot \delta E$$

thermal micro-calorimeters with transition edge sensor (TES)

δT in absorber from EC-decay
 ⇒ change in temperature T and
 resistance R of thermistor
 signal: current change measured by
 SQUID array

TES technology: NuMECS

 full ~1000 channel array operation expected for 2018

Heidelberg (Univ., MPIK), U Mainz, U Tübingen, TU Dresden, U Frankfurt, HU Berlin, ILL Grenoble, PNPI St Petersburg, U Bratislava, IIT Roorkee, Saha Inst. Kolkata

Technology

- magnetic micro-calorimeter (MMC) arrays with microwave SQUID multiplexing readout
- fast rise time (~130 ns) and excellent linearity & resolution ($\Delta E \sim 5 \text{ eV}$)
- isotope production: ${}^{162}Er(n,\gamma){}^{163}Ho$ offline mass separation

 Phase II: ECHo-1M array of 10⁵ detectors 50 x 2000 pixel x 10 Bq, 2 years: sub-eV sensitivity

Energy E [keV]

Direct v-mass determination: status and outlook

Current achievements

- **KATRIN:** under commissioning, preparing for first tritium runs
- Project 8: successful prototype, CRES proof of principle with ^{83m}Kr

Next goals

- Long-term data-taking in integral spectroscopy mode (0.2 eV)
- Develop CRES towards first tritium measurement (2 eV) and beyond

ECHo, HOLMES, NuMECS:

- Detector development (MMC, TES)
- Test of scalable arrays
- High-purity ¹⁶³Ho production and implantation into absorber

NECKARZIMMERN 2017

- Operate medium-size arrays (~10¹⁰ counts) for 10 eV sens.
- Prepare large arrays (~10¹⁴ counts) for sub-eV range

NEUTRINO 2018

Bonus material: Other fun physics with precision β spectroscopy

KATRIN: v-mass sensitivity ... and more:

Explore physics potential

close to the spectral endpoint E₀:

RH currents **Constraining local CvB** Search for eV-scale sterile v Bonn et al. (2011) e.g. Kaboth & Formaggio (2010), e.g. Formaggio & Barrett (2011) Fässler et al. (2013) Violation of Lorentz symmetry e.g. Blasone et al. (2005) Diaz, Kostelecky & Lehnert (2013) $\sim 1 \text{ eV}^2$ 0.8 capture of 0.6 $m(v_e) = 0 eV$ relic v on 0.4 **B-instable** 0.2 $m(v_e) = 1 e^{1}$ nuclei E - E₀ [eV]

and further away from E₀:

search for keV-mass scale sterile v as WDM candidates

Mertens et al. (2015); see also Barry, Heek & Rodejohann (2014)

non-standard operation, novel detector concepts

Why sterile neutrinos?

Hints of eV-scale sterile neutrinos?

May explain anomalous short-baseline oscillation results:

Hints of keV-scale sterile neutrinos?

Pure (Λ)CDM scenario:

- Missing satellites problem
- Too-big-to-fail problem
- "Core" vs. "cusp" problem

Why sterile neutrinos?

5 Δχ²

10

[°]10[°]

Hints of eV-scale sterile neutrinos?

May explain anomalous oscillation results from

10

[G. Mention et al. (2011), updated in White Paper (2014)]

sin²(2θ_{new})

- Short baseline accelerator experiments
- Gallium experiments
- Reactor experiments

10

^**∠**X²

 10^{2}

10

10

10

, 10⁻³

∆m²_{new} (eV²)

2 dof As

contour

1 dof $\Delta \chi^2$ profile

່10⁻²

Hints of keV-scale sterile neutrinos?

Well motivated as natural extension of Standard Model (vMSM)

Both scales accessible in tritium β decay [e.g., Canetti, Drewes, Shaposhnikov (2013)]

> In agreement with cosmological observations from small to large scales [e.g., Shi & Fuller (1999)]

Recent indirect hints from X-ray astronomy?

Imprint of sterile neutrinos on β spectrum

Shape modification below E_0 by active $(m_a)^2$ and sterile $(m_s)^2$ neutrinos:

-3

electron energy $E-E_{o}$ (eV)

-4

-2

additional kink in β spectrum at E = E₀ – m_s

10

-1

0

16

20

E [keV]

0

-5

Search for eV-scale sterile v with direct mass

Search for eV-scale sterile v with direct mass experiments

Combined sensitivity of direct neutrino mass exp. and SBL oscillation searches

Search for keV-scale sterile v with direct mass experiments

- First measurements with KATRIN "baseline" set-up at reduced source strength
- Prototyping and sensitivity studies for upgraded detector system under way
- Sensitivity of holmium experiments restricted by low Q-value (2.8 keV)

Search for keV-scale sterile v with KATRIN

The challenge:

- High count rates at ~few keV below endpoint
- Tiny sterile admixture $sin^2(\theta_s)$ expected
- Best sensitivity for differential measurement (energy or ToF)
- Development of new techniques necessary!

Differential detection option: novel detector required

TRISTAN* design study:

- 10⁸ cps (> 10 000 pixels)
- FWHM 300 eV @ 20 keV
- > 20 cm diameter

[Mertens et al. (2015)]

*) TRitium beta decay Investigation on Sterile To Active Neutrino mixing

[Steinbrink et al. (2013), Robertson et al. (in prep.)]

Effect of RH current contributions

$$\frac{\mathrm{d}\Gamma}{\mathrm{d}E\,\mathrm{d}t} \propto E_{\nu}\sqrt{E_{\nu}^{2}-m_{\nu}^{2}}\left(1+b'\frac{m_{\nu}}{E_{\nu}}\right) \qquad \text{[J. Bonn et al., Phys. Lett. B 703 (2011) 310]}$$
Fierz-like parameter b'
enters differential rate
$$b' \approx -2\frac{\Re(L_{V}R_{V}^{*}+L_{V}R_{S}^{*})|\mathcal{M}_{F}|^{2}+\Re(L_{A}R_{A}^{*}+L_{A}R_{T}^{*})|\mathcal{M}_{GT}|^{2}}{|L_{V}|^{2}|\mathcal{M}_{F}|^{2}+|L_{A}|^{2}|\mathcal{M}_{GT}|^{2}}$$

Imprint on integrated spectrum:

- Only small sensitivity on b' if endpoint E₀ left free in fit
 - → good for determination of $m^2(v_e)$
- Improvement of present bounds on b' with KATRIN for small m(ve) if
 - external E₀ value with accuracy
 50 meV as input*
 - absolute energy scale in KATRIN
 U_{spec} U_{source} known to same
 accuracy of < 50 meV

*) 70 meV accuracy:

E. G. Myers et al., PRL 114 (2015) 013003

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Probing Lorentz invariance in β decay

[Diaz, Kostelecky, & Lehnert 2013]

Standard Model Extension (SME) framework:

Neutrinos satisfy Dirac-like equation

 $(i\mathbf{\Gamma}^{\alpha}\partial_{\alpha}-\mathbf{M})\,\psi=0$

with Γ , **M** including momentumdependent coefficients

electron kinetic energy

- Modified energy dependence of decay rate
- Spectral shape dependent on sidereal time and experiment orientation
- Effective dim-3 coefficient: osc. shift of endpoint $T_{0,eff}$ with $\omega_{sidereal}$
- Effective dim-2 coefficient: osc. of m² parameter (can mimic tachyonic v)