

IX. Flavor oscillation and CP violation

1. Quark mixing and the CKM matrix
2. Flavor oscillations: Mixing of neutral mesons
3. CP violation
4. Neutrino oscillations

1. Quark mixing and CKM matrix

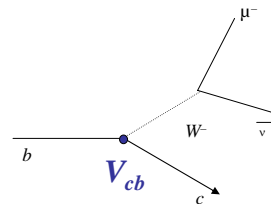
1.1 Quark mixing:

Mass eigenstates are not equal to the weak eigenstates:
Quark-mixing described by unitary **CKM matrix**,

$$\begin{pmatrix} d' \\ s' \\ b' \end{pmatrix} = \underbrace{\begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix}}_{V_{CKM}} \cdot \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$

For weak charged current one obtains:

$$J_{\mu}^{cc} \propto (\bar{u}, \bar{c}, \bar{t}) \gamma_{\mu} (1 - \gamma_5) V_{CKM} \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$



V_{CKM} is unitary:

$$V_{CKM} V_{CKM}^{\dagger} = V_{CKM}^{\dagger} V_{CKM} = 1$$

1.2 Cabbibo-Kobayashi-Maskawa matrix

Number of independent parameters:



18 parameter (9 complex elements)
 -5 relative quark phases (unobservable)
 -9 unitarity conditions

 =4 independent parameters: **3 angles + 1 phase**

PDG parametrization

3 Euler angles

1 Phase

$$\begin{pmatrix} d' \\ s' \\ b' \end{pmatrix} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \begin{pmatrix} c_{13} & 0 & s_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta} & 0 & c_{13} \end{pmatrix} \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$

$$\begin{pmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta} \\ -s_{12}c_{23} - c_{12}s_{23}s_{13}e^{i\delta} & c_{12}c_{23} - s_{12}s_{23}s_{13}e^{i\delta} & s_{23}c_{13} \\ s_{12}s_{23} - c_{12}c_{23}s_{13}e^{i\delta} & -c_{12}s_{23} - s_{12}c_{23}s_{13}e^{i\delta} & c_{23}c_{13} \end{pmatrix}$$

where $c_{ij} = \cos\theta_{ij}$, $s_{ij} = \sin\theta_{ij}$

Unobservable Phases

Phases of left-handed fields in J^{cc} are unobservable: possible redefinition

$$\begin{aligned} u_L &\rightarrow e^{i\phi(u)}u_L & c_L &\rightarrow e^{i\phi(c)}c_L & t_L &\rightarrow e^{i\phi(t)}t_L \\ d_L &\rightarrow e^{i\phi(d)}d_L & s_L &\rightarrow e^{i\phi(s)}s_L & b_L &\rightarrow e^{i\phi(b)}b_L \end{aligned}$$

↑
Real numbers

Under phase transformation:

$$V \rightarrow \begin{pmatrix} e^{-i\phi(u)} & 0 & 0 \\ 0 & e^{-i\phi(c)} & 0 \\ 0 & 0 & e^{-i\phi(t)} \end{pmatrix} \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \begin{pmatrix} e^{i\phi(d)} & 0 & 0 \\ 0 & e^{i\phi(s)} & 0 \\ 0 & 0 & e^{i\phi(b)} \end{pmatrix}$$

$$V_{\alpha j} \rightarrow \exp[i(\phi(j) - \phi(\alpha))] V_{\alpha j}$$

5 unobservable phase differences !

Wolfenstein Parametrization λ, A, ρ, η

$$V_{CKM} = \begin{pmatrix} 1 - \frac{\lambda^2}{2} & \lambda & V_{ub} = |V_{ub}| e^{-i\gamma} \\ -\lambda & 1 - \frac{\lambda^2}{2} & A\lambda^2 \\ A\lambda^3(1 - \rho - i\eta) & -A\lambda^2 & 1 \\ V_{td} = |V_{td}| e^{-i\beta} & & \end{pmatrix} + O(\lambda^4)$$

Magnitude of the CKM elements: $|V_{ij}|$

$$\begin{pmatrix} 0.97383^{+0.00024}_{-0.00023} & 0.2272^{+0.0010}_{-0.0010} & (3.96^{+0.09}_{-0.09}) \times 10^{-3} \\ 0.2271^{+0.0010}_{-0.0010} & 0.97296^{+0.00024}_{-0.00024} & (42.21^{+0.10}_{-0.80}) \times 10^{-3} \\ (8.14^{+0.32}_{-0.64}) \times 10^{-3} & (41.61^{+0.12}_{-0.78}) \times 10^{-3} & 0.999100^{+0.000034}_{-0.000004} \end{pmatrix}$$

Complex CKM elements and CP violation

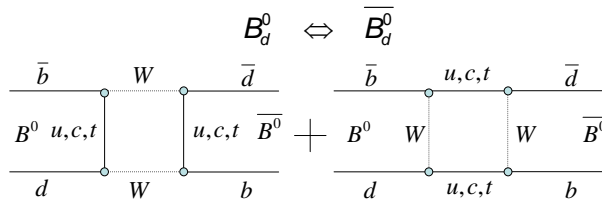
CP (T) violation $\Leftrightarrow V_{ji} \neq V_{ji}^*$
i.e. Complex elements

Remark: For 2 quark generations the mixing is described by the **real 2x2** Cabbibo matrix \rightarrow **no CP violation !!**. To explain **CPV** in the SM Kobayashi and Maskawa have predicted a **third quark generation**.

2. Flavor oscillation: Mixing of neutral mesons

Neutral mesons: $|P^0\rangle: K^0 = |d\bar{s}\rangle \quad D^0 = |\bar{u}c\rangle \quad B_d^0 = |d\bar{b}\rangle \quad B_s^0 = |s\bar{b}\rangle$
 $|\bar{P}^0\rangle: \bar{K}^0 = |\bar{d}s\rangle \quad \bar{D}^0 = |u\bar{c}\rangle \quad \bar{B}_d^0 = |\bar{d}b\rangle \quad \bar{B}_s^0 = |s\bar{b}\rangle$
 discovery 1960 2007 1987 2006

Standard Model predicts oscillations of neutral mesons:



Transition can be described by matrix element: $\langle \bar{B}_d^0 | H_W | B_d^0 \rangle$

2.1 Phenomenological description of mixing

Schrödinger equation for unstable meson:

$$i \frac{d}{dt} |\psi\rangle = H |\psi\rangle = \left(m - \frac{i}{2} \Gamma \right) |\psi\rangle \Rightarrow \begin{cases} |\psi(t)\rangle = |\psi_0\rangle \cdot e^{-imt} \cdot e^{-\frac{1}{2}\Gamma t} \\ \|\psi(t)\|^2 = \|\psi_0\|^2 \cdot e^{-\Gamma t} \end{cases}$$

For neutral mesons $(K^0, \bar{K}^0), (D^0, \bar{D}^0), (B^0, \bar{B}^0), (B_s^0, \bar{B}_s^0)$ consider 2 components

$$i \frac{d}{dt} \begin{pmatrix} P^0 \\ \bar{P}^0 \end{pmatrix} = \mathbf{H} \begin{pmatrix} P^0 \\ \bar{P}^0 \end{pmatrix} = \begin{pmatrix} H_{11} & H_{21} \\ H_{12} & H_{22} \end{pmatrix} \begin{pmatrix} P^0 \\ \bar{P}^0 \end{pmatrix} = \left(\mathbf{M} - \frac{i}{2} \mathbf{\Gamma} \right) \begin{pmatrix} P^0 \\ \bar{P}^0 \end{pmatrix} = \underbrace{\begin{pmatrix} m_{11} - \frac{i}{2} \Gamma_{11} & m_{12} - \frac{i}{2} \Gamma_{12} \\ m_{12} - \frac{i}{2} \Gamma_{12} & m_{22} - \frac{i}{2} \Gamma_{22} \end{pmatrix}}_{\langle P^0 | H_W | \bar{P}^0 \rangle} \begin{pmatrix} P^0 \\ \bar{P}^0 \end{pmatrix}$$

CP $\Rightarrow H_{12} = H_{21}$ CPT $\Rightarrow \begin{matrix} H_{11} = H_{22} = H \\ m_{11} = m_{22} = m \\ \Gamma_{11} = \Gamma_{22} = \Gamma \end{matrix}$

\mathbf{M} and $\mathbf{\Gamma}$ hermitian:

$$\begin{matrix} m_{21} = m_{12} \\ \Gamma_{21} = \Gamma_{12} \end{matrix}$$

Mass eigenstates (by diagonalizing matrix)

Heavy and light mass eigenstate:

$$|P_L\rangle = p|P^0\rangle + q|\bar{P}^0\rangle \quad \text{with } m_L, \Gamma_L$$

$$|P_H\rangle = p|P^0\rangle - q|\bar{P}^0\rangle \quad \text{with } m_H, \Gamma_H$$

$|p|^2 + |q|^2 = 1$ complex coefficients

↓

Flavor eigenstates

$$|P^0\rangle = \frac{1}{2p}(|P_L\rangle + |P_H\rangle)$$

$$|\bar{P}^0\rangle = \frac{1}{2q}(|P_L\rangle - |P_H\rangle)$$

$|P_L\rangle = |P_1\rangle$ and $|P_H\rangle = |P_2\rangle$

Parameters of the mass states

$$m_{H,L} = m \pm \text{Re}\sqrt{H_{12}H_{21}}$$

$$\Gamma_{H,L} = \Gamma \mp 2\text{Im}\sqrt{H_{12}H_{21}}$$

$$\Delta m = m_H - m_L = 2\text{Re}\sqrt{H_{12}H_{21}}$$

$$\Delta\Gamma = \Gamma_H - \Gamma_L = -4\text{Im}\sqrt{H_{12}H_{21}}$$

$$x \equiv \frac{\Delta m}{\Gamma} \quad \text{und} \quad y \equiv \frac{\Delta\Gamma}{2\Gamma}$$

Neutral B mesons

$$\Delta m_d = 0.502 \pm 0.007 \text{ ps}^{-1} \text{ (PDG)}$$

$$\Delta m_s = 17.8 \pm 0.1 \text{ ps}^{-1} \text{ (CDF 2006)}$$

$$\Delta\Gamma / \Gamma < 0.07 \text{ (90\% CL)}$$

$$\Delta\Gamma_s / \Gamma_s = 0.09 \pm 0.64 \text{ (Moriond07)}$$

Time evolution

Generic particle (before P) →

$$|B_{H,L}, t\rangle = b_{H,L}(t) |B_{H,L}\rangle \quad \text{mit} \quad b_{H,L}(t) = e^{-\Gamma_{H,L}t} e^{-im_{H,L}t}$$

$$|\psi_{B^0}(t)\rangle = \frac{|B_L, t\rangle + |B_H, t\rangle}{2p} = \frac{1}{2p} \left(b_L(t) \cdot (p|B^0\rangle + q|\bar{B}^0\rangle) + b_H(t) \cdot (p|B^0\rangle - q|\bar{B}^0\rangle) \right)$$

$$= f_+(t) \cdot |B^0\rangle - \frac{q}{p} f_-(t) \cdot |\bar{B}^0\rangle$$

$$|\psi_{\bar{B}^0}(t)\rangle = f_+(t) \cdot |\bar{B}^0\rangle + \frac{p}{q} f_-(t) \cdot |B^0\rangle$$

$$f_{\pm}(t) = \frac{1}{2} \cdot \left[e^{-im_H t} e^{-\Gamma_H t/2} \pm e^{-im_L t} e^{-\Gamma_L t/2} \right]$$

B^0

$P(B^0 \rightarrow B^0, t) = |f_+(t)|^2$
 $P(B^0 \rightarrow \bar{B}^0, t) = \left| \frac{q}{p} \right|^2 |f_-(t)|^2$

\bar{B}^0

$P(\bar{B}^0 \rightarrow \bar{B}^0, t) = |f_+(t)|^2$
 $P(\bar{B}^0 \rightarrow B^0, t) = \left| \frac{p}{q} \right|^2 |f_-(t)|^2$

Mixing of neutral mesons

Time evolution of rates:

$$P(B^0 \rightarrow B^0) = P(\bar{B}^0 \rightarrow \bar{B}^0) = \frac{1}{4} \left[e^{-\Gamma_L t} + e^{-\Gamma_H t} + 2e^{-(\Gamma_L + \Gamma_H)t/2} \cos \Delta m t \right]$$

CPT

$$P(B^0 \rightarrow \bar{B}^0) = \frac{1}{4} \left| \frac{q}{p} \right|^2 \left[e^{-\Gamma_L t} + e^{-\Gamma_H t} - 2e^{-(\Gamma_L + \Gamma_H)t/2} \cos \Delta m t \right]$$

$$P(\bar{B}^0 \rightarrow B^0) = \frac{1}{4} \left| \frac{p}{q} \right|^2 \left[e^{-\Gamma_L t} + e^{-\Gamma_H t} - 2e^{-(\Gamma_L + \Gamma_H)t/2} \cos \Delta m t \right]$$

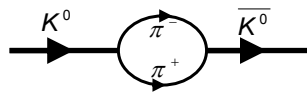
Two mixing mechanisms:

- Mixing through decays $\longrightarrow y = \frac{\Delta\Gamma}{2\Gamma} \approx O(1)$
 - Mixing through oscillation $\longrightarrow x = \frac{\Delta m}{\Gamma} \approx O(1)$
- $(K^0, \bar{K}^0), (D^0, \bar{D}^0), (B^0, \bar{B}^0), (B_s^0, \bar{B}_s^0)$
show different oscillation behavior

CP, T- violation in mixing:

$$P(B^0 \rightarrow \bar{B}^0) \neq P(\bar{B}^0 \rightarrow B^0) \Rightarrow \left| \frac{q}{p} \right| \neq 1$$

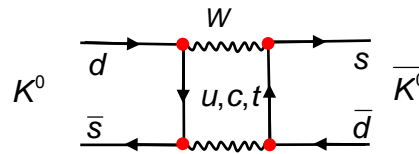
Mixing mechanisms: e.g. K^0



„long distant, on-shell states“

For K^0 important, for B^0 negligible

$$\Rightarrow \Delta\Gamma$$



„short distant, virtual states“

$$\Rightarrow \Delta m$$

Table 1. Parameters of the four neutral oscillating meson pairs [9].

	K^0/\bar{K}^0	D^0/\bar{D}^0	B^0/\bar{B}^0
τ [ps]	$89.4 \pm 0.1;$ 51700 ± 400	$0.413 \pm .003$	1.548 ± 0.021
Γ [s^{-1}]	$5.61 \cdot 10^9$	$2.4 \cdot 10^{12}$	$(6.41 \pm 0.16) \cdot 10^{11}$
$y = \frac{\Delta\Gamma}{2\Gamma}$	-0.9966	$ y < 0.06$	$ y \lesssim 0.01^*$
Δm [s^{-1}]	$(5.300 \pm 0.012) \cdot 10^9$	$< 7 \cdot 10^{10}$	$(4.89 \pm 0.09) \cdot 10^{11}$
Δm [eV]	$(3.49 \pm 0.01) \cdot 10^{-6}$	$< 5 \cdot 10^{-6}$	$(3.2 \pm 0.1) \cdot 10^{-4}$
$x = \frac{\Delta m}{\Gamma}$	0.945 ± 0.002	< 0.03	0.76 ± 0.02

2.2 Neutral kaons

Observation of two neutral kaons K_L (long) and K_S (short) with different lifetimes:

$$\tau(K_L^0) = (51.7 \pm 0.4) \text{ ns} \gg \tau(K_S^0) = (0.089 \pm 0.001) \text{ ns}$$

$$K_L^0 \rightarrow 3\pi$$

$$K_S^0 \rightarrow 2\pi$$

$$CP = -1$$

$$CP = +1$$

K_L and K_S can be identified with the mass eigenstates (ignoring CP violation)

$$|K_L\rangle = |K_2\rangle = \frac{1}{\sqrt{2}}(|K^0\rangle - |\bar{K}^0\rangle)$$

$$CP|K_2\rangle = -|K_2\rangle$$

Phase convention:

$$CP|K^0\rangle = |\bar{K}^0\rangle$$

$$|K_S\rangle = |K_1\rangle = \frac{1}{\sqrt{2}}(|K^0\rangle + |\bar{K}^0\rangle)$$

$$CP|K_1\rangle = +|K_1\rangle$$

$$CP|\bar{K}^0\rangle = |K^0\rangle$$

Large differences between lifetimes

$$\Delta m = (0.5303 \pm 0.0009) \cdot 10^{10} \text{ } \hbar\text{s}^{-1}$$

$$= (3.49 \pm 0.006) \cdot 10^{-12} \text{ MeV}$$

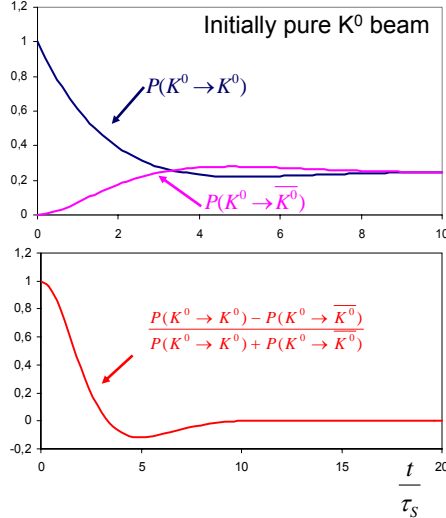
$$\Delta\Gamma = -11.182 \cdot 10^9 \text{ } \hbar\text{s}^{-1}$$

$$x = \frac{\Delta m}{\Gamma} = 0.942 \pm 0.007$$

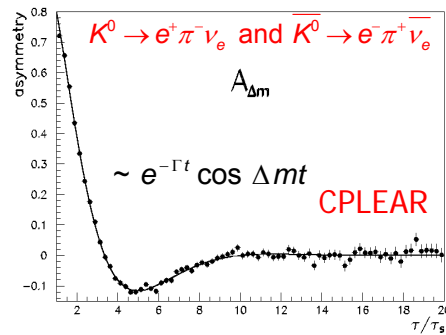
$$y = \frac{\Delta\Gamma}{2\Gamma} = -0.9966$$

Neutral kaon system

Expectation



Measurement



$K^0 - \bar{K}^0$ (strangeness) oscillation in the SM

Short range effects

Long range effects

Oscillation frequency Δm :

$$\Delta m \sim \frac{G_F^2}{4\pi} m_K f_K^2 \sum_{q=u,c,t} m_q^2 |V_{qs} V_{qd}|^2 \approx \frac{G_F^2}{4\pi} m_K f_K^2 m_c^2 |V_{cs} V_{cd}|^2$$

c quark contribution dominant: although m_t^2 is very large, the factor $|V_{ts} V_{td}|^2 \sim \lambda^5$ is very small !

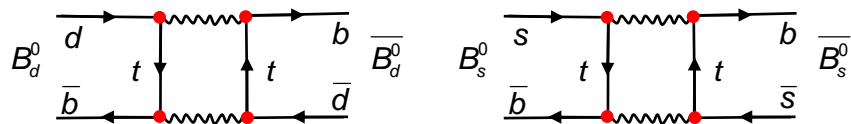
2.3 Neutral B Meson

Mixing mechanisms:

- **Mixing through decay:** many possible hadronic decays $\rightarrow \Gamma$ is large

$$\Rightarrow y = \frac{\Delta\Gamma}{2\Gamma} \text{ is small} = \begin{cases} \approx 0 \text{ for } B_d^0 \\ \approx O(0.1) \text{ for } B_s^0 \end{cases} \Rightarrow \text{don't expect mixing via decay}$$

- **Mixing through oscillation** \rightarrow Significant contribution only from top loop



$$\Delta m \sim m_t^2 |V_{tb} V_{td}|^2 \sim m_t^2 \cdot O(\lambda^6) \quad \Delta m \sim m_t^2 |V_{tb} V_{ts}|^2 \sim m_t^2 \cdot O(\lambda^4)$$

Large $\Delta m_{s,d}$: $\Delta m_s \sim 1/\lambda^2 \Delta m_d \rightarrow B_s$ osc. is about 35 times faster than B_d osc.

