

f) Discrete Symmetries

parity P , charge conjugation C , time reversal T

ii) Parity: $x = (x_0, \vec{x}) \mapsto x_p = (x_0, -\vec{x})$

\hat{P} : unitary operator

* scalar field: $\hat{P} \phi(x) \hat{P}^{-1} = \eta_P \phi(x_p)$

η_P : intrinsic parity, $|\eta_P| = 1$

ϕ hermitian: $\eta_P = +1$ scalar
 $\eta_P = -1$ pseudo-scalar

(for non-hermitian fields, the hermiticity is restored on a class of ϕ)

in modes: $\hat{P} a(p) \hat{P}^{-1} = \eta_P a(p_p)$
 $\hat{P} b^\dagger(p) \hat{P}^{-1} = \eta_P b^\dagger(p_p)$

$\hat{P} |0\rangle = |0\rangle \Rightarrow \hat{P} |p\rangle = \eta_P^* |p_p\rangle$

* vector fields: $\hat{P} V^\mu(x) \hat{P}^{-1} = \eta_P V_\mu(x_p)$

V hermitian: $\eta_P = 1$ vector
 $\eta_P = -1$ pseudo/axial vector

* Dirac field: $\hat{P} \psi(x) \hat{P}^{-1} = \eta_P \gamma^0 \psi(x_p)$

$\gamma^0 \leftrightarrow$ compatibility with Dirac eqn.

Bilinears:

$\bar{\psi}(x) \psi(x)$	$\rightarrow \bar{\psi}(x_p) \psi(x_p)$	scalar
$\bar{\psi}(x) \gamma_5 \psi(x)$	$\rightarrow -\bar{\psi}(x_p) \gamma_5 \psi(x_p)$	pseudo-scalar
$\bar{\psi}(x) \gamma^\mu \psi(x)$	$\rightarrow \bar{\psi}(x_p) \gamma_\mu \psi(x_p)$	vector
$\bar{\psi}(x) \gamma^\mu \gamma_5 \psi(x)$	$\rightarrow -\bar{\psi}(x_p) \gamma_\mu \gamma_5 \psi(x_p)$	axial vector

ii) Charge conjugation

$$\hat{C} |p, \text{particle}\rangle = \eta_c^* |p, \text{anti-particle}\rangle$$

\hat{C} : unitary
 $|\eta_c| = 1$

* scalar: $\hat{C} \phi(x) \hat{C}^{-1} = \eta_c \phi(x)^\dagger$
 $\hat{C} \phi^\dagger(x) \hat{C}^{-1} = \eta_c^* \phi(x)$

ϕ convention: $\eta_c = \pm 1$ (otherwise η_c absorbed into phase of $\underline{\phi}$)

* Dirac fermion

Ansatz: $\psi(x) \rightarrow \eta_c \psi^c(x) = \eta_c C \bar{\psi}(x)^t$

w. C s.t. ψ^c satisfies Dirac eqn.

Result: need $C (\gamma^\mu)^t C^{-1} = -\gamma^\mu$

$\Rightarrow C^T = C^{-1}, C^T = -C$ (see tutorial)

\hookrightarrow Dirac repr.: $C = i\gamma^0\gamma^2 = \begin{pmatrix} 0 & i\sigma^2 \\ i\sigma^2 & 0 \end{pmatrix}$

$\rightarrow \hat{C} \psi(x) \hat{C}^{-1} = \eta_c C \bar{\psi}^t$

$\hat{C} \bar{\psi}(x) \hat{C}^{-1} = -\eta_c^* \psi^t C^{-1}$

Find: $\hat{C} a(p, \lambda) \hat{C}^{-1} = \eta_c b(p, \lambda)$

$\hat{C} b^\dagger(p, \lambda) \hat{C}^{-1} = \eta_c a^\dagger(p, \lambda)$

Majorana spinor: $\psi(x) = \psi^c(x)$

particle its
anti-particle

Find: current $\bar{\psi}(x) \gamma_\mu \psi(x) = j_\mu$

$\hat{C} j_\mu(x) \hat{C}^{-1} = -j_\mu(x)$

* Vector: $\hat{C} A_\mu(x) \hat{C}^{-1} = -A_\mu$ so that QED
C invariant

C \leftrightarrow helicity:

$$(\psi_L)^c = C \bar{\psi}_L^t = C \cdot \frac{1}{2} (1 + \gamma_5^t) \bar{\psi}^t = \frac{1}{2} (1 + \gamma_5) \psi^c$$
$$(C \gamma_5^t C^{-1} = \gamma_5)$$

$$\boxed{(\psi_L)^c = (\psi^c)_R}$$

charge conjugation,
reverses parity.

Useful identities:

$$C^\dagger = C^{-1}, \quad C^t = -C$$

$$(\gamma^\mu)^t = -C^{-1} \gamma^\mu C$$

$$\psi^c = C \bar{\psi}^*$$

Time reversal

- * exchanges initial & final states with identical position & opposite velocities/momenta
- * Consistent action & S-matrix requires \hat{T} to be anti-linear: $\hat{T}(a|\psi\rangle) = a^* \hat{T}|\psi\rangle$

* Results: * $\hat{T} \phi(x) \hat{T}^{-1} = \eta_T \phi(x_T) \quad x_T = (-x^0, x^2)$

$$\Leftrightarrow \hat{T} a(p) \hat{T}^{-1} = \eta_T a(p_P)$$

* $\hat{T} \psi(x) \hat{T}^{-1} = \eta_T B^{-1} \psi(x_T) \quad B = \gamma^5 C$

$$\Rightarrow \hat{T} \bar{\psi}(x) \psi(x) \hat{T}^{-1} = \bar{\psi}(x_T) \psi(x_T) \quad \text{scalars}$$

$$\hat{T} j^\mu(x) \hat{T}^{-1} = \bar{\psi}(x_T) B(\gamma^\mu)^* B^{-1} \psi(x_T)$$

$$\Rightarrow \hat{T} j^0(x) \hat{T}^{-1} = j^0(x_T)$$

$$\hat{T} j^2 \hat{T}^{-1} = -j^2(x_T)$$

← CPT theorem

Any Lorentz invariant Lagrangian $\mathcal{L}(x)$ formed from products of quantum fields at point x (locality) is invariant under CPT.

- * The axial interaction of the SM violates C & P separately:

$$i.g. \mathcal{L}_+ = j^\mu \not{A}_\mu(x) - j^\mu_A(x) A_\mu(x)$$

$$\xrightarrow{C} j^\mu \not{A}_\mu + j^\mu_A A_\mu$$

* The complex phase of the CKM matrix is responsible for violation of CP:

$$\mathcal{L}_F = -\frac{g}{2\sqrt{2}} \left(\bar{u}_i \gamma^\mu (1-\gamma^5) V_{ij} d_j W_\mu^- + \bar{d}_j \gamma^\mu (1-\gamma^5) V_{ij}^* u_i W_\mu^+ \right)$$

CP: $\bar{u}_i \gamma^\mu (1-\gamma^5) d_j W_\mu^- \rightarrow \bar{d}_j \gamma^\mu (1-\gamma^5) u_i W_\mu^+$
 $\bar{d}_j \gamma^\mu (1-\gamma^5) u_i W_\mu^+ \rightarrow \bar{u}_i \gamma^\mu (1-\gamma^5) d_j W_\mu^-$

$$\rightarrow \mathcal{L}_F \rightarrow -\frac{g}{2\sqrt{2}} \left(\bar{u}_i \gamma^\mu (1-\gamma^5) V_{ij}^* d_j W_\mu^- + \bar{d}_j \gamma^\mu (1-\gamma^5) V_{ij} u_i W_\mu^+ \right)$$

Since $V_{ij} \neq V_{ij}^* \iff$ CP violation