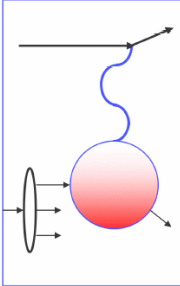


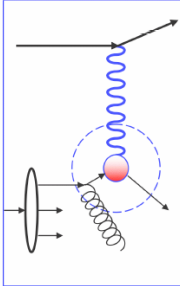
### QCD explains observed scaling violation

Large x: valence quarks



$Q^2 \uparrow \Rightarrow F_2 \downarrow$  for fixed x

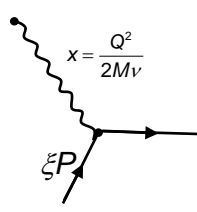
Small x: Gluons, sea quarks



$Q^2 \uparrow \Rightarrow F_2 \uparrow$  for fixed x

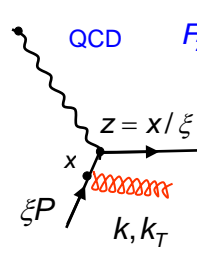
### Quantitative description of scaling violation

Quark Parton Model



$$F_2(x) = x \sum_i e_i^2 \int_0^1 q_i(\xi) \cdot \delta(x - \xi) d\xi = x \sum_i e_i^2 q_i(x)$$



QCD

$$F_2(x, Q^2) = x \sum_i e_i^2 \int_0^1 \frac{d\xi}{\xi} q_i(\xi) \cdot \left[ \delta\left(1 - \frac{x}{\xi}\right) + \frac{\alpha_s}{2\pi} P_{qq}\left(\frac{x}{\xi}\right) \log \frac{Q^2}{\mu_0^2} \right]$$

$$\hat{\sigma}(\gamma^* q \rightarrow qg) \sim \frac{\alpha_s}{2\pi} P_{qq}(z) \int_{\mu_0^2}^{Q^2} \frac{dk_T^2}{k_T^2}$$

$$\sim \frac{\alpha_s}{2\pi} P_{qq}(z) \log\left(\frac{Q^2}{\mu_0^2}\right)$$

$P_{qq}$  probability of a quark to emit gluon and becoming a quark with momentum reduced by fraction z.  
 $\mu_0$  cutoff parameter

Changing to the quark densities:

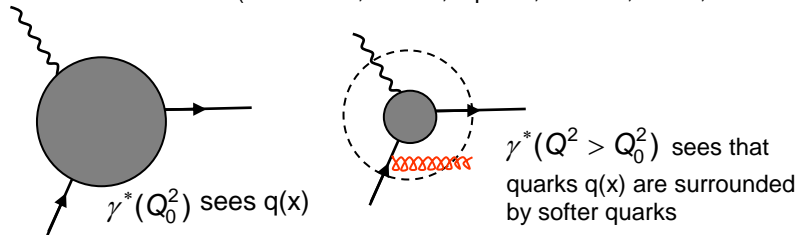
$$q_i(x, Q^2) = q_i(x) + \underbrace{\frac{\alpha_s}{2\pi} \log \frac{Q^2}{\mu_0^2} \int_0^1 \frac{d\xi}{\xi} q_i(\xi) P_{qq}\left(\frac{x}{\xi}\right)}_{\Delta q(x, Q^2)}$$

Integro-differential equation for  $q(x, Q^2)$ :

$$\frac{d}{d \log Q^2} q(x, Q^2) = \frac{\alpha_s}{2\pi} \int_0^1 \frac{d\xi}{\xi} q(\xi, Q^2) P_{qq}\left(\frac{x}{\xi}\right)$$

**DGLAP** evolution equation

(Dokshitzer, Gribov, Lipatov, Altarelli, Parisi, 1972 – 1977)



### Evolution of parton densities (quarks and gluons)

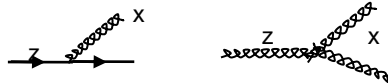
evolution of quark density with  $\ln Q^2$

$$\frac{\partial q(x, Q^2)}{\partial \ln Q^2} = \frac{\alpha_s}{2\pi} \int_x^1 \frac{dz}{z} \left[ q(z, Q^2) P_{qq}\left(\frac{x}{z}\right) + g(z, Q^2) P_{gq}\left(\frac{x}{z}\right) \right]$$



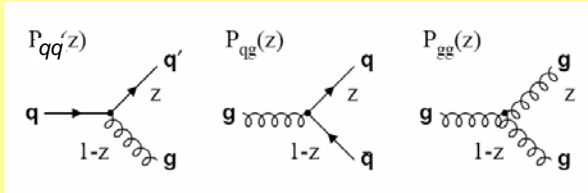
$$\frac{\partial g(x, Q^2)}{\partial \ln Q^2} = \frac{\alpha_s}{2\pi} \int_x^1 \frac{dz}{z} \left[ q(z, Q^2) P_{gq}\left(\frac{x}{z}\right) + g(z, Q^2) P_{gg}\left(\frac{x}{z}\right) \right]$$

evolution of gluon density with  $\ln Q^2$



**Splitting functions:** Probability that a parton (quark or gluon) emits a parton (q, g) with momentum fraction  $\epsilon=x/z$  of the parent parton.

Splitting functions are calculated as power series in  $\alpha_s$  up to a given order:



$$P_{ij}(z, \alpha_s) = P_{ij}^0(z) + \frac{\alpha_s}{2\pi} P_{ij}^1(z) + \dots$$

In leading order:  $P_{ij}(z, \alpha_s) \equiv P_{ij}^0(z)$

$$P_{qq}(z) = \frac{4}{3} \frac{1+z^2}{1-z}$$

$$P_{gq}(z) = \frac{4}{3} \frac{1+(1-z)^2}{z}$$

$$P_{qg}(z) = \frac{z^2 + (1-z)^2}{2}$$

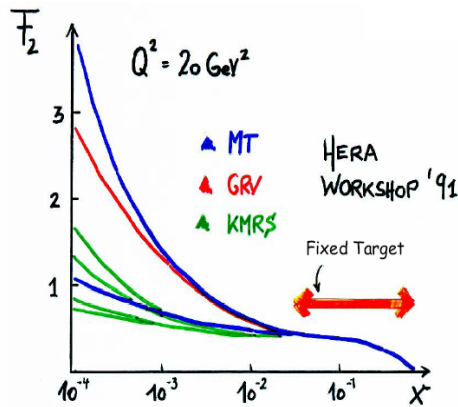
$$P_{gg}(z) = 6 \left( \frac{z}{1-z} + \frac{1-z}{z} + z(1-z) \right)$$

DGLAP Evolution (“symbolic”):

$$\frac{\partial}{\partial \log Q^2} \begin{bmatrix} q(x, Q^2) \\ g(x, Q^2) \end{bmatrix} = \frac{\alpha_s}{2\pi} \begin{bmatrix} P_{qq} \left[ \frac{x}{z} \right] & P_{qg} \left[ \frac{x}{z} \right] \\ P_{gq} \left[ \frac{x}{z} \right] & P_{gg} \left[ \frac{x}{z} \right] \end{bmatrix} \otimes \begin{bmatrix} q(x, Q^2) \\ g(x, Q^2) \end{bmatrix}$$

$$P \otimes f(x, Q^2) = \int_x^1 \frac{dz}{z} P\left(\frac{x}{z}\right) f(z, Q^2)$$

Bjorken x dependence of parton densities:



DGLAP:

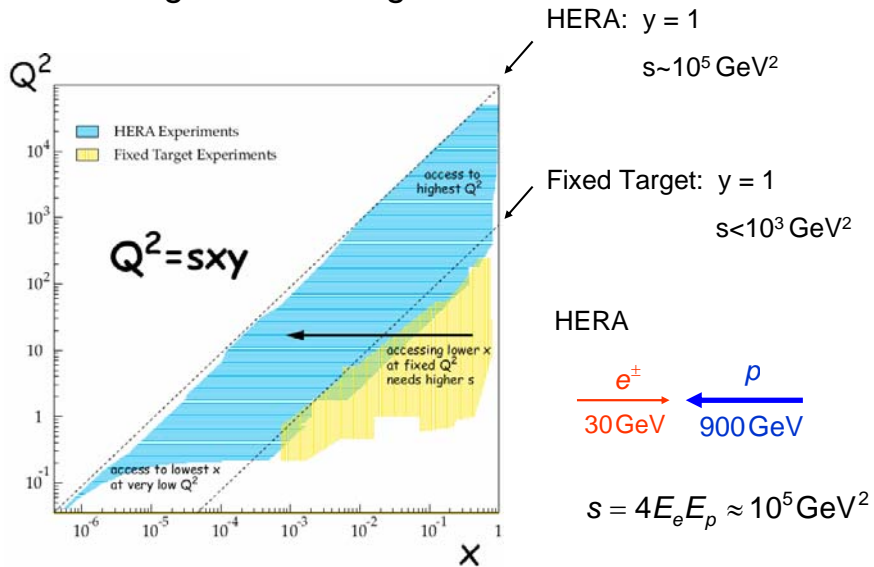
No prediction for the x dependence of the parton densities.

Status in 1991 (pre HERA): Data limited to a small x region. Models to extrapolate to smaller x differ significantly.

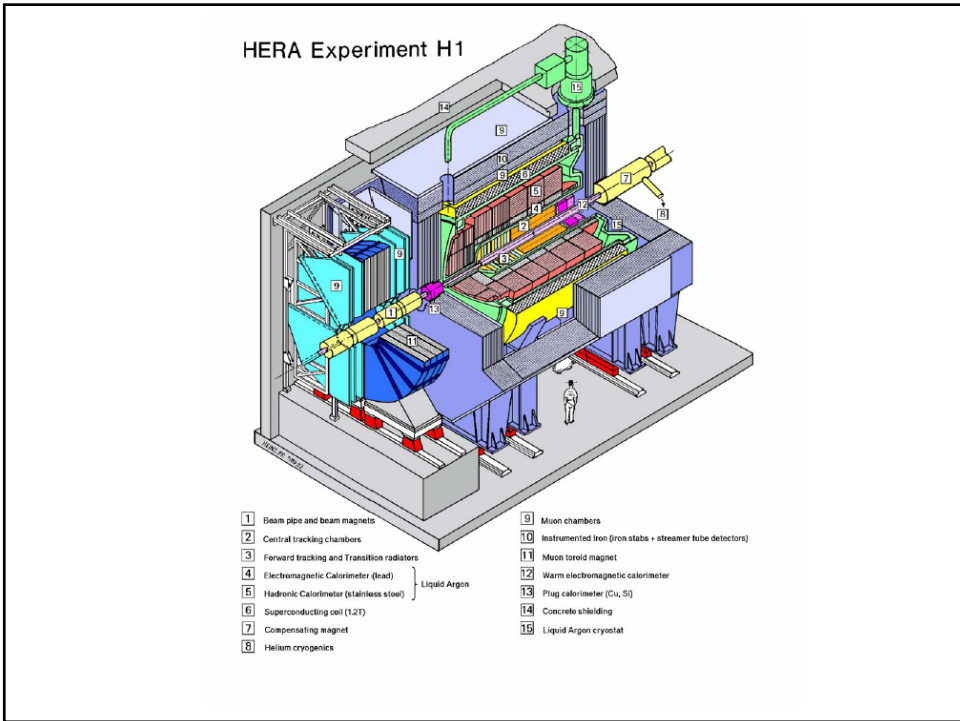
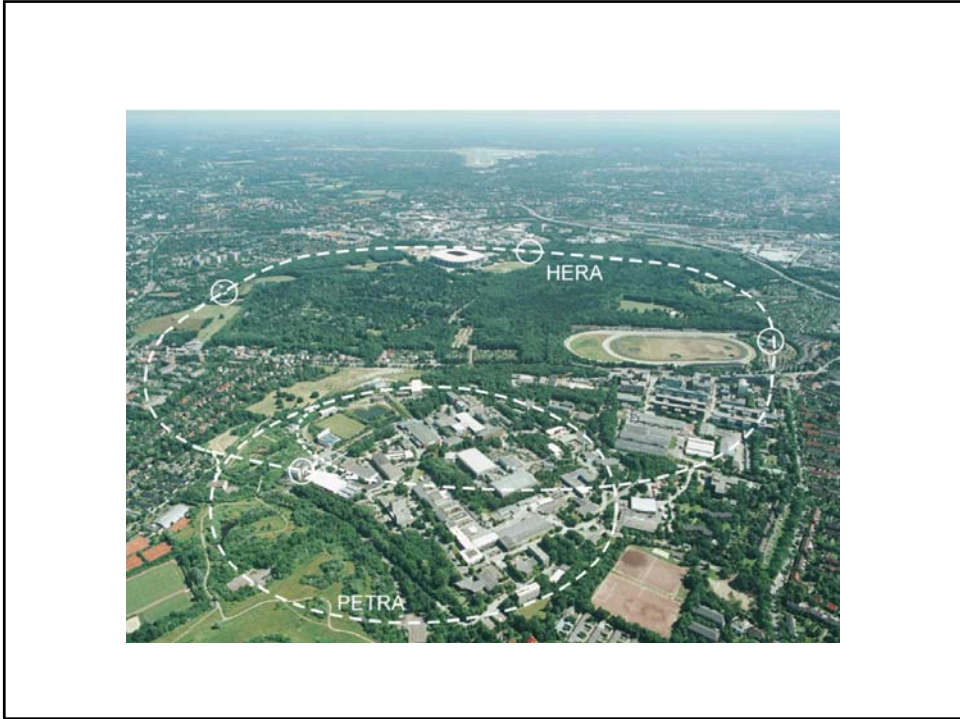


Measure structure functions (parton densities) at low x.

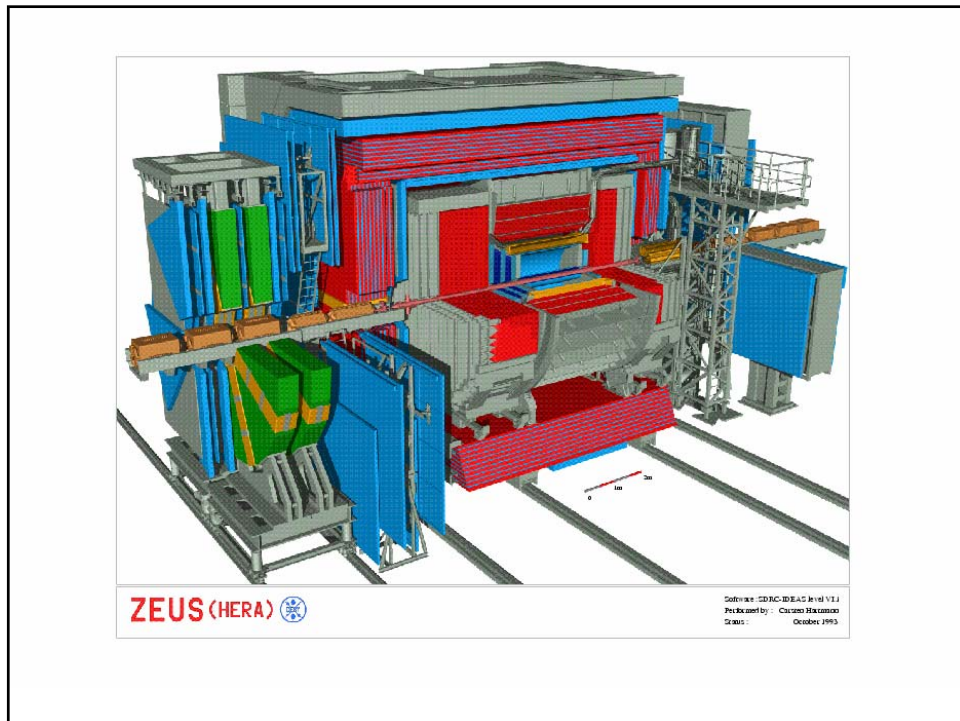
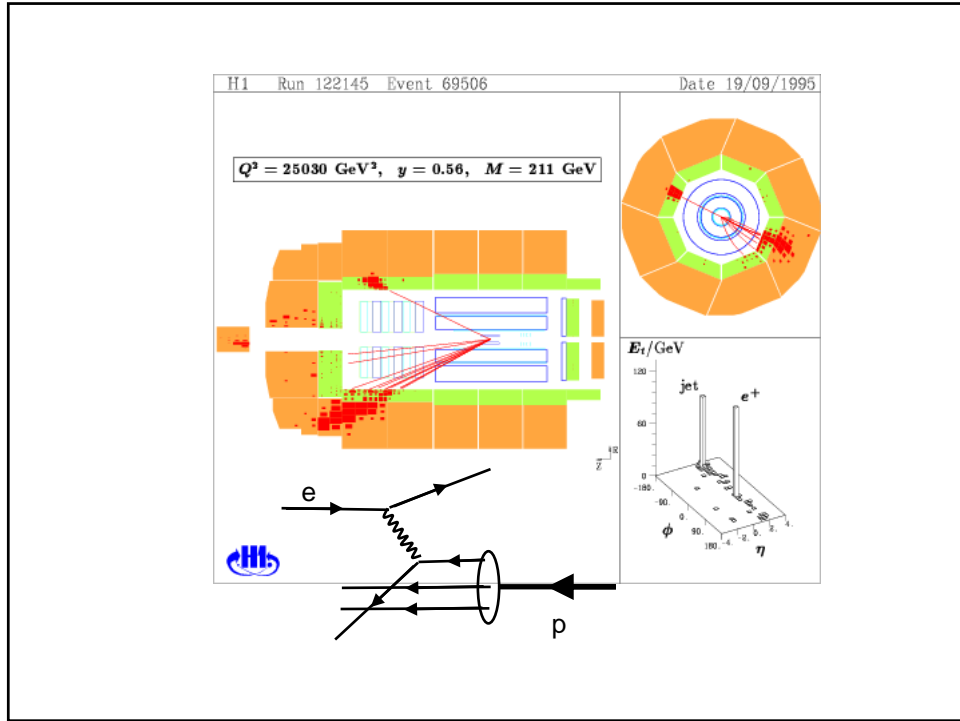
Accessing the low x region:

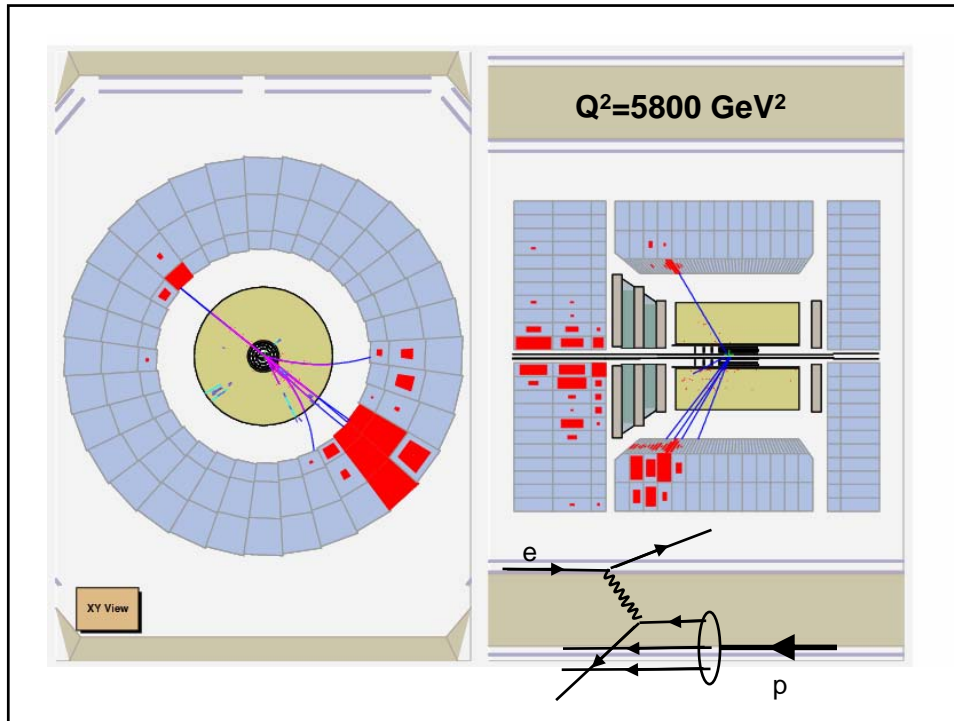


# Advanced Particle Physics: V. Experimental studies of QCD



# Advanced Particle Physics: V. Experimental studies of QCD





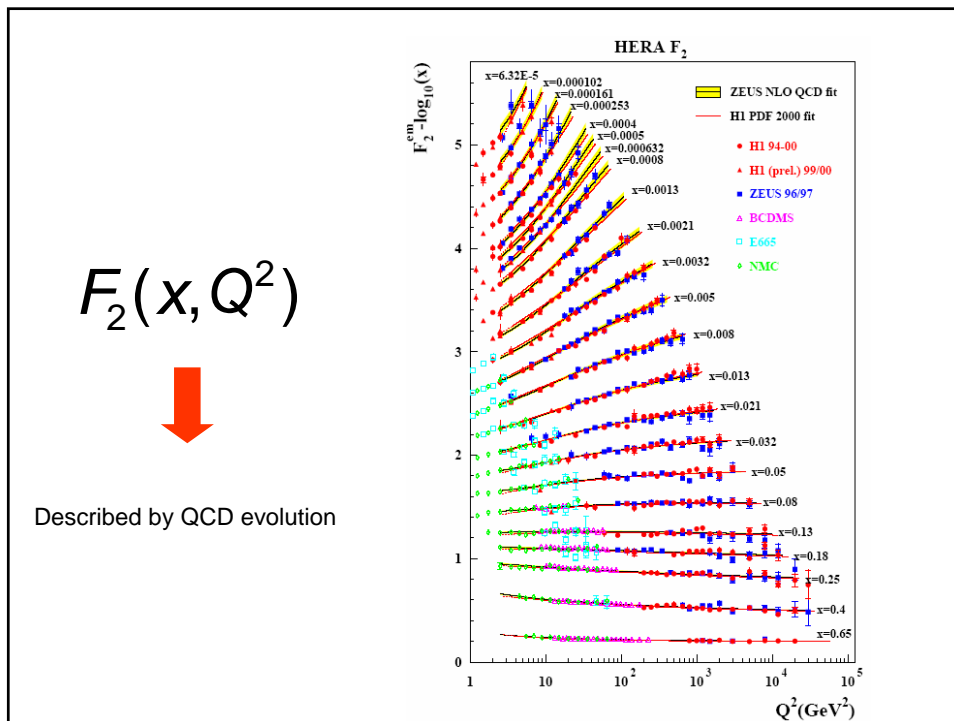
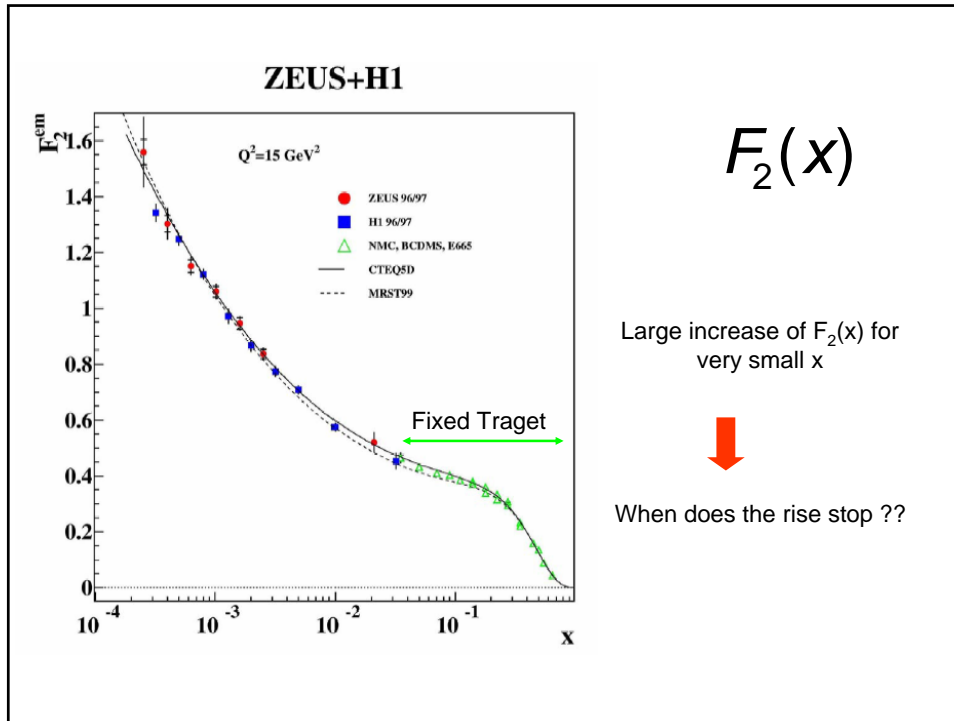
Measurement of the parton densities /  $F_2$

$$\frac{d^2\sigma}{dx dQ^2} = \left( \frac{2\pi\alpha^2}{xQ^4} \right) \cdot (2 \cdot (1-y)F_2(x, Q^2) + y^2F_2(x, Q^2))$$

↓ for  $y=1$

$$\frac{d^2\sigma}{dx dQ^2} = \left( \frac{2\pi\alpha^2}{xQ^4} \right) \cdot F_2(x, Q^2)$$

$$F_2(x, Q^2) = x \sum_q e_q^2 [q(x, Q^2) + \bar{q}(x, Q^2)]$$





### Experimental determination of the gluon density

Using the DGLAP evolution eq. one finds for  $F_2(x, Q^2)$ :

$$\frac{dF_2(x, Q^2)}{d \ln Q^2} = x \sum_i e_i^2 \frac{\alpha_s(Q^2)}{2\pi} \cdot \int_x^1 \frac{dz}{z} \left[ P_{qq}\left(\frac{x}{z}\right) q_i(z, Q^2) + P_{qg}\left(\frac{x}{z}\right) g(z, Q^2) \right]$$

For small  $x$  ( $x < 10^{-2}$ ):

quark pair production through gluon splitting dominant ( $1/x$  gluon spectrum):

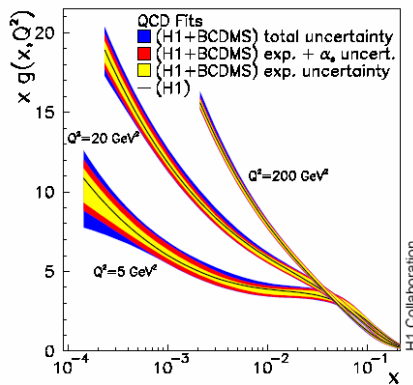
$$\rightarrow P_{qg}\left(\frac{x}{z}\right) g(z, Q^2) \text{ dominant}$$

As an approximation one finds:

$$x \cdot g(x, Q^2) \approx \frac{27\pi}{10\alpha_s(Q^2)} \cdot \frac{dF_2(x, Q^2)}{d \ln Q^2}$$

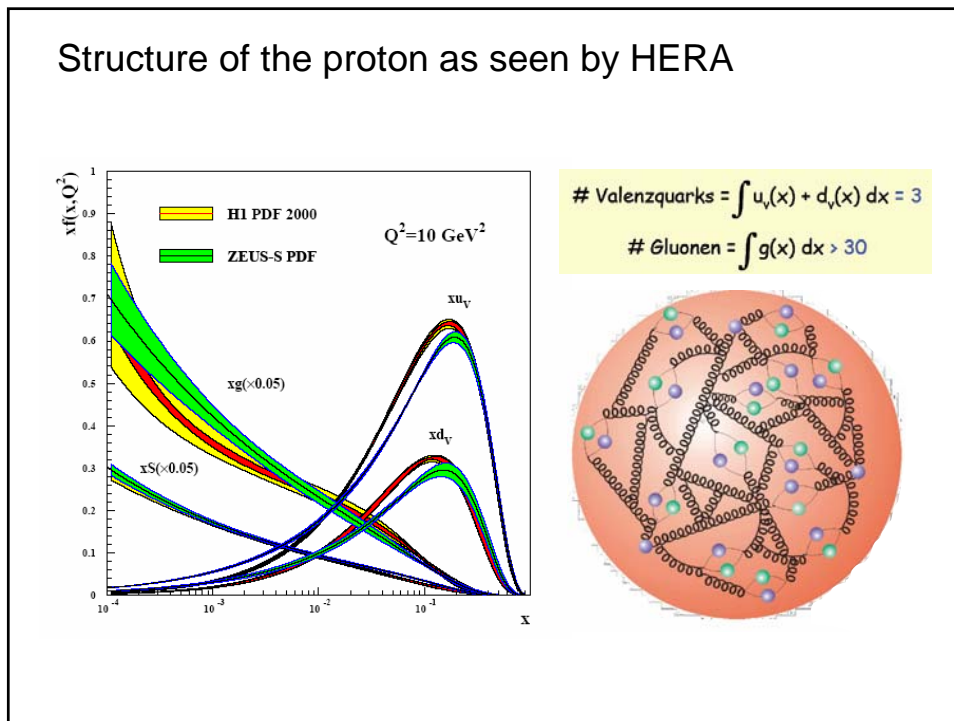
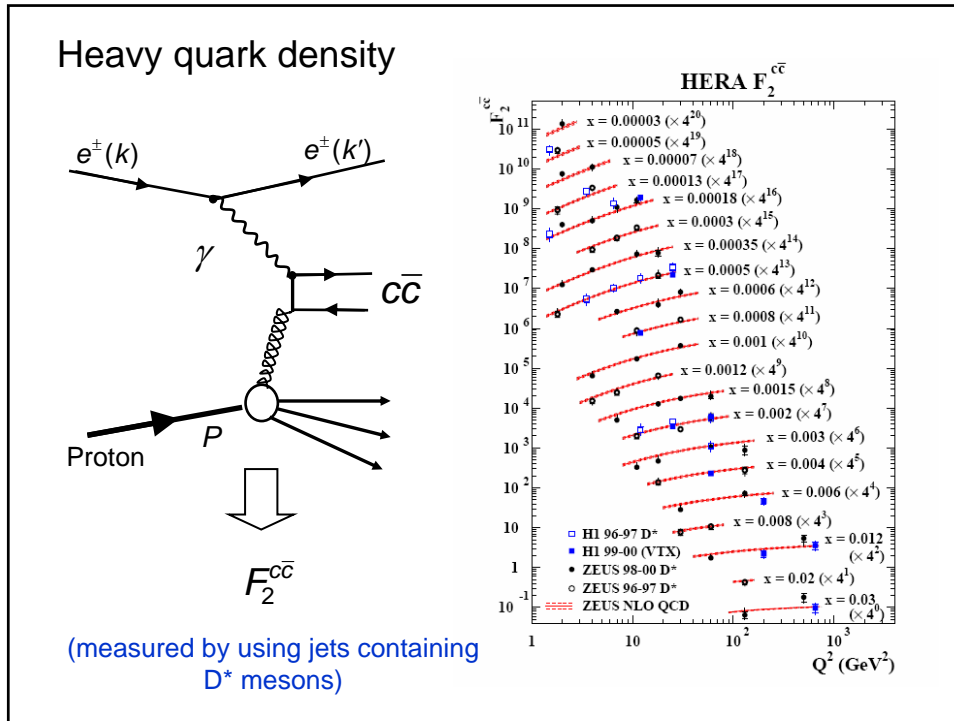
i.e. scaling violation of  $F_2$  at small  $x$  measures the gluon density.

### Gluon density $g(x, Q^2)$ :

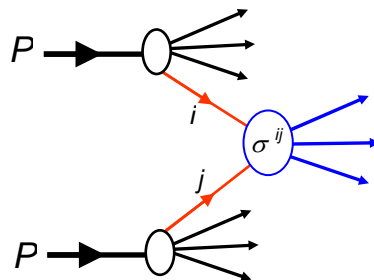


In practice one makes a global DGLAP fit to the measured  $F_2(x, Q^2)$

$$\alpha_s(M_Z^2) = 0.115$$



#### 4. Factorization – QCD at LHC



Generalized cross section for n parton production can be **factorized**:

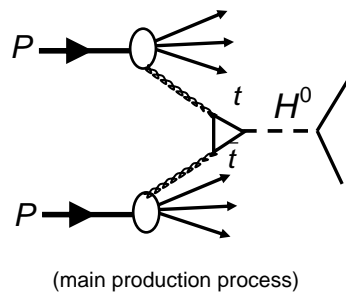
- PDFs  $f_i(x, \mu^2)$  and  $f_j(x, \mu^2)$
- “short distance” cross section,  $\hat{\sigma}^{ij \rightarrow k_1 \dots k_n}$ , perturbatively calculable

$$\sigma^n = \sum_{i,j,k_1,\dots,k_n=q,\bar{q},g} \int dx_1 dx_2 f_i(x_1, \mu^2) f_j(x_2, \mu^2) \hat{\sigma}^{ij \rightarrow k_1 \dots k_n}$$

Factorization scale is the value  $\mu$  at which the structure functions are defined and at which the distinction between short//long range parts is done.

Parton densities important to predict signal and background at LHC:

Higgs-production at LHC



Can we extrapolate from HERA?

