

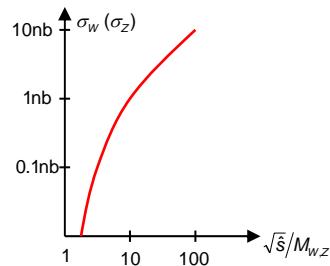
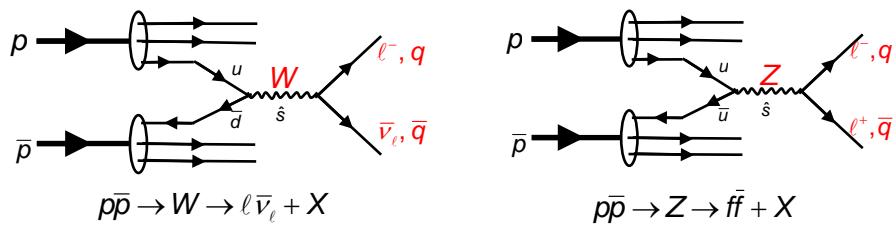
VIII. Experimental tests of the Standard Model

1. Discovery of W and Z boson
2. Precision tests of the Z sector
3. Precision test of the W sector
4. Radiative corrections and prediction of the Higgs mass
5. Higgs searches at the LHC

1. Discovery of the W and Z boson

1983 at CERN SppS accelerator,
 $\sqrt{s} \approx 540$ GeV, UA-1/2 experiments

1.1 Boson production in $p\bar{p}$ interactions



Similar to Drell-Yan: (photon instead of W)

$$\hat{s} = x_q x_{\bar{q}} s \quad \text{mit} \quad \langle x_q \rangle \approx 0.12$$

$$\hat{s} = \langle x_q \rangle^2 s \approx 0.014 s = (65 \text{ GeV})^2$$

→ Cross section is small !

1.2 UA-1 Detector

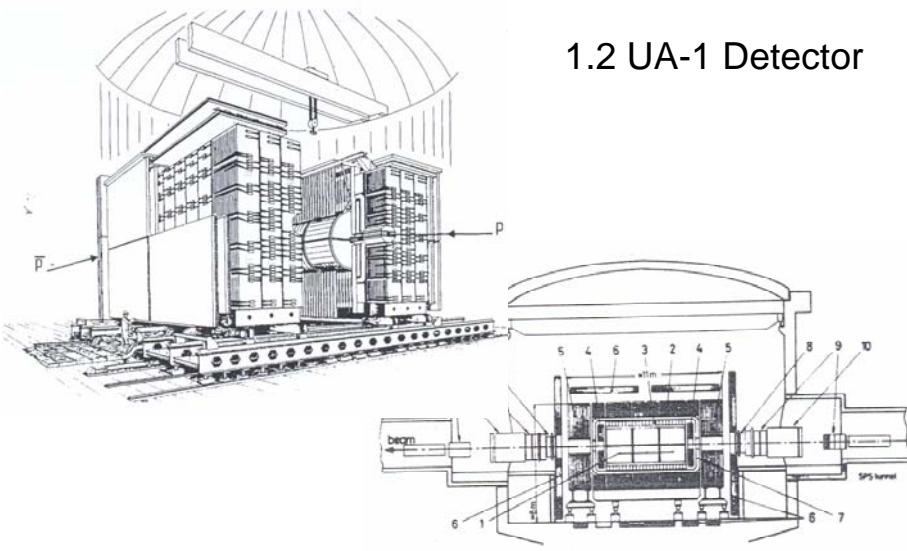
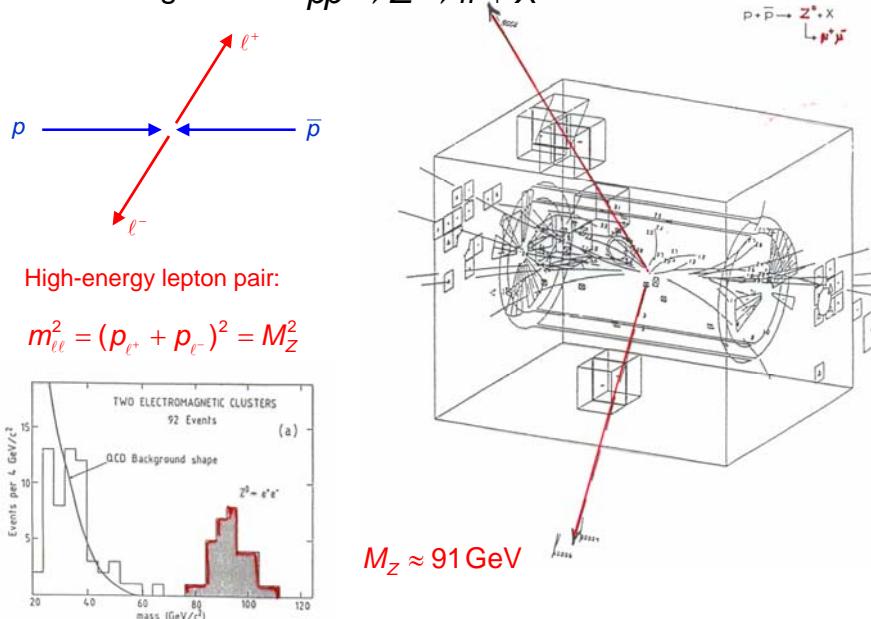


Fig. 8.16: Seitenansicht des UA1-Detektors zum Nachweis von Proton-Antiproton-Wechselwirkungen bei 540 GeV Schwerpunktsenergie: 1. Zentraldetektor, 2. und 5. Hadron-Kalorimeter, 3. und 4. Elektron-Photonen-Schauerzähler, 6. Myon-Detektor, 7. Spule für Dipolfeld, 8. und 9. Kleinwinkeldetektor mit Kammern und Kalorimetern, 10. Kompassator-Magnete [UA1].

1.3 Event signature: $p\bar{p} \rightarrow Z \rightarrow f\bar{f} + X$



1.4 Event signature: $p\bar{p} \rightarrow W \rightarrow \ell \bar{\nu}_\ell + X$

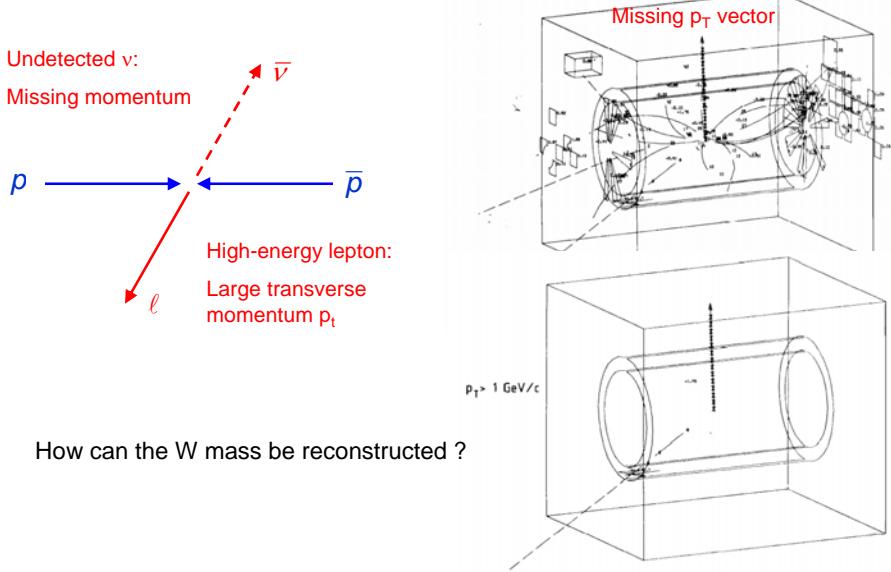
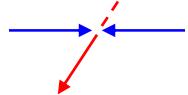


Fig. 16b. The same as picture (a), except that now only particles with $p_T > 1 \text{ GeV}/c$ and calorimeters with $E_T > 1 \text{ GeV}$ are shown.

W mass measurement

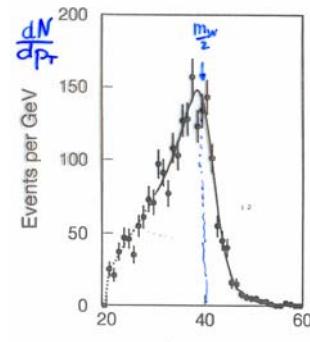
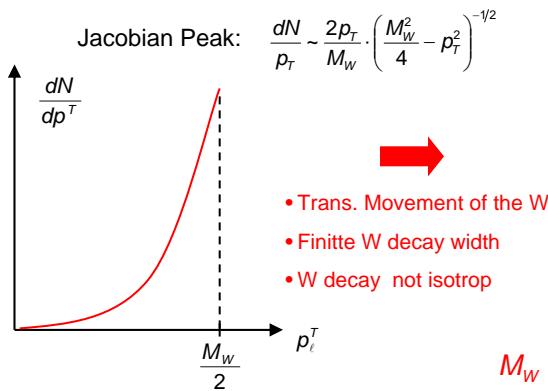
In the W rest frame:

- $|\vec{p}_\ell| = |\vec{p}_\nu| = \frac{M_W}{2}$
- $|p_\ell^T| \leq \frac{M_W}{2}$



In the lab system:

- W system boosted only along z axis
- p_T distribution is conserved





The Nobel Prize in Physics 1984



Carlo Rubbia

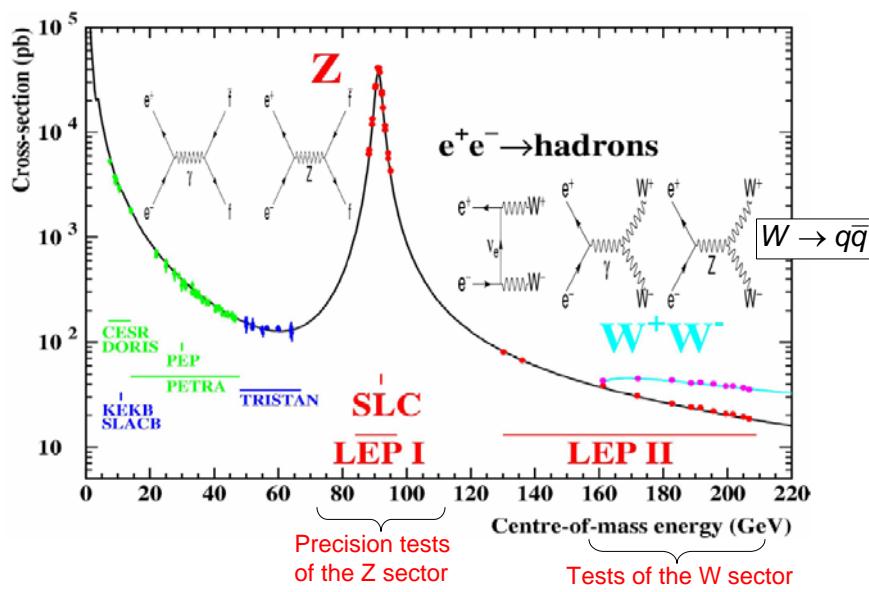
Simon van der Meer

"for their decisive contributions to the large project, which led to the discovery of the field particles W and Z, communicators of weak interaction"

S. van der Meer

One of the achievements to allow high-intensity $p\bar{p}$ collisions, is stochastic cooling of the \bar{p} beams before inserting them into SPS.

1.5 Production of Z and W bosons in e^+e^- annihilation



2. Precision tests of the Z sector

2.1 Cross section for $e^+ e^- \rightarrow \gamma / Z \rightarrow f\bar{f}$

(LEP and SLC)

$\sim 4.5M$ Z decays / experiment

$$|M|^2 = \left| \begin{array}{c} \text{Feynman diagram for } \gamma \\ \text{Feynman diagram for } Z \end{array} \right|^2$$

for $e^+ e^- \rightarrow \mu^+ \mu^-$

$$M_\gamma = -e^2 (\bar{\mu} \gamma^\mu \mu) \frac{1}{q^2} (\bar{e} \gamma^\mu e)$$

$$M_Z = -\frac{g^2}{\cos^2 \theta_W} \left[\bar{\mu} \gamma^\nu \frac{1}{2} (g_V^\mu - g_A^\mu \gamma^5) \mu \right] \underbrace{\frac{g_{\nu\rho} - q_\nu q_\rho/M_Z^2}{(q^2 - M_Z^2) + iM_Z \Gamma_Z}}_{Z \text{ propagator considering a finite } Z \text{ width}} \left[\bar{e} \gamma^\rho \frac{1}{2} (g_V^e - g_A^e \gamma^5) e \right]$$

Z propagator considering
a finite Z width

One finds for the differential cross section:

$$\frac{d\sigma}{d\cos\theta} = \frac{\pi \alpha^2}{2s} \left[F_\gamma(\cos\theta) + F_{\gamma Z}(\cos\theta) \underbrace{\frac{s(s-M_Z^2)}{(s-M_Z^2)^2 + M_Z^2 \Gamma_Z^2}}_{\gamma/Z \text{ interference}} + F_Z(\cos\theta) \underbrace{\frac{s^2}{(s-M_Z^2)^2 + M_Z^2 \Gamma_Z^2}}_{\text{Vanishes at } \sqrt{s} \approx M_Z} \right]$$

$$F_\gamma(\cos\theta) = Q_e^2 Q_\mu^2 (1 + \cos^2\theta) = (1 + \cos^2\theta)$$

$$F_{\gamma Z}(\cos\theta) = \frac{Q_e Q_\mu}{4 \sin^2 \theta_W \cos^2 \theta_W} [2g_V^e g_V^\mu (1 + \cos^2\theta) + 4g_A^e g_A^\mu \cos\theta]$$

$$F_Z(\cos\theta) = \frac{1}{16 \sin^4 \theta_W \cos^4 \theta_W} [(g_V^{e^2} + g_A^{e^2})(g_V^{\mu^2} + g_A^{\mu^2}) (1 + \cos^2\theta) + 8g_V^e g_A^e g_V^\mu g_A^\mu \cos\theta]$$

Forward-backward asymmetry

$$\frac{d\sigma}{d\cos\theta} \sim (1 + \cos^2\theta) + \frac{8}{3} A_{FB} \cos\theta \quad \text{with} \quad A_{FB} = \frac{\sigma_F - \sigma_B}{\sigma_F + \sigma_B}$$

At the Z-pole $\sqrt{s} \approx M_Z$ → Z contribution is dominant
 → interference vanishes

$$\sigma_{tot} \approx \sigma_Z = \frac{4\pi}{3s} \frac{\alpha^2}{16 \sin^4 \theta_w \cos^4 \theta_w} \cdot [g_V^{e^2} + g_A^{e^2}] [g_V^{e^2} + g_A^{e^2}] \cdot \underbrace{\frac{s^2}{(s - M_Z^2)^2 + (M_Z \Gamma_Z)^2}}$$

“Massive propagator” → 1 for $M_Z \rightarrow 0$

$$\sigma_Z(\sqrt{s} = M_Z) = \frac{12\pi}{M_Z^2} \frac{\Gamma_e \Gamma_\mu}{\Gamma_Z^2}$$

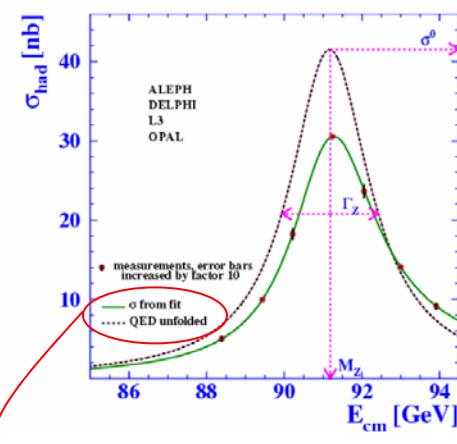
With partial and total widths:

$$\Gamma_f = \frac{\alpha M_Z}{4 \sin^2 \theta_w \cos^2 \theta_w} \cdot [g_V^{f^2} + g_A^{f^2}]$$

$$\Gamma_Z = \sum_i \Gamma_i$$

Cross sections and widths can be calculated within the Standard Model if all parameters are known

2.2 Measurement of the Z lineshape



Resonance curve:

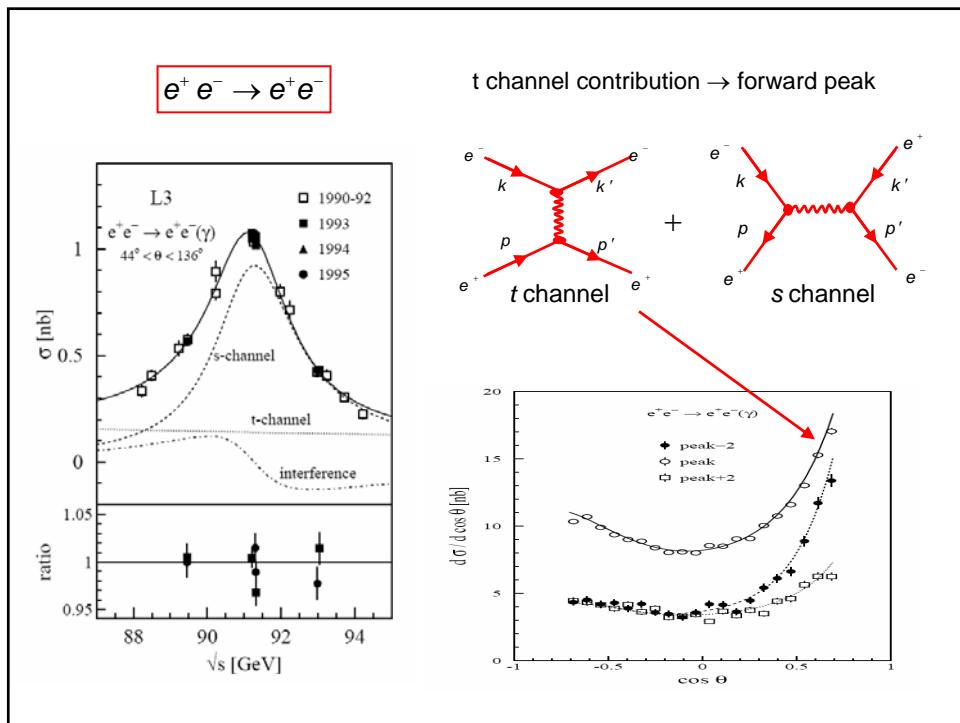
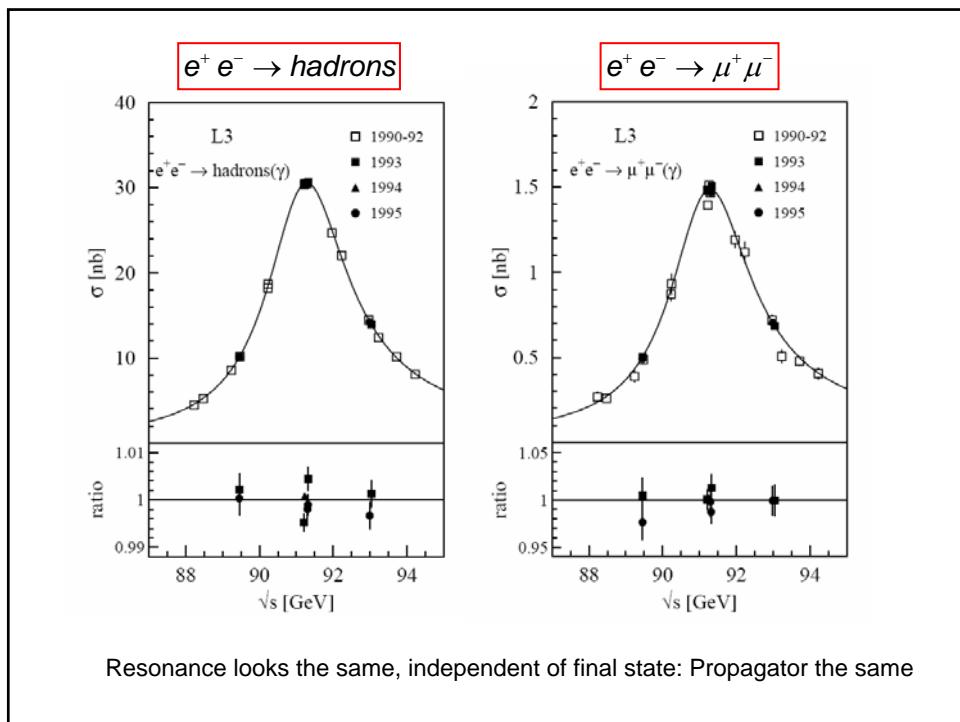
$$\sigma(s) = 12\pi \frac{\Gamma_e \Gamma_\mu}{M_Z^2} \cdot \frac{s}{(s - M_Z^2)^2 + M_Z^2 \Gamma_Z^2}$$

$$\text{Peak: } \sigma_0 = \frac{12\pi}{M_Z^2} \frac{\Gamma_e \Gamma_\mu}{\Gamma_Z^2}$$

- Resonance position → M_Z
- Height → $\Gamma_e \Gamma_\mu$
- Width → Γ_Z

Initial state Bremsstrahlung corrections

$$\sigma_{ff(\gamma)} = \int_{4m_\gamma^2/s}^1 G(z) \sigma_{ff}^0(zs) dz \quad z = 1 - \frac{2E_\gamma}{\sqrt{s}}$$



Z line shape parameters (LEP average)

$$M_Z = 91.1876 \pm 0.0021 \text{ GeV} \quad \pm 23 \text{ ppm (*)}$$

Γ_Z	$= 2.4952 \pm 0.0023 \text{ GeV}$	}	±0.09 % 3 leptons are treated independently
Γ_{had}	$= 1.7458 \pm 0.0027 \text{ GeV}$		
Γ_e	$= 0.08392 \pm 0.00012 \text{ GeV}$		
Γ_μ	$= 0.08399 \pm 0.00018 \text{ GeV}$		
Γ_τ	$= 0.08408 \pm 0.00022 \text{ GeV}$		

Γ_Z	$= 2.4952 \pm 0.0023 \text{ GeV}$	}	Assuming lepton universality: $\Gamma_e = \Gamma_\mu = \Gamma_\tau$
Γ_{had}	$= 1.7444 \pm 0.0022 \text{ GeV}$		
Γ_e	$= 0.083985 \pm 0.000086 \text{ GeV}$		

↑
↓

test of lepton universality

*) error of the LEP energy determination: ±1.7 MeV (19 ppm)

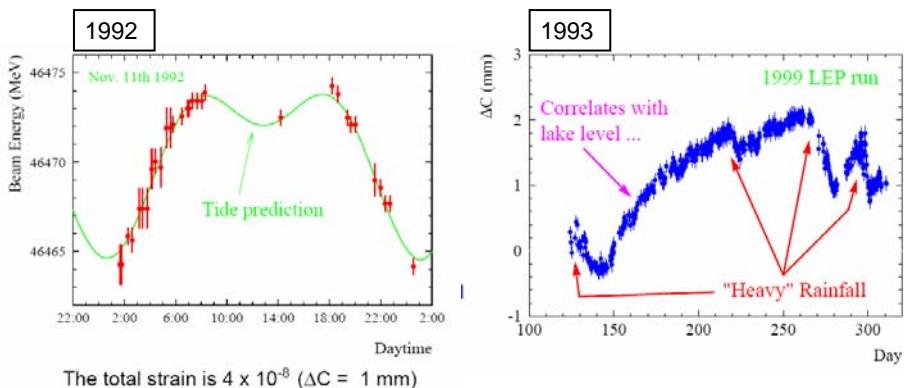
<http://lepewwg.web.cern.ch/> (Summer 2005)

LEP energy calibration: Hunting for ppm effects

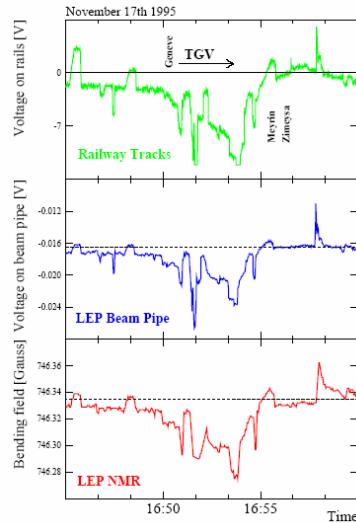
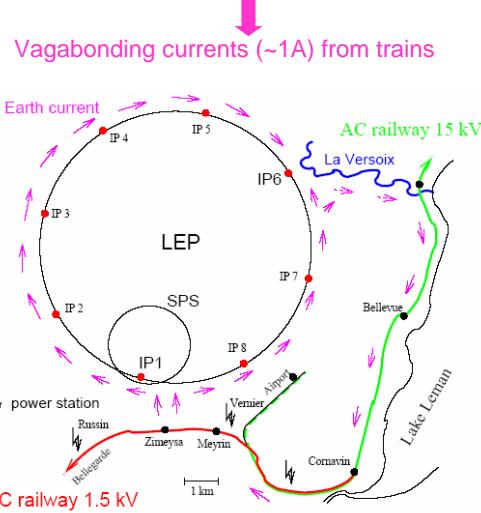
Changes of the circumference of the LEP ring changes the energy of the electrons:

- tide effects
- water level in lake Geneva

Changes of LEP circumference
 $\Delta C = 1...2 \text{ mm}/27\text{km} (4...8 \times 10^{-8})$



Effect of the French “Train a Grande Vitesse” (TGV)



In conclusion: Measurements at the ppm level are difficult to perform. Many effects must be considered!

2.3 Number of light neutrino generations

In the Standard Model:

$$\Gamma_Z = \Gamma_Z + 3 \cdot \Gamma_\ell + \underbrace{N_\nu \cdot \Gamma_\nu}_{\text{invisible}} \rightarrow \left\{ \begin{array}{l} e^+ e^- \rightarrow Z \rightarrow \nu_e \bar{\nu}_e \\ e^+ e^- \rightarrow Z \rightarrow \nu_\mu \bar{\nu}_\mu \\ e^+ e^- \rightarrow Z \rightarrow \nu_\tau \bar{\nu}_\tau \end{array} \right.$$

$$\Gamma_{inv} = 0.4990 \pm 0.0015 \text{ GeV}$$

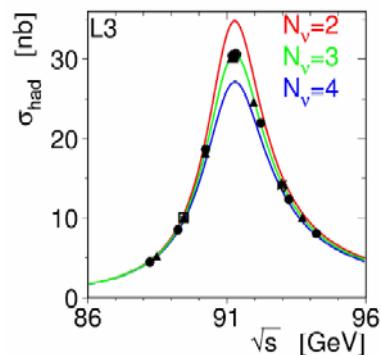
To determine the number of light neutrino generations:

$$N_\nu = \underbrace{\left(\frac{\Gamma_{inv}}{\Gamma_\ell} \right)}_{\exp} \cdot \underbrace{\left(\frac{\Gamma_\ell}{\Gamma_\nu} \right)}_{SM}$$

$$5.9431 \pm 0.0163 = 1.991 \pm 0.001 \text{ (small theo. uncertainties from } m_{top} M_H)$$

$$N_\nu = 2.9840 \pm 0.0082$$

No room for new physics: $Z \rightarrow \text{new}$



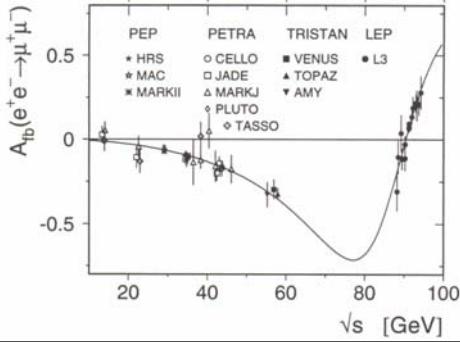
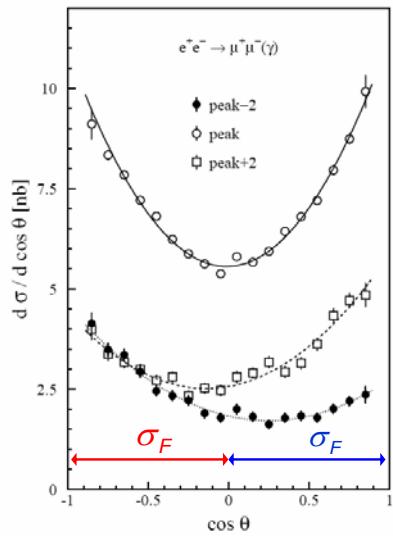
2.4 Forward-backward asymmetry and fermion couplings to Z

$$e^+ e^- \rightarrow Z \rightarrow \mu^+ \mu^-$$

$$\frac{d\sigma}{d\cos\theta} \sim (1 + \cos^2\theta) + \frac{8}{3} A_{FB} \cos\theta$$

with $A_{FB} = \frac{\sigma_F - \sigma_B}{\sigma_F + \sigma_B}$

$$\sigma_{F(B)} = \int_{0(-1)}^{1(0)} \frac{d\sigma}{d\cos\theta} d\cos\theta$$



Fermion couplings

Forward-backward asymmetry

- Away from the resonance A_{FB} large
→ interference term dominates

$$A_{FB} \sim g_A^e g_A^f \cdot \frac{s(s - M_Z^2)}{(s - M_Z^2)^2 + M_Z^2 \Gamma_Z^2}$$

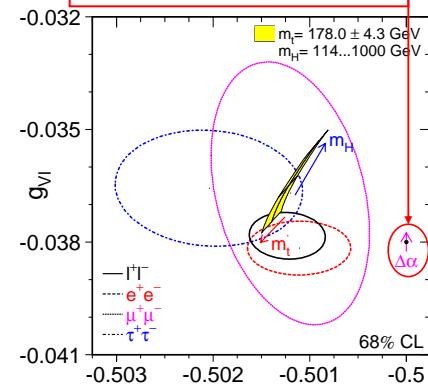
- At the Z pole: Interference = 0

$$A_{FB} \sim g_A^e g_V^e g_A^f g_V^f$$

→ very small because g_V^f small in SM

Lowest order SM prediction:

$$g_V = T_3 - 2q \sin^2 \theta_W \quad g_A = T_3$$



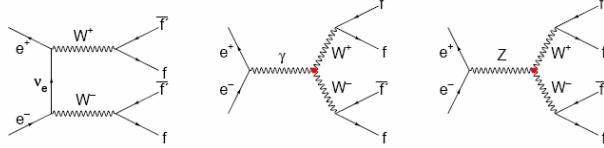
Asymmetries together with cross sections allow the determination of the fermion couplings g_A and g_V

Confirms lepton universality
Higher order corrections seen

3. Precision tests of the W sector (LEP and Tevatron)

$$e^+ e^- \rightarrow WW \rightarrow f\bar{f}f\bar{f}$$

↑ ~10K WW events / experiment

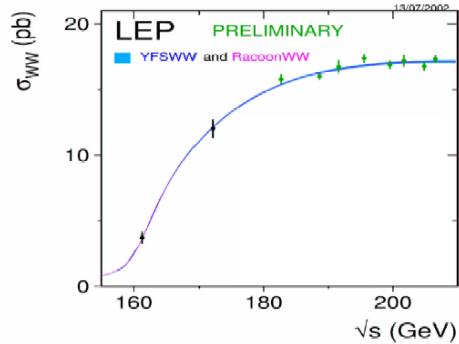


Threshold behavior of the cross section (phase space) for $e^+e^- \rightarrow WW$ production:

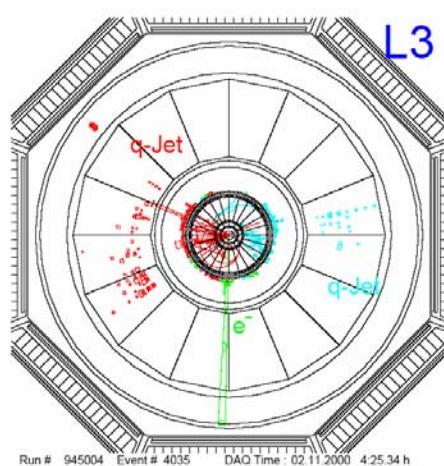


Phase space factor = $f(M_W, \sqrt{s})$:

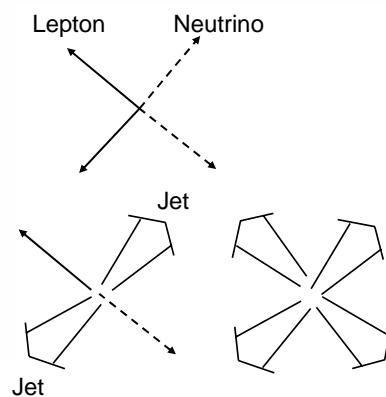
→ Allows determination of M_W



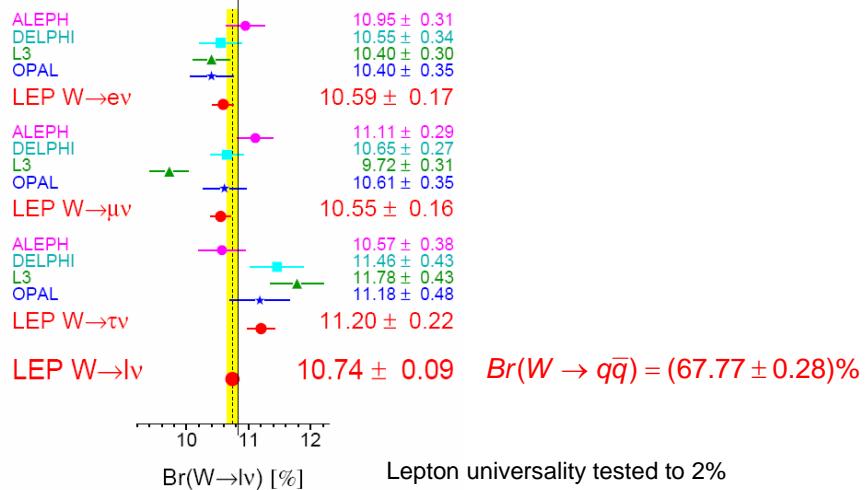
W decays



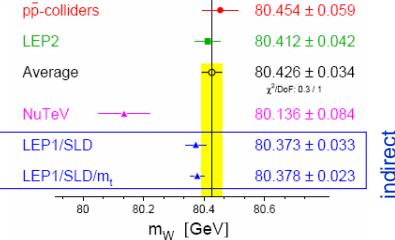
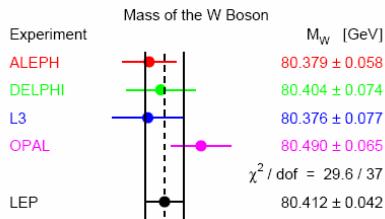
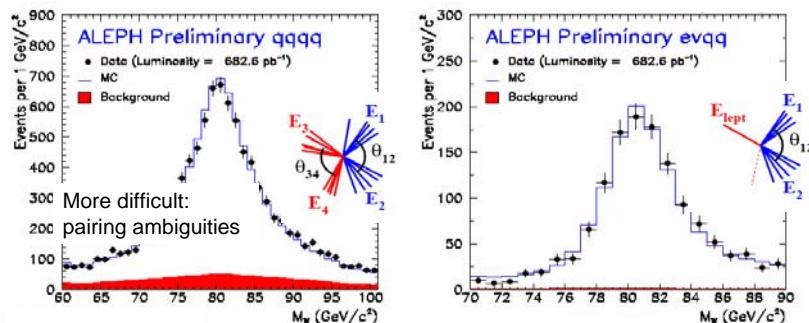
$$WW \rightarrow \begin{cases} q\bar{q}\ell\nu & 44\% \\ q\bar{q}q\bar{q} & 45\% \\ \ell\nu\ell\nu & 11\% \end{cases}$$



W branching ratios



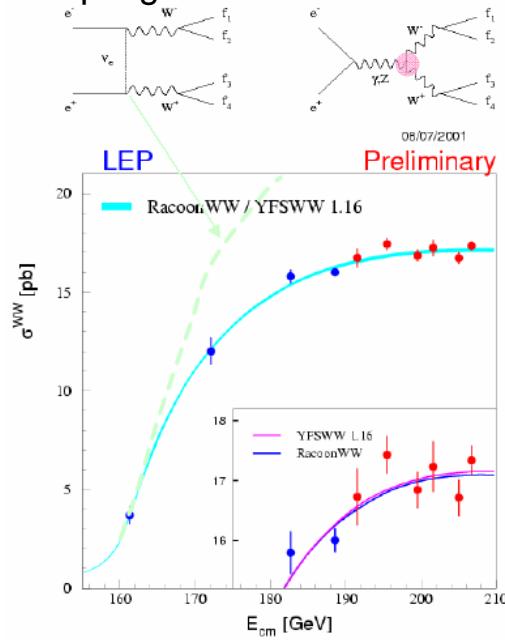
Invariant W mass reconstruction



Effect of triple gauge coupling



Data confirms the existence of the γ/ZWW triple gauge boson vertex



4. Higher order corrections and the Higgs mass

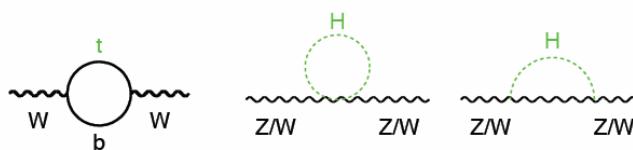
Lowest order SM predictions

$$\begin{aligned} \rho &= \frac{m_W^2}{m_Z^2 \cos^2 \theta_W} = 1 \\ \sin^2 \theta_W &= 1 - \frac{m_W^2}{m_Z^2} \\ m_W^2 &= \frac{\pi \alpha}{\sqrt{2} \sin^2 \theta_W G_F} \end{aligned} \quad \begin{aligned} \Rightarrow & \bar{\rho} = 1 + \Delta \rho \\ \Rightarrow & \sin^2 \theta_{\text{eff}} = (1 + \Delta \kappa) \sin^2 \theta_W \\ \Rightarrow & m_W^2 = \frac{\pi \alpha}{\sqrt{2} \sin^2 \theta_W G_F} (1 + \Delta r) \\ \Rightarrow & \alpha(m_Z^2) = \frac{\alpha(0)}{1 - \Delta \alpha} \end{aligned}$$

Including radiative corrections

with : $\Delta \alpha = \Delta \alpha_{\text{lept}} + \Delta \alpha_{\text{top}} + \Delta \alpha_{\text{had}}^{(5)}$

$$\Delta \rho, \Delta \kappa, \Delta r = f(m_t^2, \log(m_H), \dots)$$



Top mass prediction from radiative corrections

The measurement of the radiative corrections:

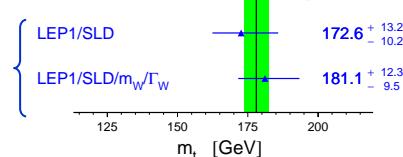
$$\sin^2 \theta_{\text{eff}} \equiv \frac{1}{4} (1 - \bar{g}_V / \bar{g}_A)$$

$$\sin^2 \theta_{\text{eff}} = (1 + \Delta \kappa) \sin^2 \theta_w$$

Allows the indirect determination of the unknown parameters m_t and M_H .

Top-Quark Mass [GeV]

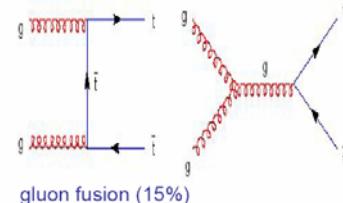
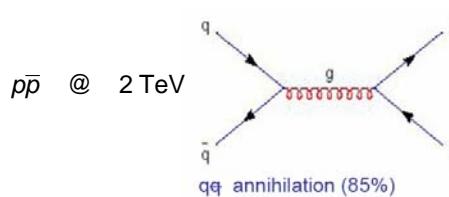
CDF
DØ
Average



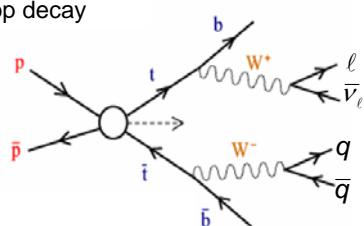
Prediction of m_t by LEP before the discovery of the top at TEVATRON.

Good agreement between the indirect prediction of m_t and the value obtained in direct measurements confirm the radiative corrections of the SM

Observation of the top quark at TEVATRON (1995)

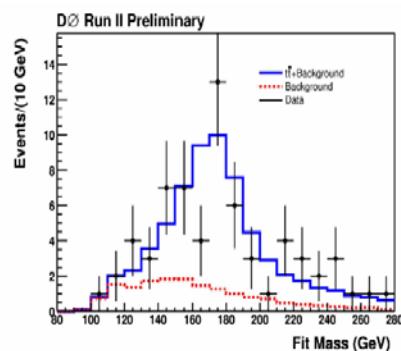


Top decay

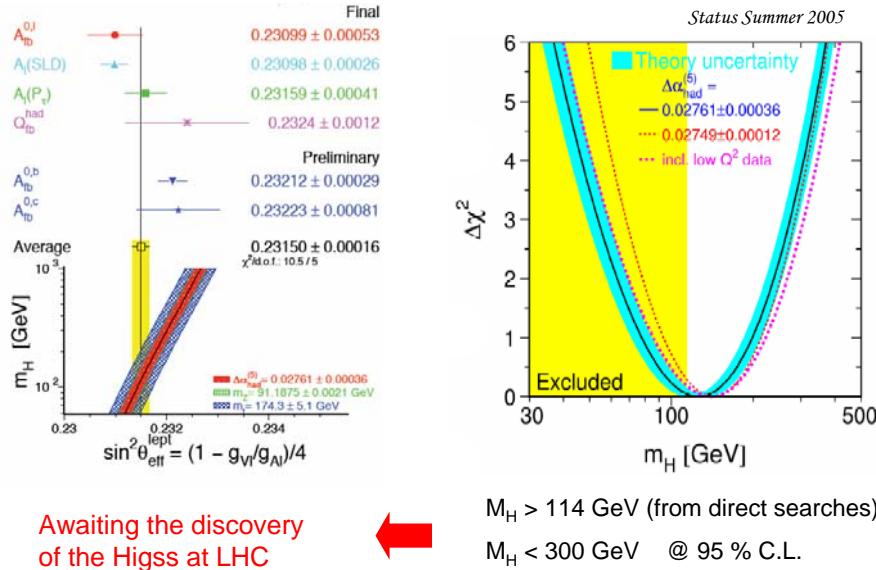


Channel used for mass reconstruction:

$$m_t = m_{\text{inv}}(b - \text{jet}, W \rightarrow \text{jet} + \text{jet})$$

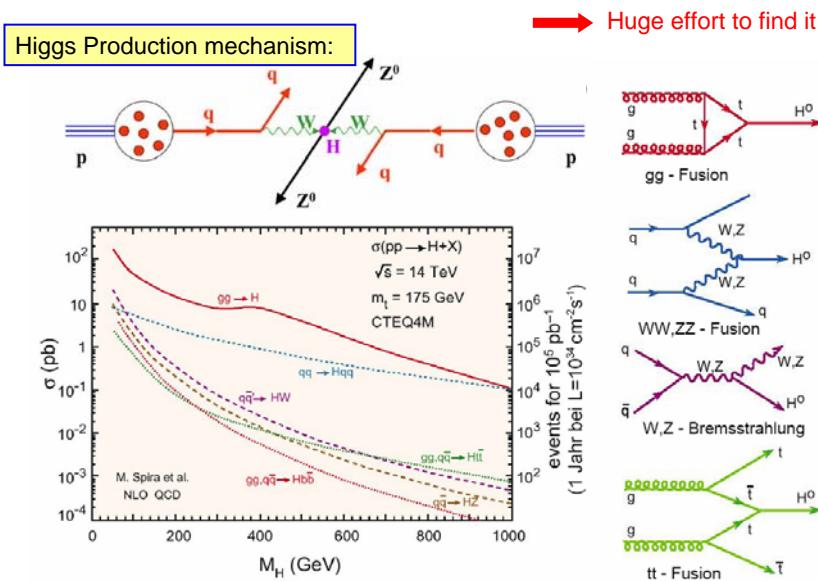


Higgs mass prediction from radiative corrections



5. Higgs searches at the LHC (pp collider @ 14 TeV)

Only missing ingredient of the Standard Model: Higgs-Boson



Higgs decay channels

At LEP: Searches were done using

$$H \rightarrow b\bar{b} \quad M_H > 114 \text{ GeV}$$

("golden" Higgs decay channel at LEP energies because of large b mass
 → too much background at LHC)

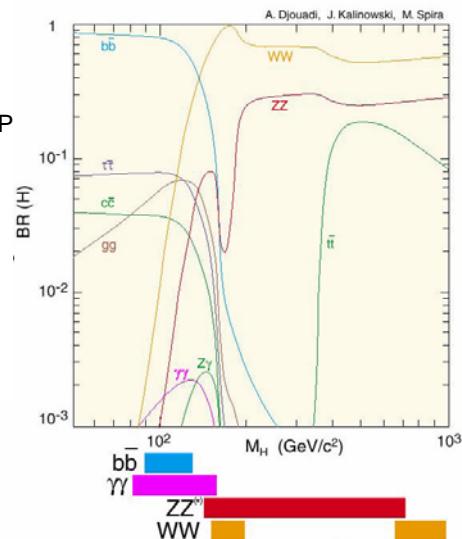
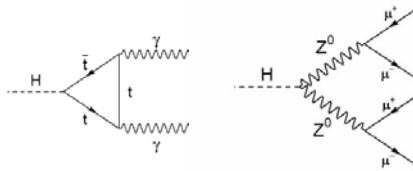
At LHC:

- $m_H < 150 \text{ GeV}: \quad H \rightarrow \gamma\gamma$

- $150 \text{ GeV} < m_H < 1 \text{ TeV}$

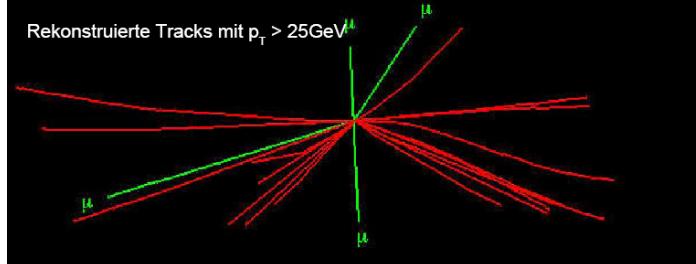
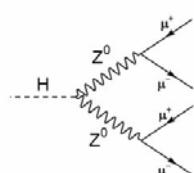
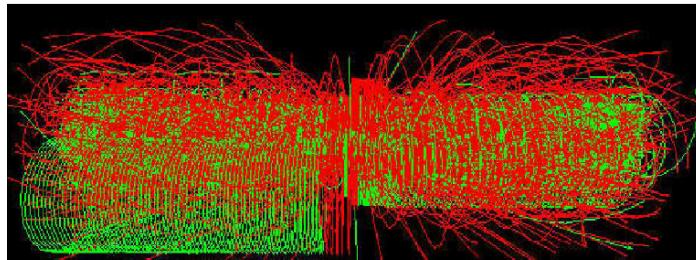
$$H \rightarrow ZZ^{(*)}$$

$$H \rightarrow W^+W^-$$



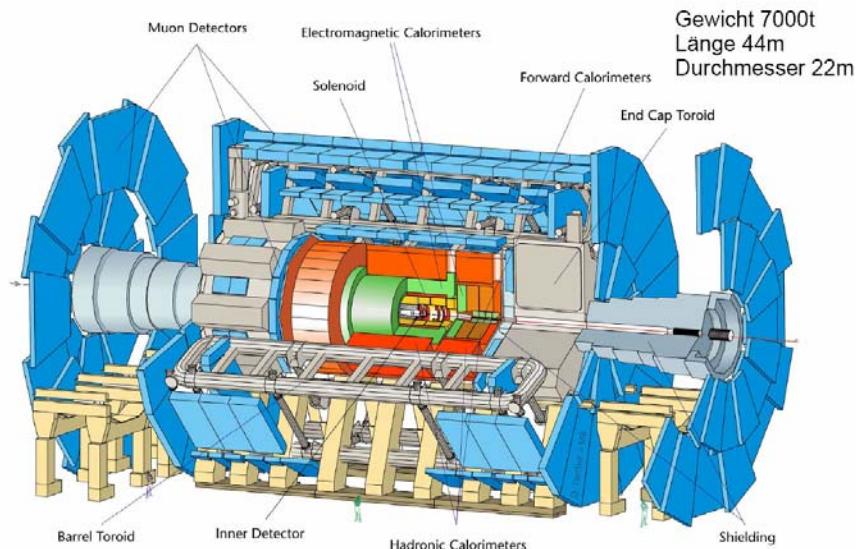
Simulated $H \rightarrow ZZ \rightarrow 4\mu$ event at LHC

- 20 pp interaction / event
- Large number of particles



To trigger and to reconstruct these events is an exp. challenge.

The ATLAS Experiment at LHC



ATLAS Construction



Higgs discovery potential (ATLAS experiment)

