

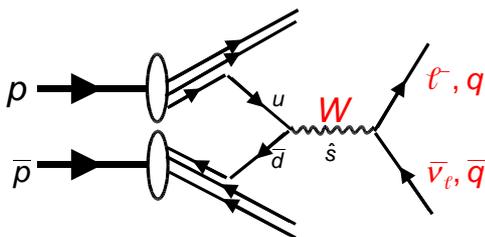
Experimental tests of the Standard Model

- Discovery of the W and Z bosons
- Precision tests of the Z sector
- Precision tests of the W sector
- Electro-weak unification at HERA
- Radiative corrections and prediction of the top and Higgs mass
- Top discovery at the Tevatron
- Higgs searches at the LHC

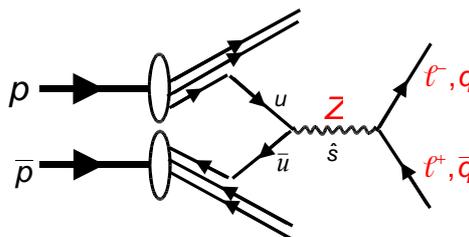
1. Discovery of the W and Z boson

1983 at CERN SppS accelerator,
 $\sqrt{s} \approx 540$ GeV, UA-1/2 experiments

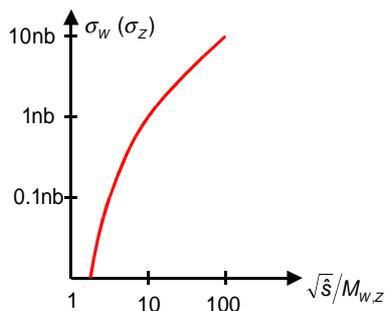
1.1 Boson production in $p\bar{p}$ interactions



$$p\bar{p} \rightarrow W \rightarrow \ell \bar{\nu}_\ell + X$$



$$p\bar{p} \rightarrow Z \rightarrow \ell \bar{\ell} + X$$



Similar to Drell-Yan: (photon instead of W)

$$\hat{s} = x_q x_{\bar{q}} s \quad \text{mit} \quad \langle x_q \rangle \approx 0.12$$

$$\hat{s} = \langle x_q \rangle^2 s \approx 0.014 s = (65 \text{ GeV})^2$$

→ Cross section is small !

1.2 UA-1 Detector

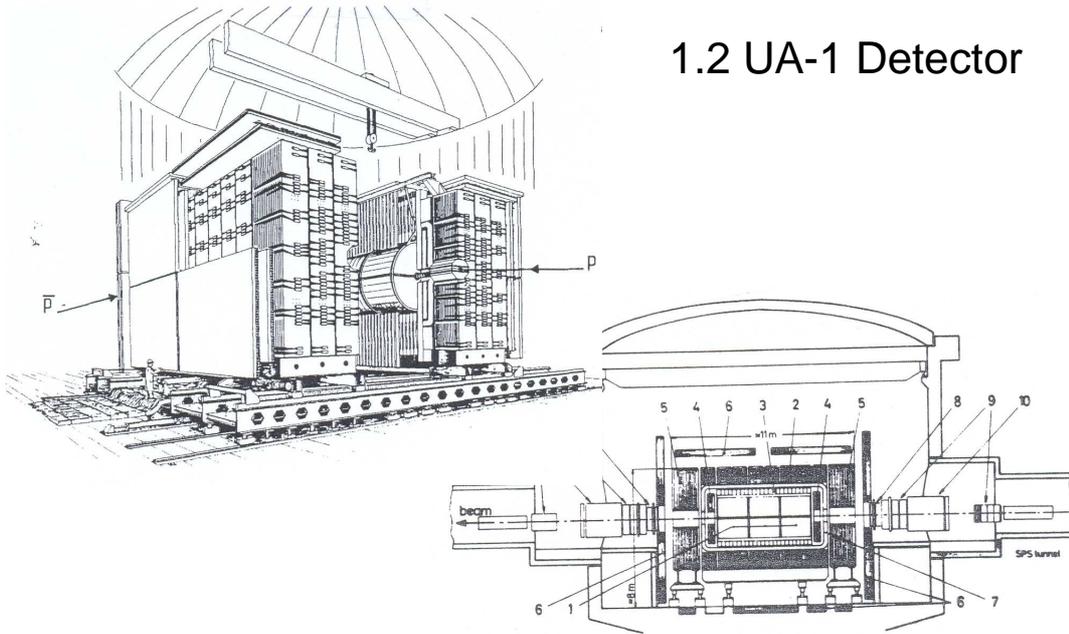
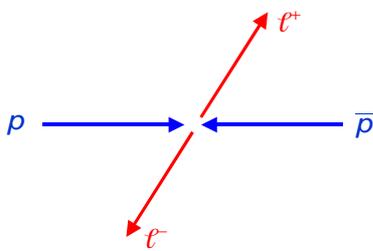


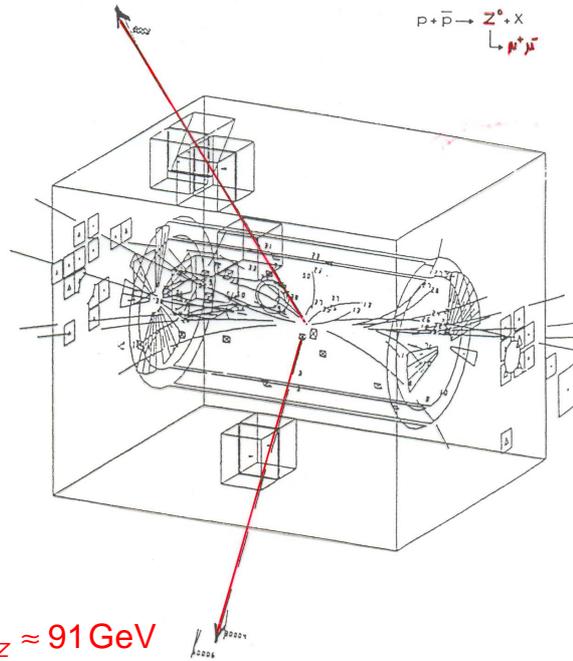
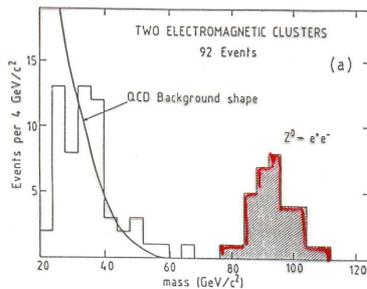
Fig.8.16: Seitenansicht des UA1-Detektors zum Nachweis von Proton-Antiproton-Wechselwirkungen bei 540 GeV Schwerpunktsenergie: 1. Zentraldetektor, 2. und 5. Hadron-Kalorimeter, 3. und 4. Elektron-Photon-Schauerzähler, 6. Myon-Detektor, 7. Spule für Dipolfeld, 8. und 9. Kleinwinkeldetektor mit Kammern und Kalorimetern, 10. Kompensator-Magnete [UA1].

1.3 Event signature: $p\bar{p} \rightarrow Z \rightarrow f\bar{f} + X$



High-energy lepton pair:

$$m_{e\bar{e}}^2 = (p_{e^+} + p_{e^-})^2 = M_Z^2$$



$M_Z \approx 91 \text{ GeV}$

1.4 Event signature: $p\bar{p} \rightarrow W \rightarrow \ell\bar{\nu}_\ell + X \quad W^- \rightarrow e\bar{\nu}$

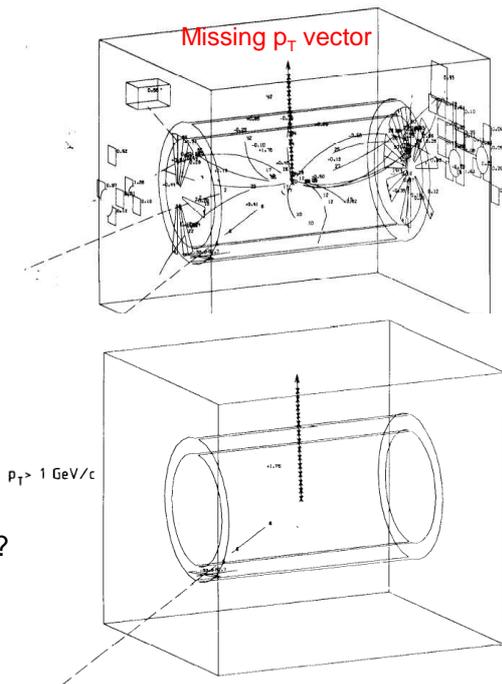
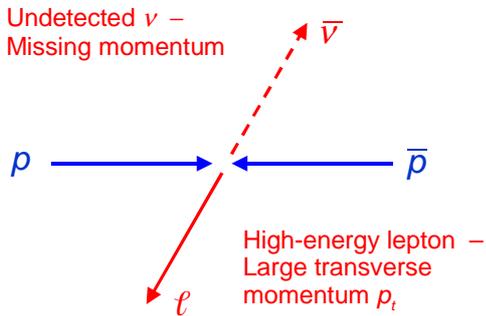
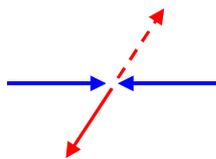


Fig. 16b. The same as picture (a), except that now only particles with $p_T > 1$ GeV/c and calorimeters with $E_T > 1$ GeV are shown.

How can the W mass be reconstructed ?

W mass measurement

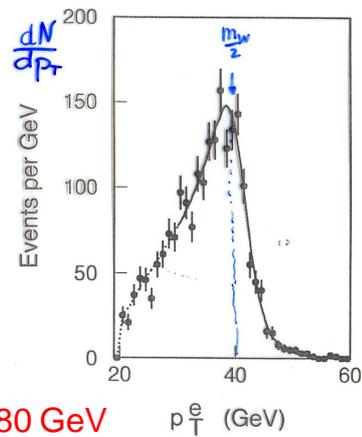
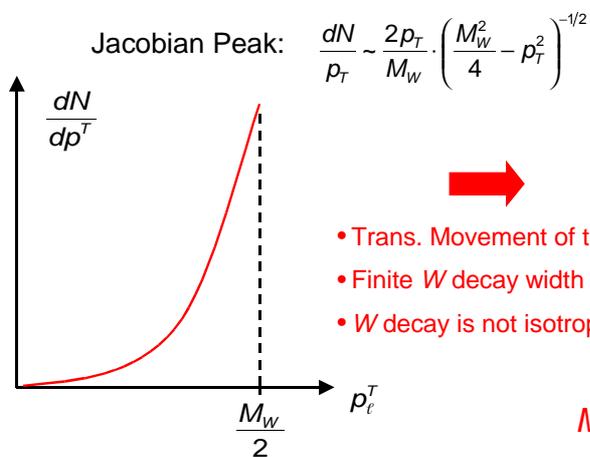


In the W rest frame:

- $|\vec{p}_\ell| = |\vec{p}_\nu| = \frac{M_W}{2}$
- $|\vec{p}_\ell^T| \leq \frac{M_W}{2}$

In the lab system:

- W system boosted only along z axis
- p_T distribution is conserved





The Nobel Prize in Physics 1984



Carlo Rubbia

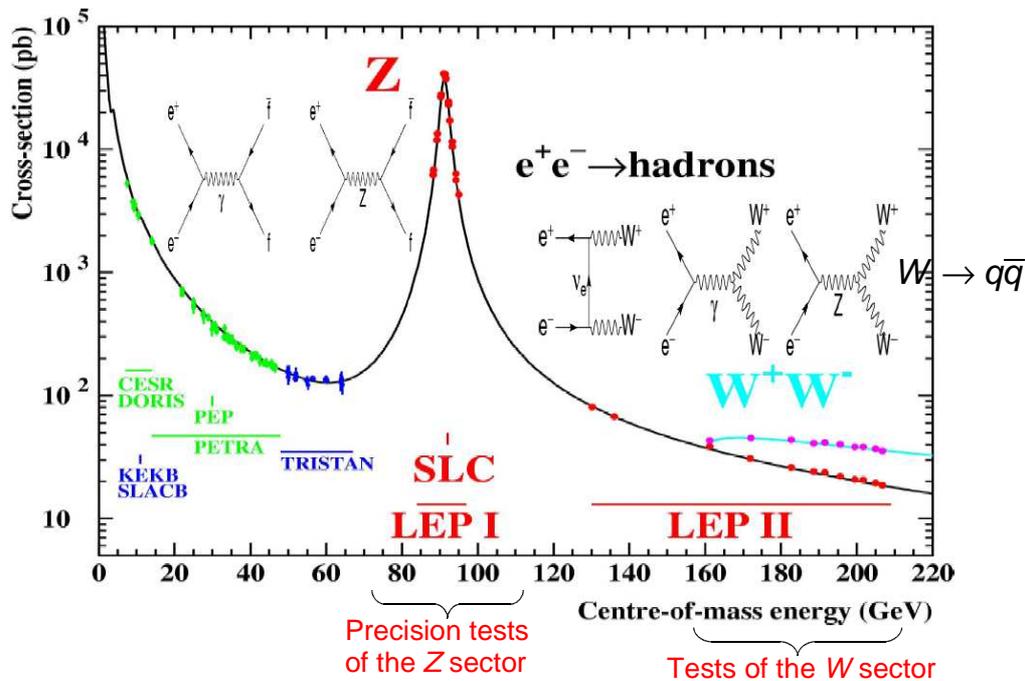
Simon van der Meer

"for their decisive contributions to the large project, which led to the discovery of the field particles W and Z, communicators of weak interaction"

S. van der Meer

One of the achievements to allow high-intensity $p\bar{p}$ collisions, is stochastic cooling of the \bar{p} beams before inserting them into SPS.

1.5 Production of Z and W bosons in e^+e^- annihilation



2. Precision tests of the Z sector

(LEP and SLC)

2.1 Cross section for $e^+ e^- \rightarrow \gamma/Z \rightarrow f\bar{f}$

$\sim 4.5M$ Z decays / experiment

$$|M|^2 = \left| \begin{array}{c} \text{diagram with } \gamma \text{ propagator} \\ + \\ \text{diagram with } Z \text{ propagator} \end{array} \right|^2$$

for $e^+ e^- \rightarrow \mu^+ \mu^-$

$$M_\gamma = -e^2 (\bar{\mu} \gamma_\mu \mu) \frac{1}{q^2} (\bar{e} \gamma^\mu e)$$

$$M_Z = -\frac{g^2}{\cos^2 \theta_W} \left[\bar{\mu} \gamma^\nu \frac{1}{2} (g_V^\mu - g_A^\mu \gamma^5) \mu \right] \underbrace{\frac{g_{V\rho} - q_\nu q_\rho / M_Z^2}{(q^2 - M_Z^2) + iM_Z \Gamma_Z}}_{Z \text{ propagator considering a finite } Z \text{ width}} \left[\bar{e} \gamma^\rho \frac{1}{2} (g_V^e - g_A^e \gamma^5) e \right]$$

Z propagator considering a finite Z width

Differential cross section for $e^+ e^- \rightarrow \mu^+ \mu^-$

$$\frac{d\sigma}{d\cos\theta} = \frac{\pi\alpha^2}{2s} \left[\underbrace{F_\gamma(\cos\theta)}_\gamma + \underbrace{F_{\gamma Z}(\cos\theta)}_{\gamma/Z \text{ interference}} \frac{s(s-M_Z^2)}{(s-M_Z^2)^2 + M_Z^2 \Gamma_Z^2} + \underbrace{F_Z(\cos\theta)}_Z \frac{s^2}{(s-M_Z^2)^2 + M_Z^2 \Gamma_Z^2} \right]$$

Vanishes at $\sqrt{s} \approx M_Z$

with

$$F_\gamma(\cos\theta) = Q_e^2 Q_\mu^2 (1 + \cos^2 \theta) = (1 + \cos^2 \theta)$$

$$F_{\gamma Z}(\cos\theta) = \frac{Q_e Q_\mu}{4 \sin^2 \theta_W \cos^2 \theta_W} \left[2g_V^e g_V^\mu (1 + \cos^2 \theta) + 4g_A^e g_A^\mu \cos \theta \right]$$

$$F_Z(\cos\theta) = \frac{1}{16 \sin^4 \theta_W \cos^4 \theta_W} \left[(g_V^{e^2} + g_A^{e^2})(g_V^{\mu^2} + g_A^{\mu^2})(1 + \cos^2 \theta) + 8g_V^e g_A^e g_V^\mu g_A^\mu \cos \theta \right]$$

At the Z-pole $\sqrt{s} \approx M_Z \rightarrow$ Z contribution is dominant
 \rightarrow interference vanishes

$$\sigma_{tot} \approx \sigma_Z = \frac{4\pi}{3s} \frac{\alpha^2}{16 \sin^4 \theta_w \cos^4 \theta_w} \cdot [(g_V^e)^2 + (g_A^e)^2] [(g_V^\mu)^2 + (g_A^\mu)^2] \cdot \frac{s^2}{(s - M_Z^2)^2 + (M_Z \Gamma_Z)^2}$$

$$\sigma(s) = 12\pi \frac{\Gamma_e \Gamma_\mu}{M_Z^2} \cdot \frac{s}{(s - M_Z^2)^2 + M_Z^2 \Gamma_Z^2}$$

$$\sigma_Z(\sqrt{s} = M_Z) = \frac{12\pi}{M_Z^2} \frac{\Gamma_e \Gamma_\mu}{\Gamma_Z^2}$$

With partial and total widths:

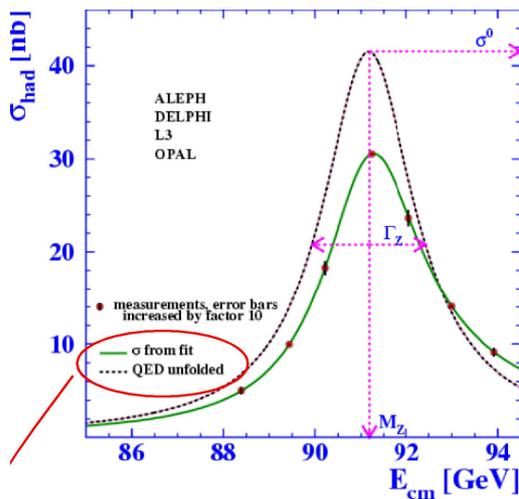
$$\Gamma_i = \frac{\alpha M_Z}{12 \sin^2 \theta_w \cos^2 \theta_w} \cdot [(g_V^f)^2 + (g_A^f)^2]$$

$$\Gamma_Z = \sum_i \Gamma_i$$

Cross sections and widths can be calculated within the Standard Model if all parameters are known



!2 Measurement of the Z lineshape



Z resonance curve:

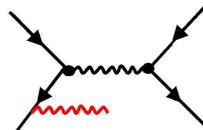
$$\sigma(s) = 12\pi \frac{\Gamma_e \Gamma_\mu}{M_Z^2} \cdot \frac{s}{(s - M_Z^2)^2 + M_Z^2 \Gamma_Z^2}$$

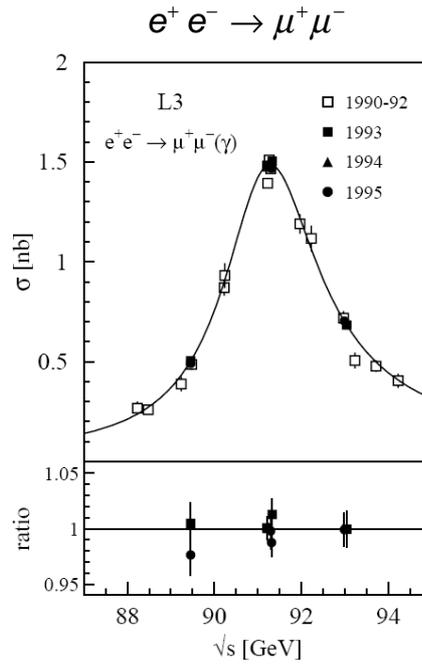
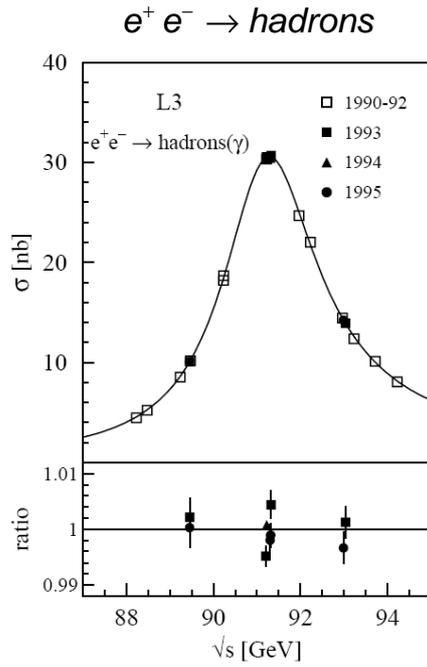
Peak: $\sigma_0 = \frac{12\pi}{M_Z^2} \frac{\Gamma_e \Gamma_\mu}{\Gamma_Z^2}$

- Resonance position $\rightarrow M_Z$
- Height $\rightarrow \Gamma_e \Gamma_\mu$
- Width $\rightarrow \Gamma_Z$

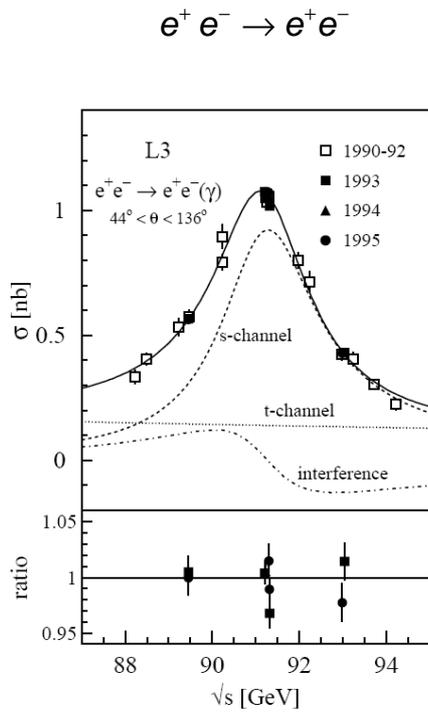
Initial state Bremsstrahlung corrections

$$\sigma_{ff(\gamma)} = \int_{4m_f^2/s}^1 G(z) \sigma_{ff}^0(zs) dz \quad z = 1 - \frac{2E_\gamma}{\sqrt{s}}$$

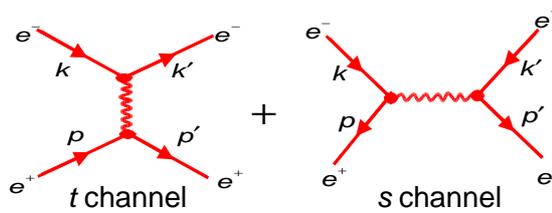




Resonance looks the same, independent of final state: Propagator is the same



t channel contribution \rightarrow forward peak



Z line shape parameters (LEP average)

$$M_Z = 91.1876 \pm 0.0021 \text{ GeV} = \pm 23 \text{ ppm} (*)$$

Γ_Z	$= 2.4952 \pm 0.0023$	GeV	} $\pm 0.09\%$	3 leptons are treated independently
Γ_{had}	$= 1.7458 \pm 0.0027$	GeV		
Γ_e	$= 0.08392 \pm 0.00012$	GeV		
Γ_μ	$= 0.08399 \pm 0.00018$	GeV		
Γ_τ	$= 0.08408 \pm 0.00022$	GeV		
<hr/>				
Γ_Z	$= 2.4952 \pm 0.0023$	GeV	} Assuming lepton universality: $\Gamma_e = \Gamma_\mu = \Gamma_\tau$	↑↓ test of lepton universality
Γ_{had}	$= 1.7444 \pm 0.0022$	GeV		
Γ_e	$= 0.083985 \pm 0.000086$	GeV		

*) error of the LEP energy determination: $\pm 1.7 \text{ MeV}$ (19 ppm)

<http://lepewwg.web.cern.ch/> (Summer 2005)

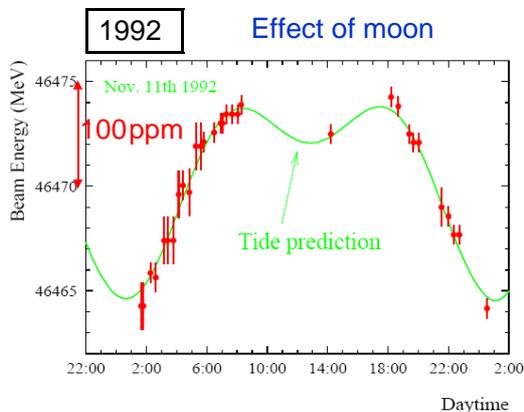
LEP energy calibration: Hunting for ppm effects

Changes of the circumference of the LEP ring changes the energy of the electrons:

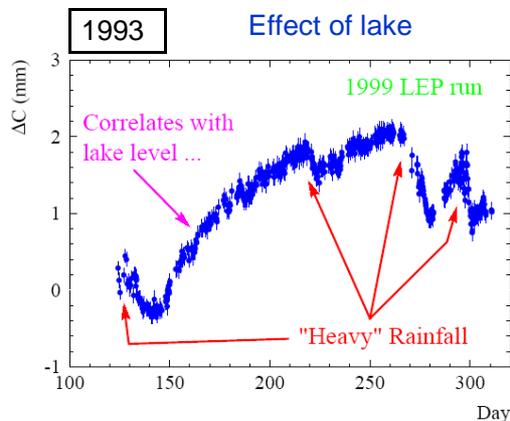
- tide effects
- water level in lake Geneva



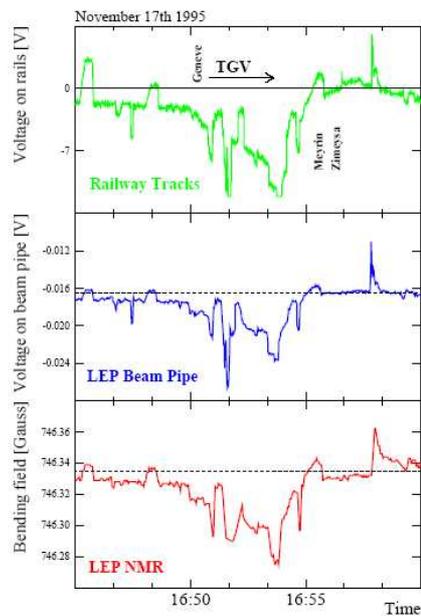
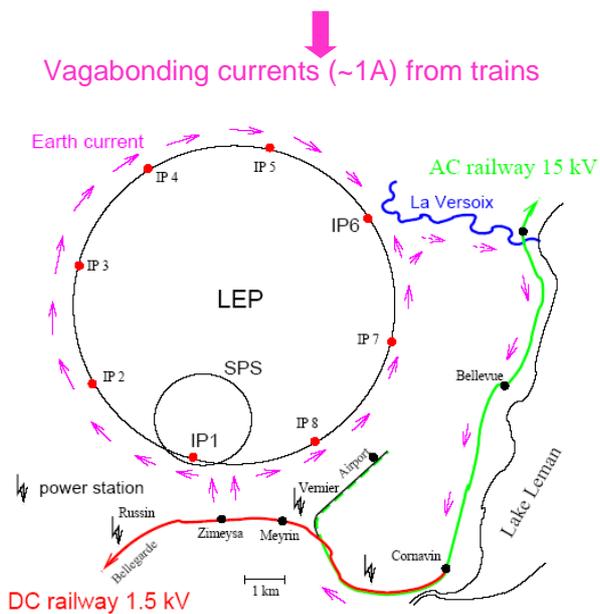
Changes of LEP circumference $\Delta C = 1 \dots 2 \text{ mm} / 27 \text{ km}$ ($4 \dots 8 \times 10^{-8}$)



The total strain is 4×10^{-8} ($\Delta C = 1 \text{ mm}$)



Effect of the French "Train a Grande Vitesse" (TGV)



In conclusion: Measurements at the ppm level are difficult to perform. Many effects must be considered!

2.3 Number of light neutrino generations

In the Standard Model:

$$\Gamma_Z = \Gamma_{\text{had}} + 3 \cdot \Gamma_\ell + \underbrace{N_\nu \cdot \Gamma_\nu}_{\text{invisible} : \Gamma_{\text{inv}}} \rightarrow \begin{cases} e^+ e^- \rightarrow Z \rightarrow \nu_e \bar{\nu}_e \\ e^+ e^- \rightarrow Z \rightarrow \nu_\mu \bar{\nu}_\mu \\ e^+ e^- \rightarrow Z \rightarrow \nu_\tau \bar{\nu}_\tau \end{cases}$$

$$\Gamma_{\text{inv}} = 0.4990 \pm 0.0015 \text{ GeV}$$

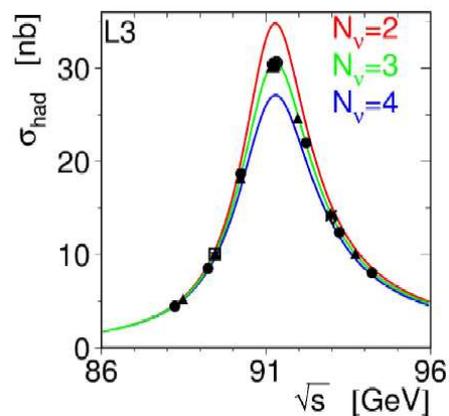
To determine the number of light neutrino generations:

$$N_\nu = \left(\frac{\Gamma_{\text{inv}}}{\Gamma_\ell} \right)_{\text{exp}} \cdot \left(\frac{\Gamma_\ell}{\Gamma_\nu} \right)_{\text{SM}}$$

$$5.9431 \pm 0.0163 = 1.991 \pm 0.001 \text{ (small theo. uncertainties from } m_{\text{top}} M_H)$$

$$N_\nu = 2.9840 \pm 0.0082$$

No room for new physics: $Z \rightarrow \text{new}$



4 Forward-backward asymmetry and fermion couplings to Z

Differential cross section

$$\frac{d\sigma}{d\cos\theta} = \frac{\pi\alpha^2}{2s} \left[\underbrace{F_\gamma(\cos\theta)}_{\gamma} + \underbrace{F_{\gamma Z}(\cos\theta)}_{\gamma/Z \text{ interference}} + \underbrace{F_Z(\cos\theta)}_Z \right]$$

Vanishes at $\sqrt{s} \approx M_Z$

$$F_\gamma(\cos\theta) = Q_e^2 Q_\mu^2 (1 + \cos^2\theta) = (1 + \cos^2\theta)$$

$$F_{\gamma Z}(\cos\theta) = \frac{Q_e Q_\mu}{4 \sin^2\theta_W \cos^2\theta_W} [2g_V^e g_V^\mu (1 + \cos^2\theta) + 4g_A^e g_A^\mu \cos\theta]$$

$$F_Z(\cos\theta) = \frac{1}{16 \sin^4\theta_W \cos^4\theta_W} [(g_V^{e^2} + g_A^{e^2})(g_V^{\mu^2} + g_A^{\mu^2})(1 + \cos^2\theta) + 8g_V^e g_A^e g_V^\mu g_A^\mu \cos\theta]$$

Forward-backward asymmetry

Forward-backward asymmetry

$$\frac{d\sigma}{d\cos\theta} \sim (1 + \cos^2\theta) + \frac{8}{3} A_{FB} \cos\theta$$

with

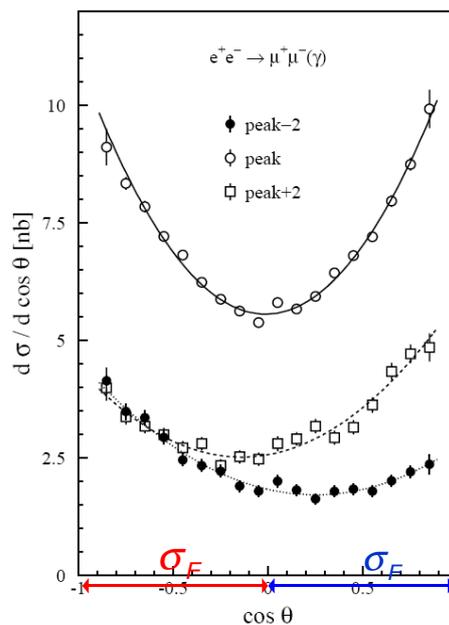
$$A_{FB} = \frac{\sigma_F - \sigma_B}{\sigma_F + \sigma_B}$$

$$\sigma_{F(B)} = \int_{0^{(-1)}}^{1^{(0)}} \frac{d\sigma}{d\cos\theta} d\cos\theta$$

At the Z-pole $\sqrt{s} \approx M_Z$

- Z contribution is dominant
- interference vanishes

$$A_{FB} = 3 \cdot \frac{g_V^e g_A^e}{(g_V^e)^2 + (g_A^e)^2} \cdot \frac{g_V^\mu g_A^\mu}{(g_V^\mu)^2 + (g_A^\mu)^2}$$



Forward-backward asymmetry

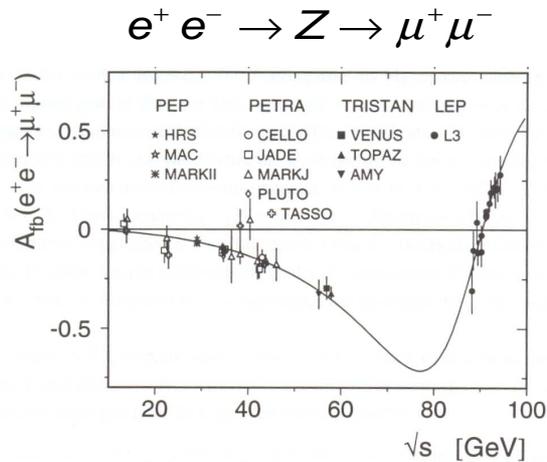
- Away from the resonance A_{FB} is large
→ interference term dominates

$$A_{FB} \sim g_A^e g_A^f \cdot \frac{s(s - M_Z^2)}{(s - M_Z^2)^2 + M_Z^2 \Gamma_Z^2}$$

- At the Z pole: Interference = 0

$$A_{FB} \sim g_A^e g_V^e g_A^f g_V^f$$

→ very small because g_V^f small in SM



Fermion couplings

- Away from the resonance A_{FB} is large
→ interference term dominates

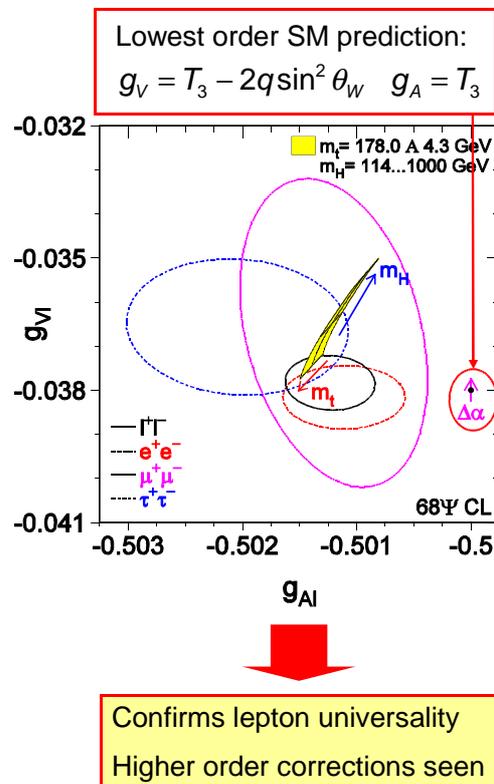
$$A_{FB} \sim g_A^e g_A^f \cdot \frac{s(s - M_Z^2)}{(s - M_Z^2)^2 + M_Z^2 \Gamma_Z^2}$$

- At the Z pole: Interference = 0

$$A_{FB} \sim g_A^e g_V^e g_A^f g_V^f$$

→ very small because g_V^f small in SM

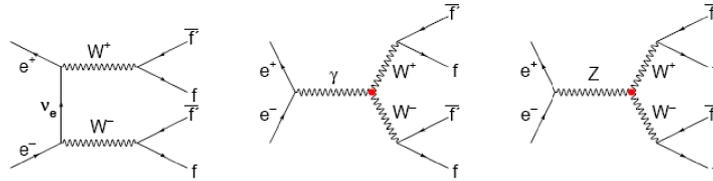
Asymmetries together with cross sections allow the determination of the fermion couplings g_A and g_V



b. Precision tests of the W sector (LEP2 and Tevatron)

$$e^+e^- \rightarrow WW \rightarrow f\bar{f}f\bar{f}$$

~10K WW events / experiment

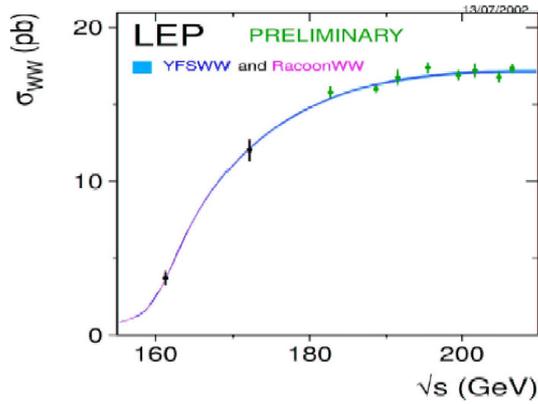


threshold behavior of the cross section (phase space) for $ee \rightarrow WW$ production:

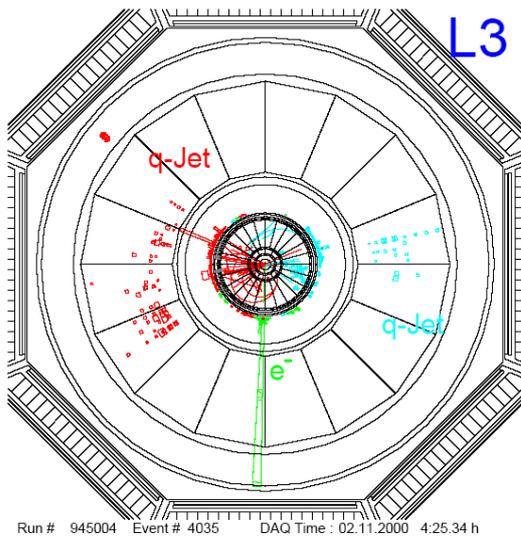


phase space factor = $f(M_W, \sqrt{s})$:

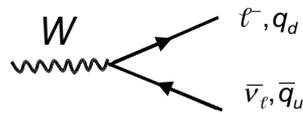
Allows determination of M_W



W decays

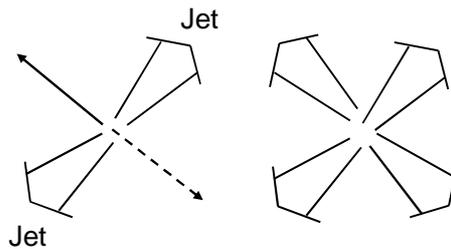


Run# 945004 Event# 4035 DAQ Time : 02.11.2000 4:25.34 h

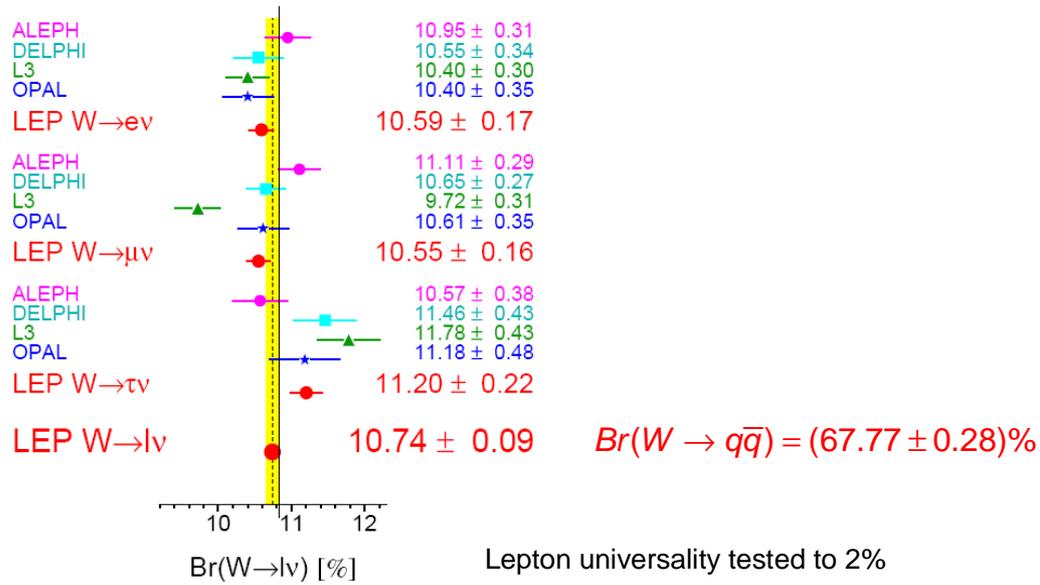


$$WW \rightarrow \begin{cases} qq\bar{l}\nu & 44\% \\ qq\bar{q}q & 45\% \\ l\nu\bar{l}\nu & 11\% \end{cases}$$

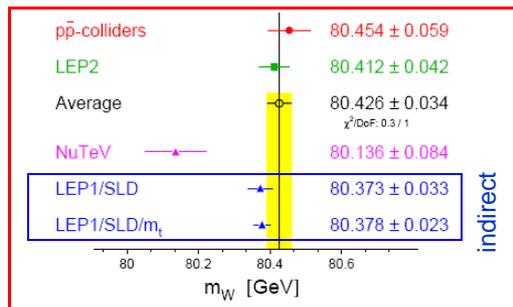
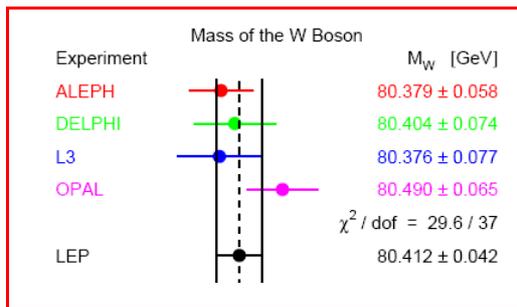
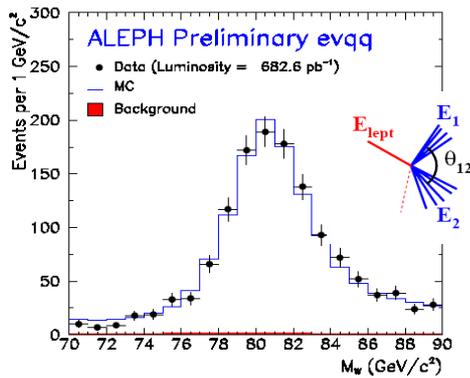
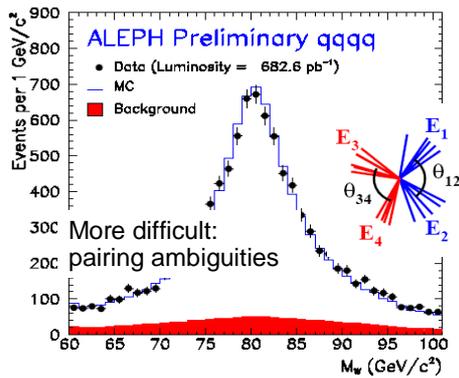
Lepton Neutrino



W branching ratios



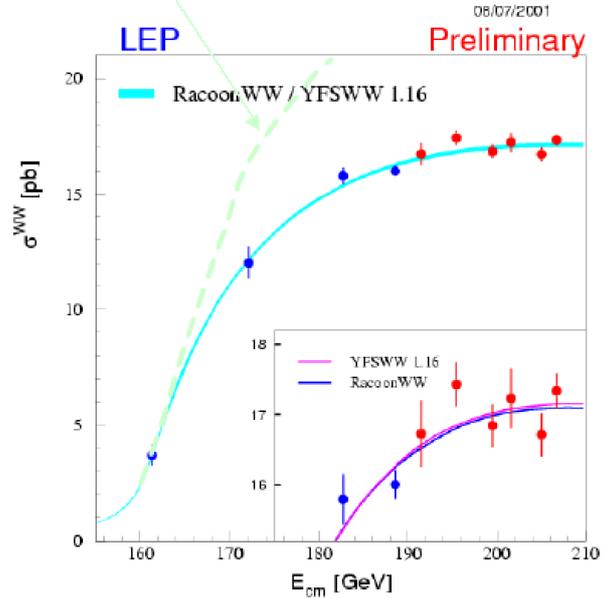
Invariant W mass reconstruction



Effect of triple gauge coupling



Data confirms the existence of the γ/ZWW triple gauge boson vertex

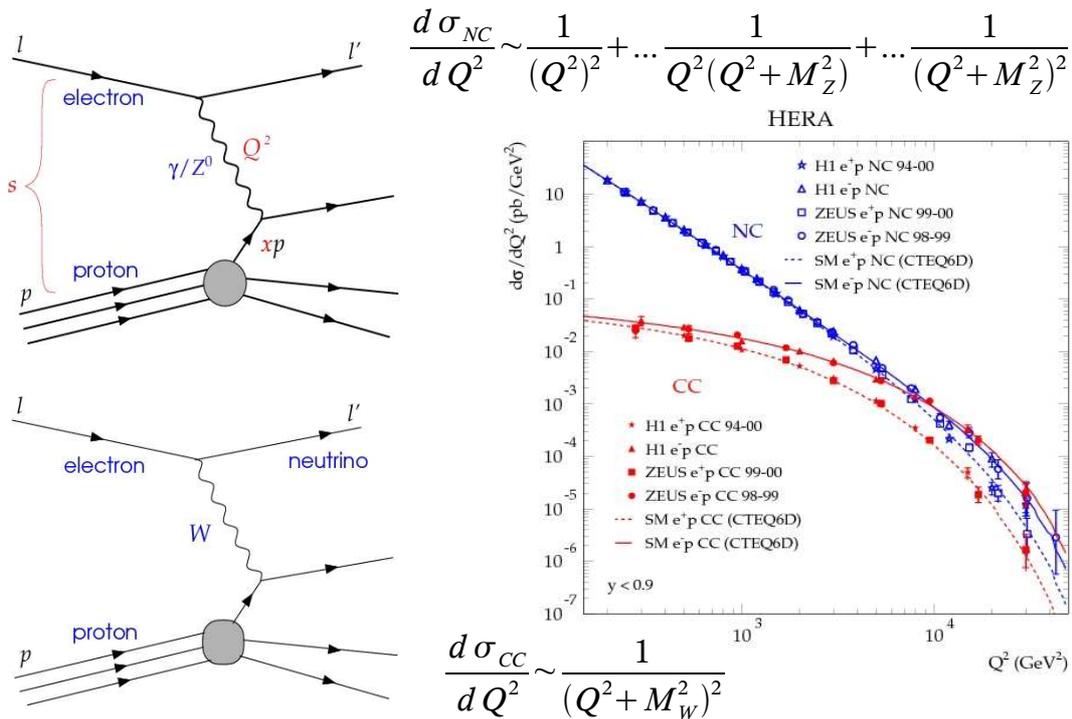


4. Electro-weak unification, as visible at HERA

$$\frac{d\sigma_{NC}}{dQ^2} \sim \underbrace{\frac{1}{(Q^2)^2}}_{\gamma} + \dots + \underbrace{\frac{1}{Q^2(Q^2 + M_Z^2)}}_{\gamma/Z} + \dots + \underbrace{\frac{1}{(Q^2 + M_Z^2)^2}}_Z$$

$$\frac{d\sigma_{CC}}{dQ^2} \sim \frac{1}{(Q^2 + M_W^2)^2}$$

4. Electro-weak unification, as visible at HERA



5. Higher order corrections and the Higgs mass

Lowest order SM predictions

$$\rho = \frac{m_W^2}{m_Z^2 \cos^2 \theta_W} = 1$$

$$\sin^2 \theta_W = 1 - \frac{m_W^2}{m_Z^2}$$

$$m_W^2 = \frac{\pi\alpha}{\sqrt{2} \sin^2 \theta_W G_F}$$

$\alpha(0)$

\Rightarrow

Including radiative corrections

$$\bar{\rho} = 1 + \Delta\rho$$

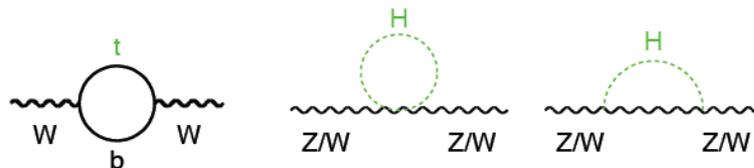
$$\sin^2 \theta_{\text{eff}} = (1 + \Delta\kappa) \sin^2 \theta_W$$

$$m_W^2 = \frac{\pi\alpha}{\sqrt{2} \sin^2 \theta_W G_F} (1 + \Delta r)$$

$$\alpha(m_Z^2) = \frac{\alpha(0)}{1 - \Delta\alpha}$$

with : $\Delta\alpha = \Delta\alpha_{\text{lept}} + \Delta\alpha_{\text{top}} + \Delta\alpha_{\text{had}}^{(5)}$

$\Delta\rho, \Delta\kappa, \Delta r = f(m_t^2, \log(m_H), \dots)$



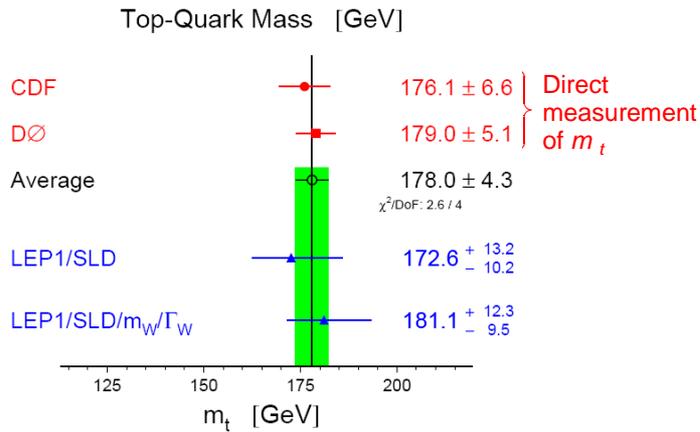
Top mass prediction from radiative corrections

The measurement of the radiative corrections:

$$\sin^2 \theta_{eff} \equiv \frac{1}{4}(1 - \bar{g}_V / \bar{g}_A)$$

$$\sin^2 \theta_{eff} = (1 + \Delta\kappa) \sin^2 \theta_w$$

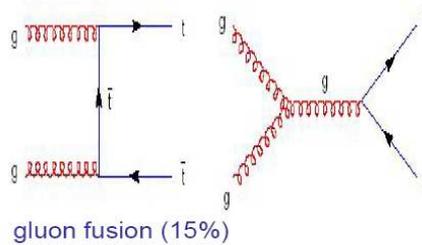
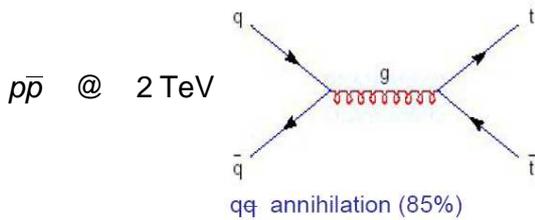
allows an indirect determination of the unknown parameters m_t and M_H .



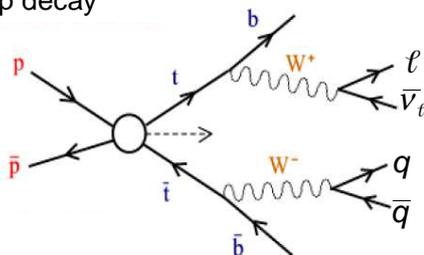
Prediction of m_t by LEP before the discovery of the top at TEVATRON.

Good agreement between the indirect prediction of m_t and the value obtained in direct measurements confirm the radiative corrections of the SM

Observation of the top quark at TEVATRON (1995)

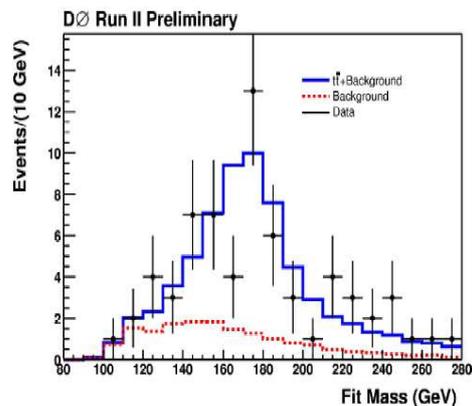


Top decay

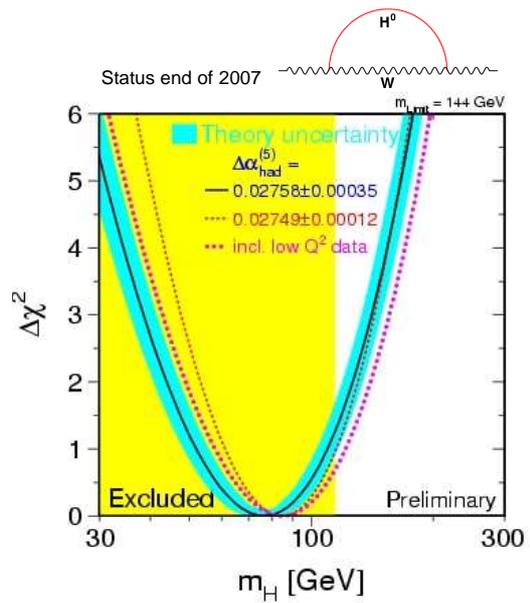
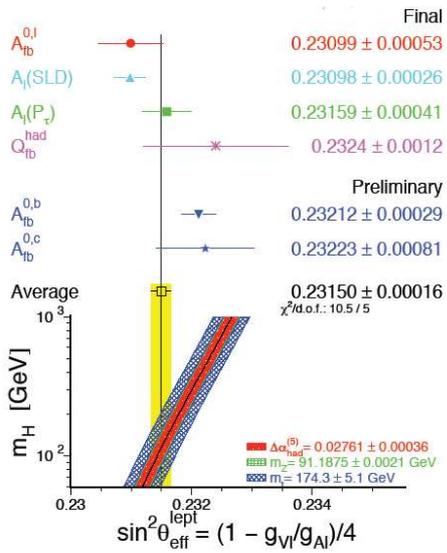


Channel used for mass reconstruction:

$$m_t = m_{inv}(b - jet, W \rightarrow jet + jet)$$



Higgs mass prediction from radiative corrections



Awaiting the discovery of the Higgs at the LHC



$M_H > 114 \text{ GeV}$ (from direct searches)

$M_H < 144 \text{ GeV}$ (from EW fits)