

Lecture:

# Standard Model of Particle Physics

Heidelberg SS 2016

## Discovery of the Higgs

# Contents

- Intro: Statistics and limit setting
- Higgs Mass Predictions
- Higgs Searches at LEP
- Higgs Searches at Hadron Colliders
- WW-Scattering Amplitude

# Statistics and Limit Setting

chi<sup>2</sup> fit:

$$\chi^2 = \sum_i \frac{(y_i - \mu_i)^2}{\sigma_i^2}$$

$y_i$  measurement

$\mu_i$  model prediction (nuisance parameter)

$\sigma_i$  uncertainty (statistical and systematical)

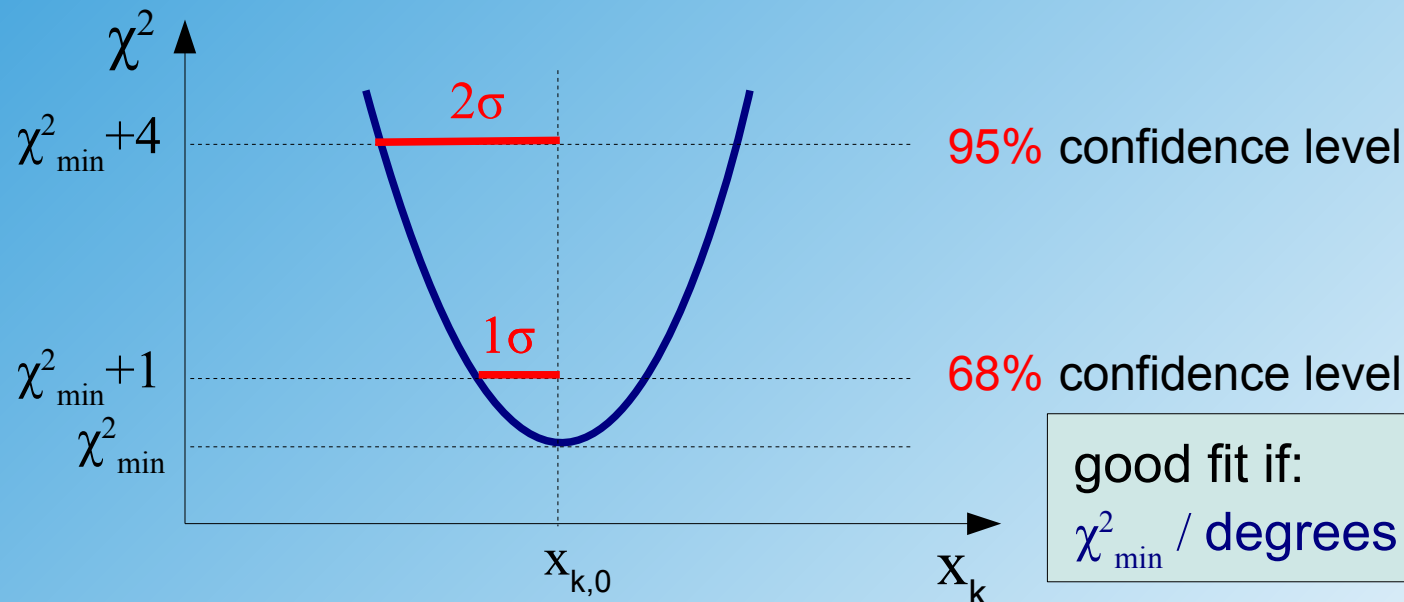
chi<sup>2</sup> fit with correlated errors:

$$\chi^2 = \sum_i \sum_j (y_i - \mu_i) \text{cov}_{ij}^{-1} (y_j - \mu_j) \quad \text{cov}_{ij} \text{ covariance (error) matrix}$$

Parameter Fit:

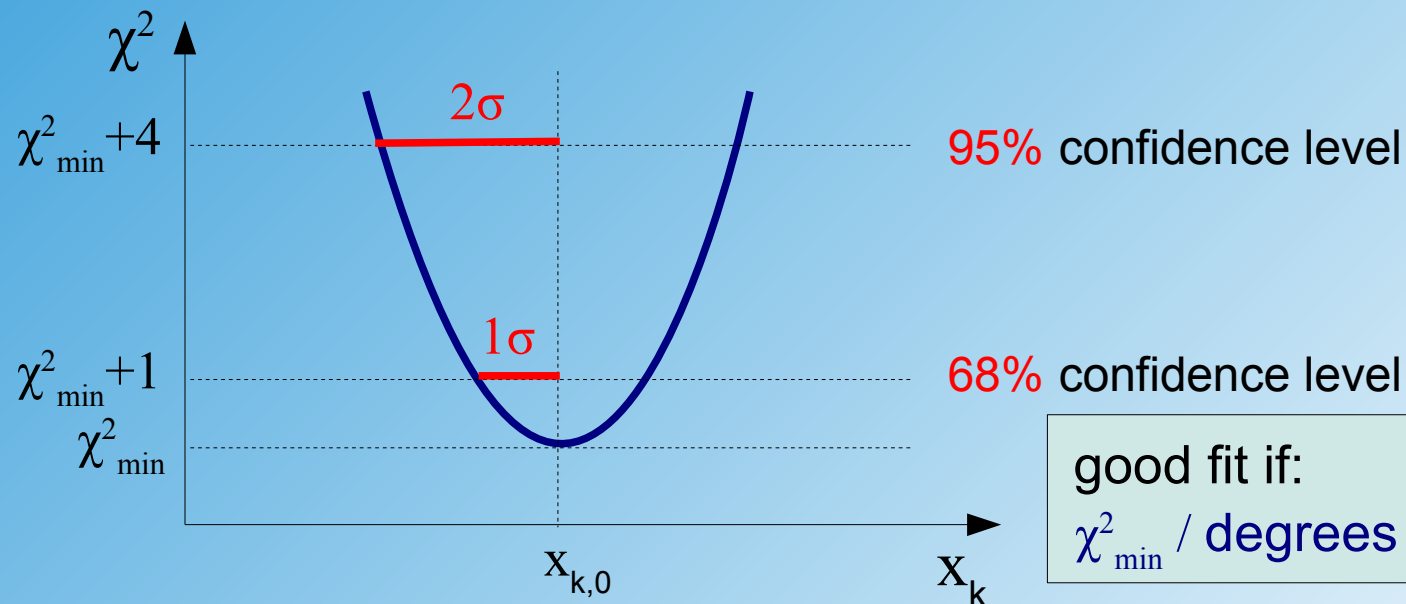
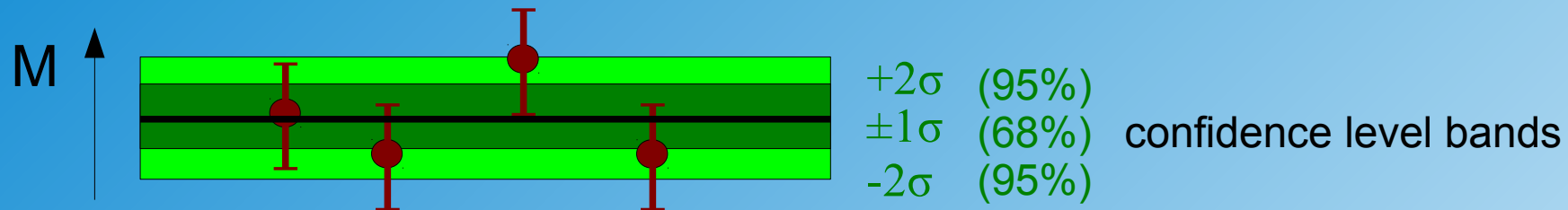
$$\mu_i = \mu_i(x_1, x_2, \dots, x_n)$$

$x_k$  model parameter, e.g. cross section



# Example Fit

Measurement of some mass (1-Parameter fit) from 4 experiments:



good fit if:

$\chi^2_{\min} / \text{degrees of freedom} \sim 1$

# Probability Densities

Gaussian (normal) distribution:

$$f(x) = \frac{1}{\sqrt{2\pi}\sigma} e^{-\frac{(x-\mu)^2}{2\sigma^2}} \rightarrow P(a < x < b) = \int_a^b f(x) dx$$

→ used for systematic uncertainties (symmetric)

Poissonian distribution:

$$P(N) = \frac{e^{-\mu} \mu^N}{N!}$$

→ used for statistical uncertainties

Note, Poisson distribution approaches Gaussian distribution for large  $\mu$

# Limit Setting Philosophies

## Bayesian Method:

- based on the **experiment posterior**, exclusion limits are calculated
- low probability models are excluded
- probabilities are assigned to models using a prior

“natural method” but choice of prior is arbitrary

## Frequentist Method:

- based on **Monte Carlo** toy experiments probabilities are assigned to all possible experimental outcomes
- exclusion limit is set by that model which excludes this experimental outcome with certain confidence interval

computationally expensive and might give unphysical results  
(e.g. negative cross sections)

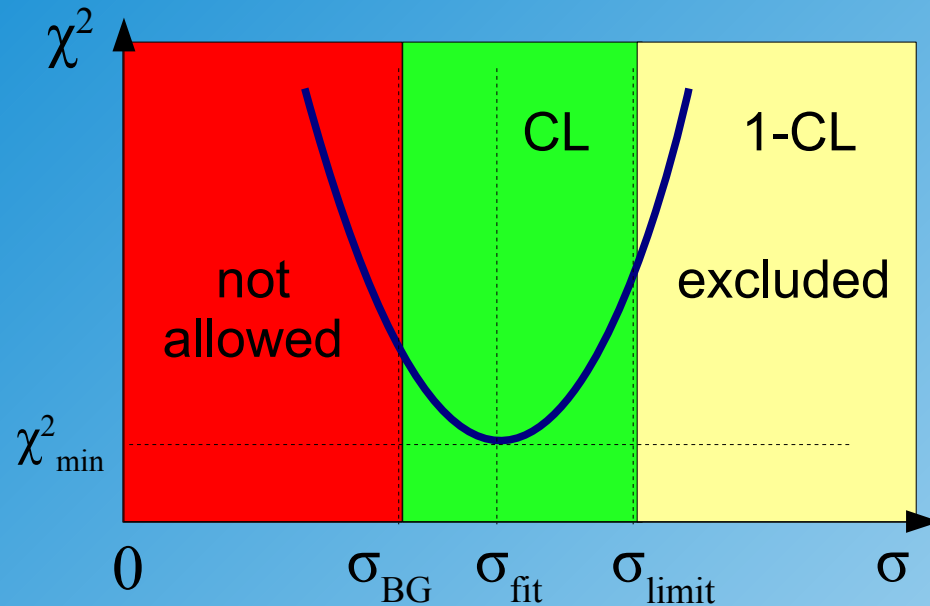
# Bayesian Method

Model:  $N = N_{BG} + N_{Sig}$

background + signal

$$\sigma = \sigma_{BG} + \sigma_{Sig} = N/L$$

choose cross section as “prior”



Probability:

$$P \propto e^{-\frac{1}{2}(\chi^2(\sigma) - \chi_{fit}^2)}$$

usually additional constraint:

$$\sigma > \sigma_{BG} \quad \text{because} \quad \sigma_{Sig} > 0$$

$$CL = \frac{\int_{\sigma > \sigma_{BG}}^{\sigma_{CL}} P(\sigma) d\sigma}{\int_{\sigma > \sigma_{BG}}^{\infty} P(\sigma) d\sigma}$$

CL = confidence level



# Choice of Prior

- Cross sections depend on couplings
- Alternatively we can choose coupling as prior

$$\sigma_{Sig} = \sigma_0 \alpha^2$$

$$\frac{d\sigma}{d\alpha} = \frac{d\sigma_{Sig}}{d\alpha} = 2\sigma_0 \alpha$$

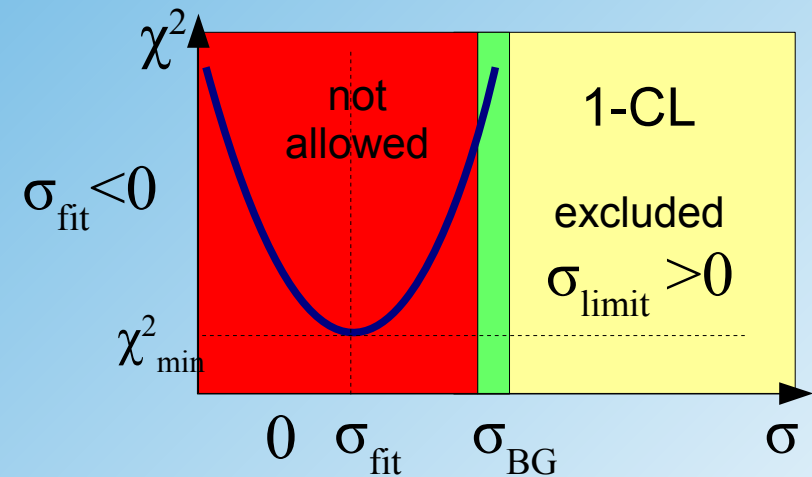
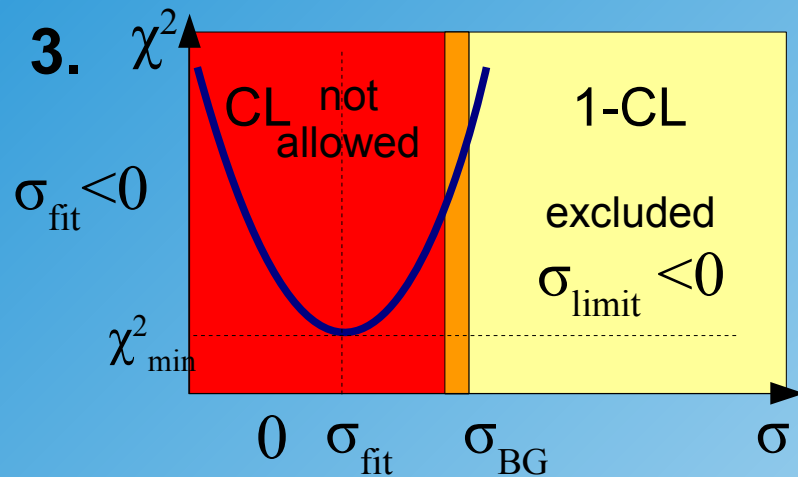
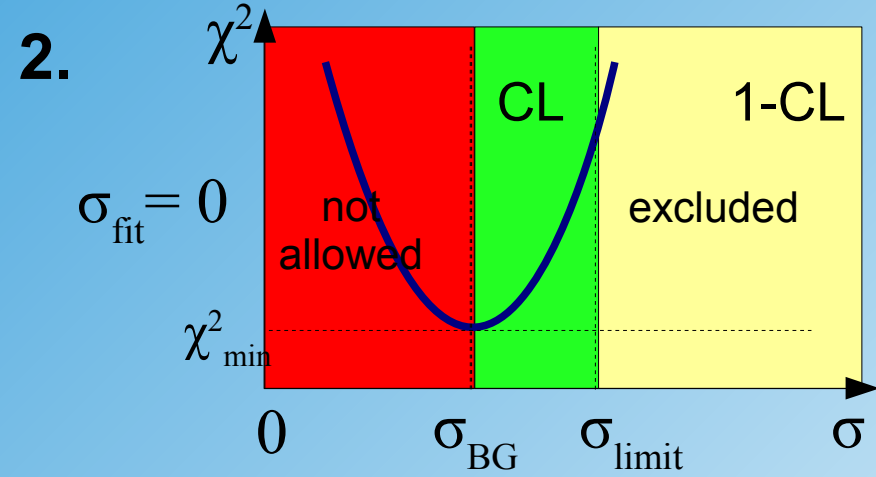
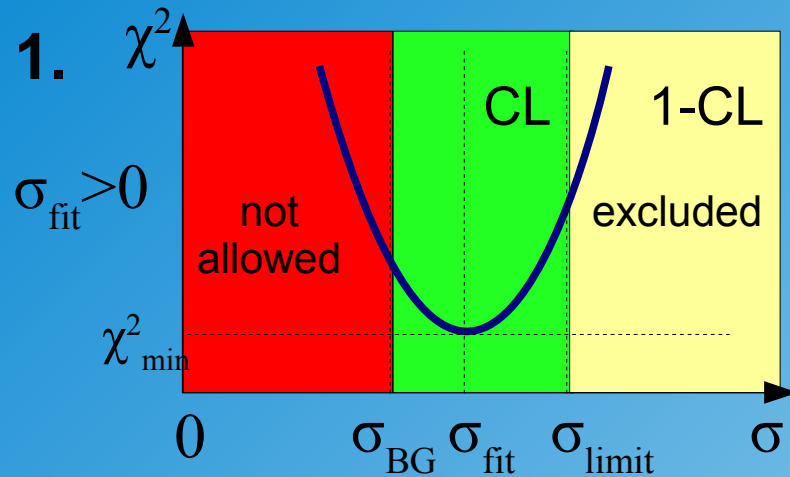
$$CL = \frac{\int_{\sigma_{Sig}>0}^{\sigma_{CL}} P(\sigma) d\sigma}{\int_{\sigma_{Sig}>0}^{\infty} P(\sigma) d\sigma} = \frac{\int_{\alpha>0}^{\sigma_{CL}} P(\alpha^2) \alpha d\alpha}{\int_{\alpha>0}^{\infty} P(\alpha^2) \alpha d\alpha} \neq \frac{\int_{\alpha>0}^{\alpha_{CL}} P(\alpha) d\alpha}{\int_{\alpha>0}^{\infty} P(\alpha) d\alpha}$$

Result depends on choice of prior!



# Negative Cross Section Limit Problem

possible outcomes::



$$CL = \frac{\int_{\sigma > -\infty}^{\sigma_{\text{CL}}} P(\sigma) d\sigma}{\int_{\sigma > -\infty}^{\infty} P(\sigma) d\sigma} \quad \text{WRONG!}$$

negative cross sections  $\rightarrow$  unphysical

$$CL = \frac{\int_{\sigma_{\text{Sig}} > 0}^{\sigma_{\text{CL}}} P(\sigma) d\sigma}{\int_{\sigma_{\text{Sig}} > 0}^{\infty} P(\sigma) d\sigma}$$

only positive cross sections!

# Frequentist Method

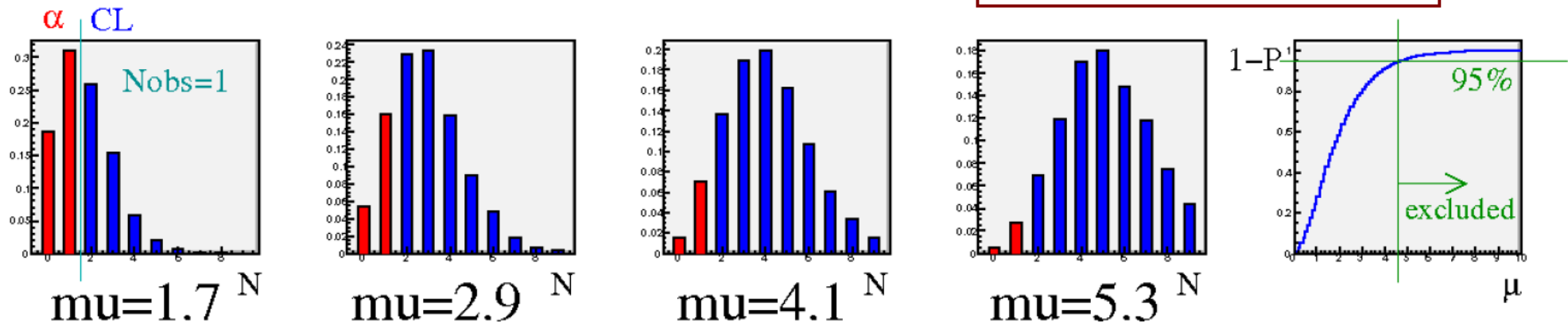
S.Schmitt

- Frequentist limit: exclude all theories which produce the data at small probability  $\alpha$  less than  $1 - \text{CL}$  (typically:  $\text{CL} = 0.95$ )

$$\alpha = P_{\mu}(N \leq N_{\text{obs}}) < 1 - \text{CL}$$

$\alpha$ : also called p-value

Frequentist limit:  
sum (integrate)  
over observations  
up to  $N_{\text{obs}}$   
Repeat for each model



- $\mu$  = expectation value (BG+signal)
- set limit with 95% confidence level for  $\mu = 4.6$
- experiment has a 5% probability to happen
- measurement uncertainty does not enter (in contrast to Bayesian method)

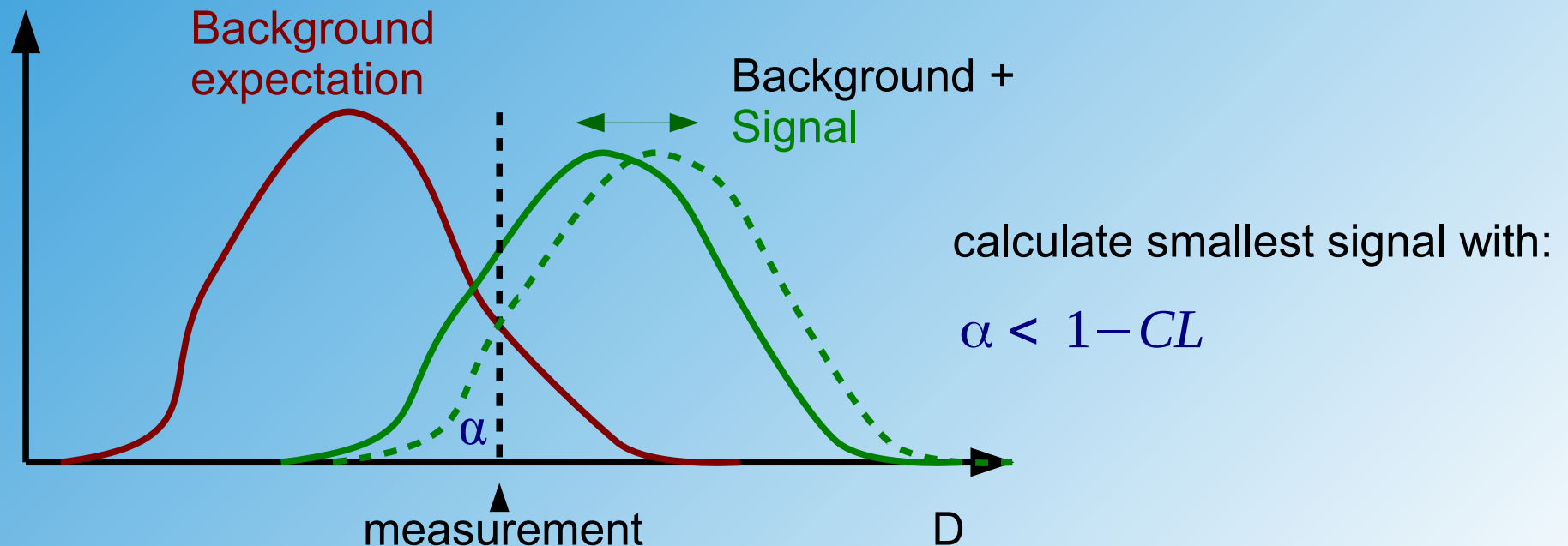
# Frequentist Method

In case of many observables  $x_k$  a combined discriminator variable is often defined:

$$D = D(x_1, x_2, x_3, \dots, X_k)$$

- large discriminator means high probability
- small discriminator means low probability

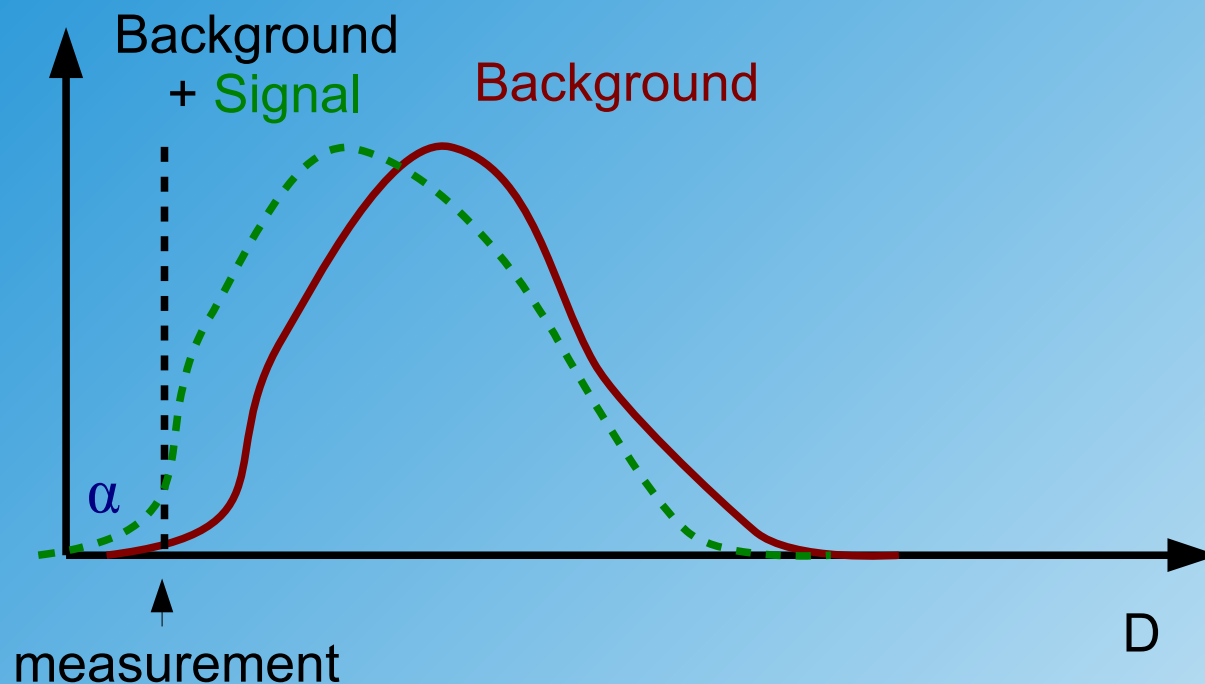
Often, the output from artificial neural nets or other multivariate methods is used as discriminator variable



# Negative Limit Problem with Frequentist Method

Problem in case of a very small measurement value with  $P(\text{BG}) < (1-\text{CL})$

- would require a negative signal cross section:



negative signal  $\rightarrow$  unphysical solution!

# CL<sub>s</sub> Method

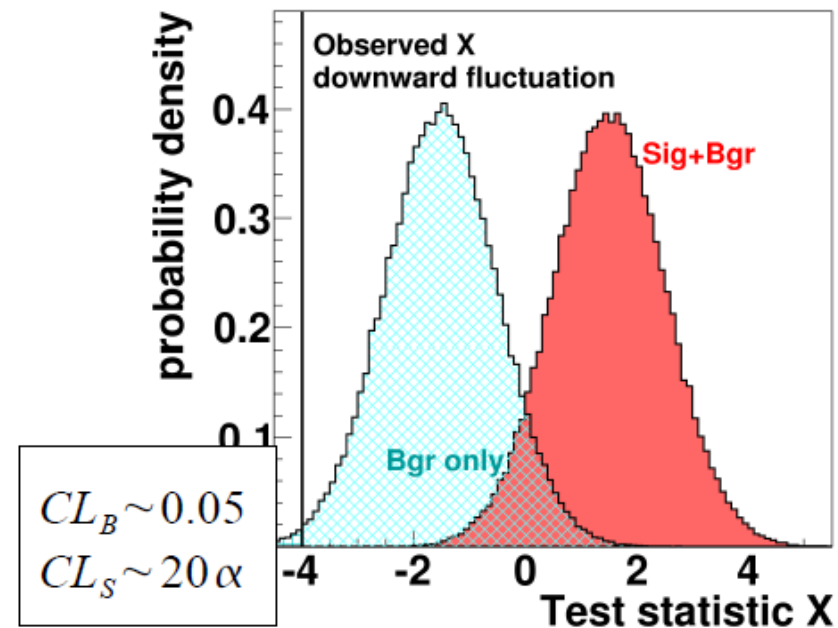
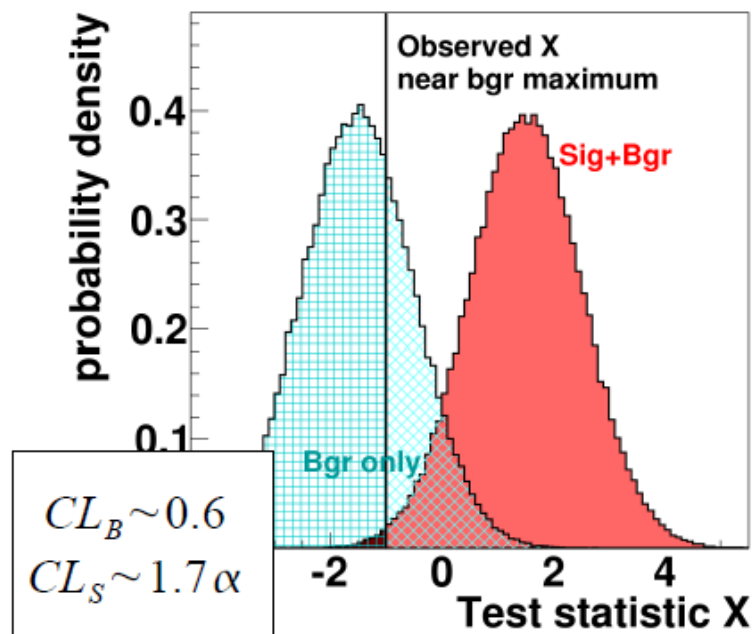
- Use ratio of two probabilities CL<sub>s</sub> instead of  $\alpha$  to test against CL

$$CL_{SB} = \alpha = \int_{X < X_{\text{obs}}} P(X | \text{signal} + \text{bgr}) dX$$

$$CL_B = \int_{X < X_{\text{obs}}} P(X | \text{bgr}) dX$$

$$CL_s = \frac{CL_{SB}}{CL_B}$$

- Standard model has CL<sub>s</sub>=1 and is never excluded



S.Schmitt

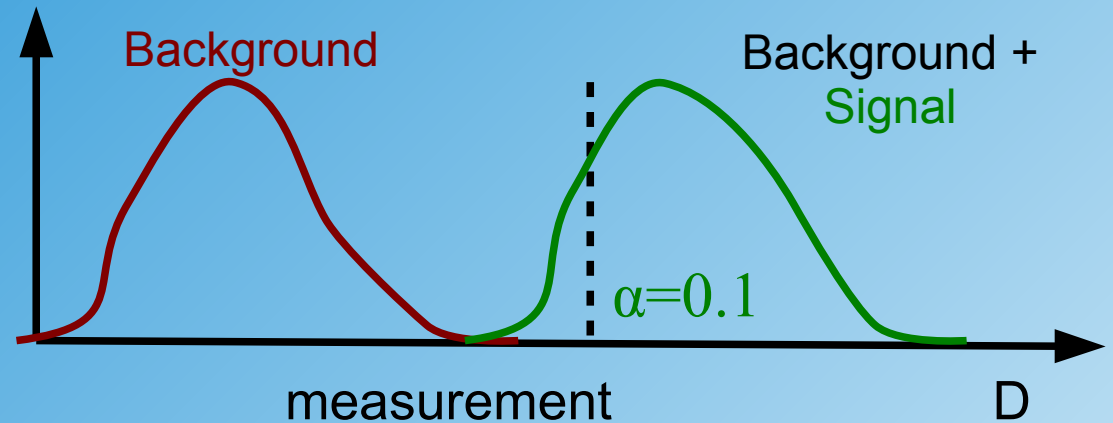
CL<sub>s</sub> > CL<sub>SB</sub> by definition!

# Another Example $CL_S$ Method

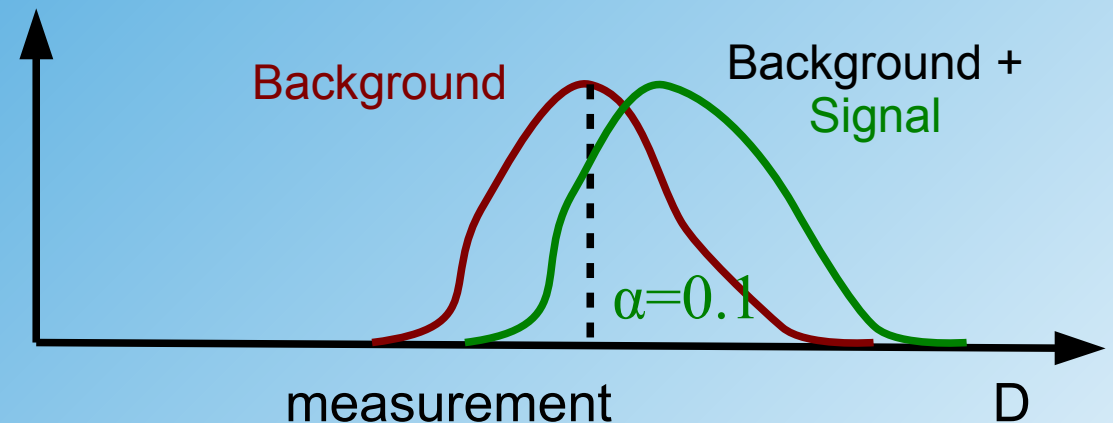
Definition:

$$CL_S = \frac{CL_{SB}}{CL_B}$$

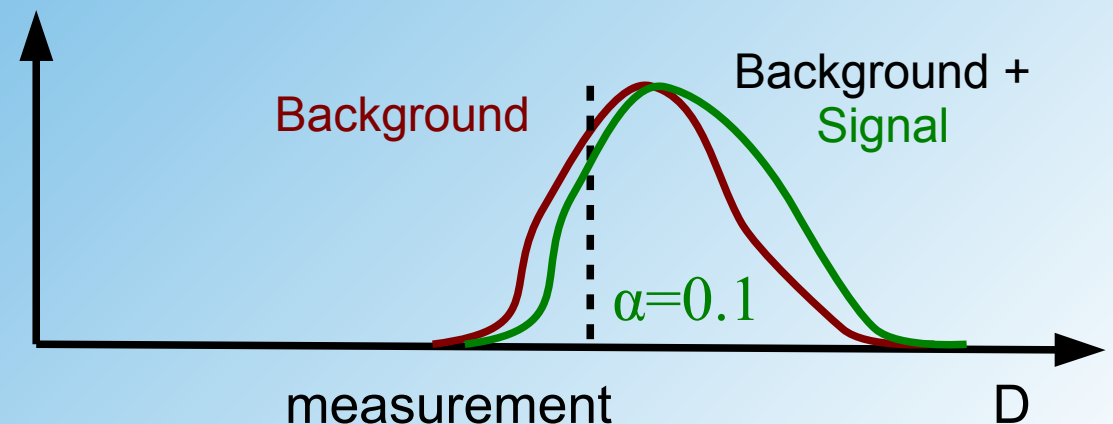
$$0.1 = \frac{0.1}{1.0}$$



$$0.2 = \frac{0.1}{0.5}$$



$$0.5 = \frac{0.1}{0.2}$$



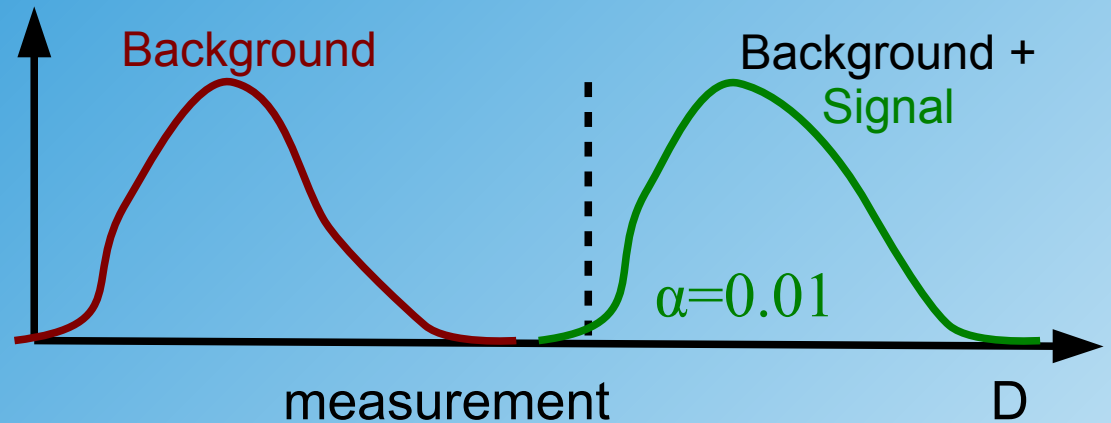


# Another Example $CL_S$ Method

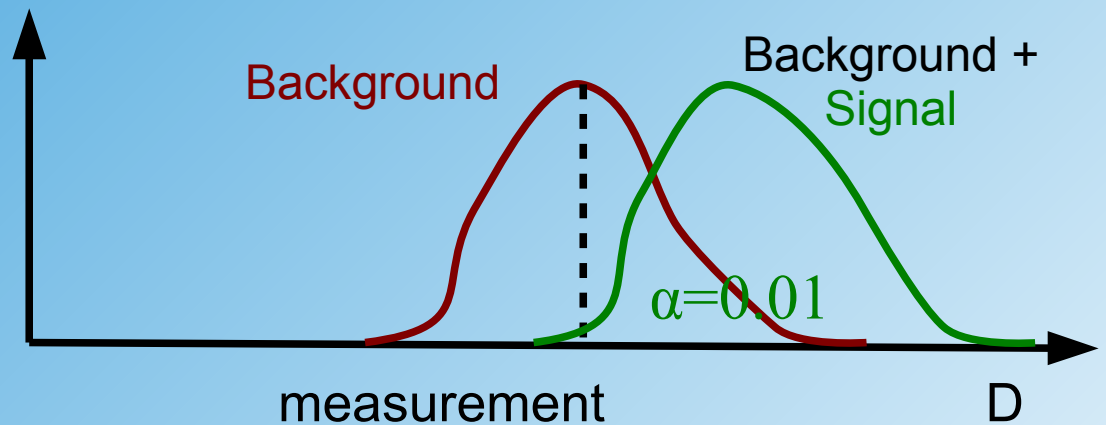
Definition:

$$CL_S = \frac{CL_{SB}}{CL_B}$$

$$0.01 = \frac{0.01}{1.0}$$

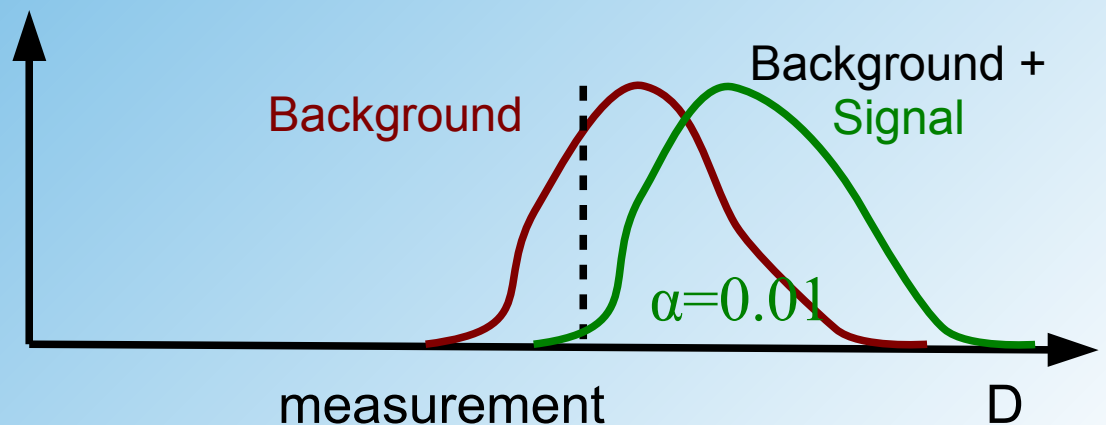


$$0.02 = \frac{0.01}{0.5}$$



$$0.05 = \frac{0.01}{0.2}$$

excluded at 95% CL





# Reminder: Indirect Higgs Mass Constraints

# Radiative Corrections and Indirect Higgs Constraints

$$\sin^2 \theta_w = 1 - \frac{M_W^2}{M_Z^2} \quad \sin \theta_w = \frac{e}{g}$$

$$\rho = \frac{m_W^2}{m_Z^2 \cos^2 \theta_W} = 1$$

$$\sin^2 \theta_W = 1 - \frac{m_W^2}{m_Z^2}$$

$$m_W^2 = \frac{\pi \alpha}{\sqrt{2} \sin^2 \theta_W G_F}$$

$\alpha(0)$

Lowest order  
SM predictions

$\Rightarrow$

$\Rightarrow$

$\Rightarrow$

$\Rightarrow$

$$\bar{\rho} = 1 + \Delta\rho$$

$$\sin^2 \theta_{\text{eff}} = (1 + \Delta\kappa) \sin^2 \theta_W$$

$$m_W^2 = \frac{\pi \alpha}{\sqrt{2} \sin^2 \theta_W G_F} (1 + \Delta r)$$

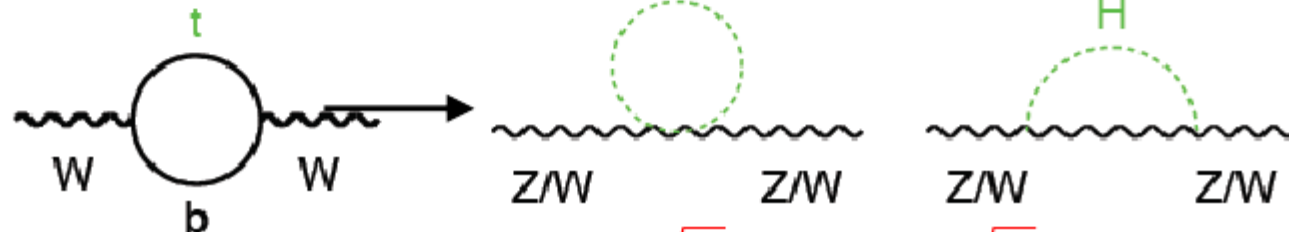
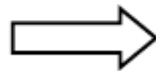
$$\alpha(m_Z^2) = \frac{\alpha(0)}{1 - \Delta\alpha}$$

with :  $\Delta\alpha = \Delta\alpha_{\text{lept}} + \Delta\alpha_{\text{top}} + \Delta\alpha_{\text{had}}^{(5)}$

Including radiative  
corrections

$$\sin^2 \theta_w$$

$$g_A, g_V$$



$$\Delta\rho, \Delta\kappa, \Delta r = f(m_t^2, \log(m_H), \dots)$$

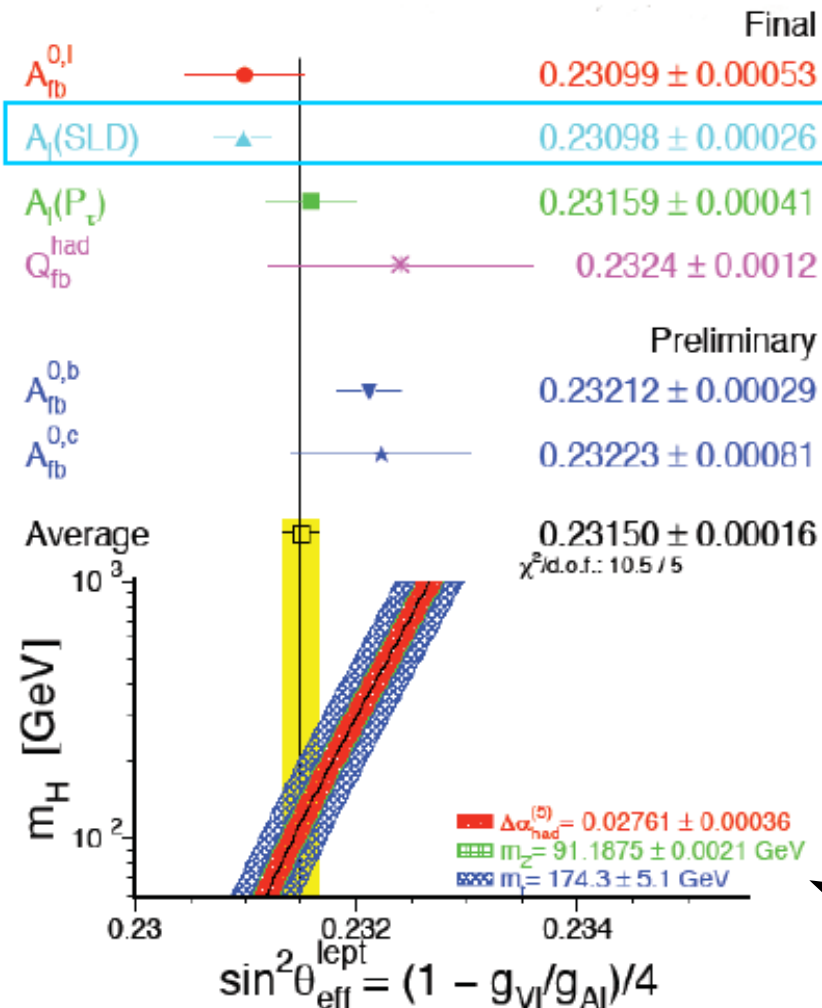
$$\sin^2 \theta_{\text{eff}}$$

$$\bar{g}_A, \bar{g}_V$$

$$\bar{g}_A = \sqrt{\bar{\rho}} T^3 \quad \bar{g}_V = \sqrt{\bar{\rho}} (T^3 - 2Q \sin^2 \theta_{\text{eff}})$$

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# Indirect Higgs Mass Prediction



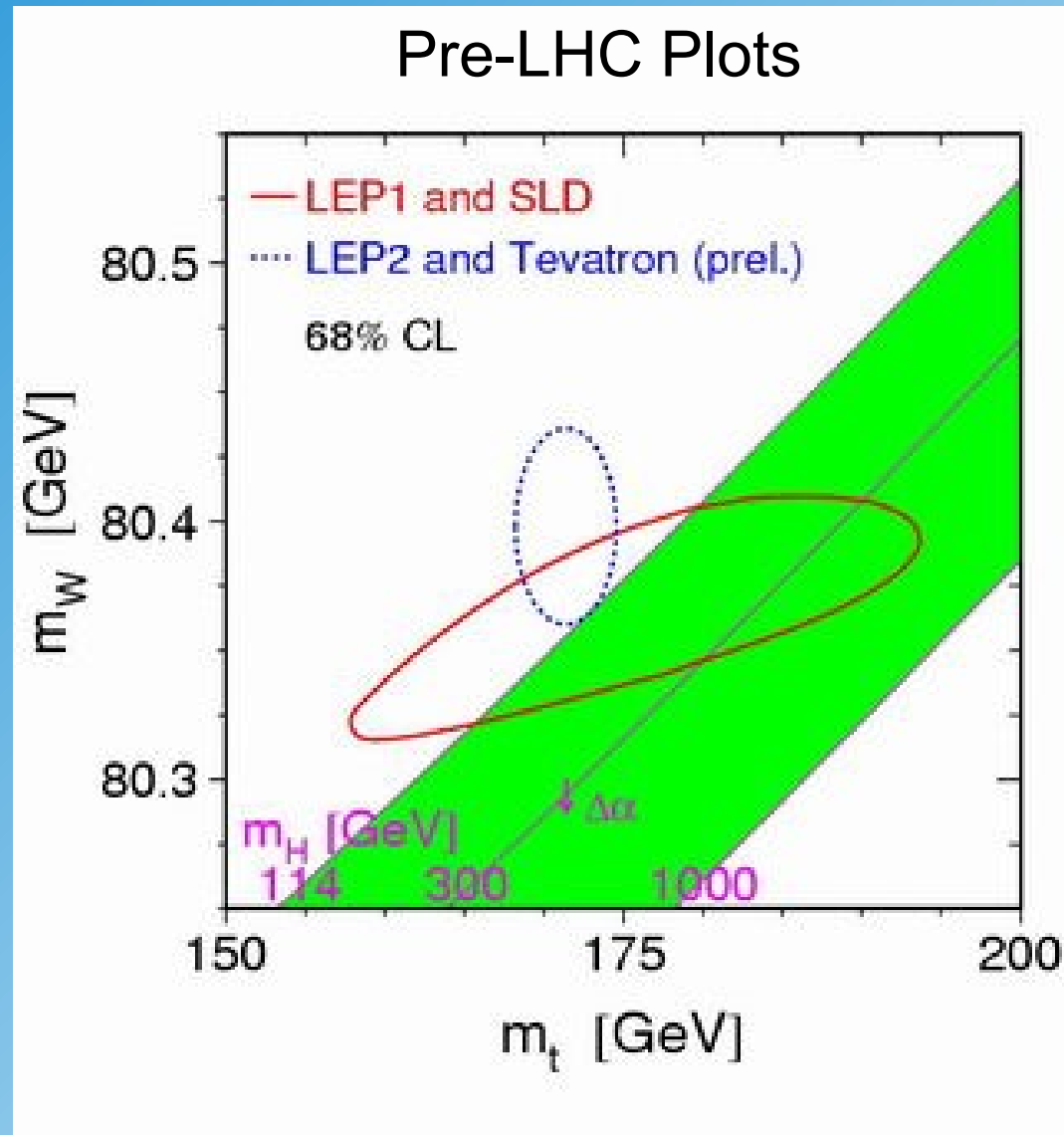
Take the top mass from direct measurements and use the radiative corrections to determine the Higgs mass.

$$\Delta r(m_t, M_H) = -\frac{3\alpha \cos^2 \theta_w}{16\pi \sin^4 \theta_w} \frac{m_t^2}{M_W^2} - \frac{11\alpha}{48\pi \sin^2 \theta_w} \ln \frac{M_H^2}{M_W^2} \dots$$

outdated

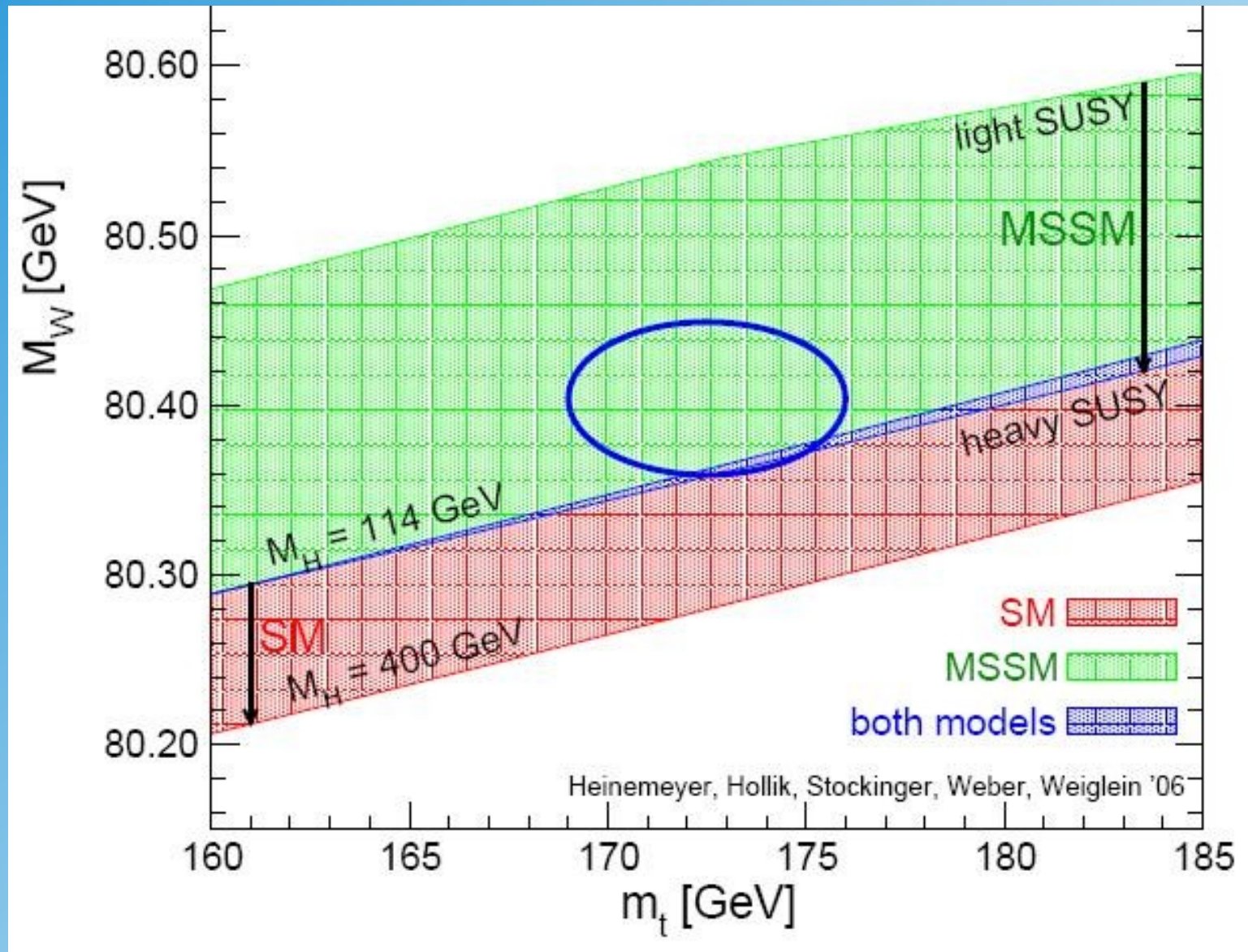
Theoretical prediction of  $\sin^2 \theta_{\text{eff}}$  as function of the Higgs mass.

# SM Higgs Mass Predictions



# W-Top-Higgs Mass Relation and Supersymmetry (SUSY)

SUSY Higgs  
is light!





# Higgs Direct Searches at LEP

Higgs-Fermion Coupling to b- and t-quarks:

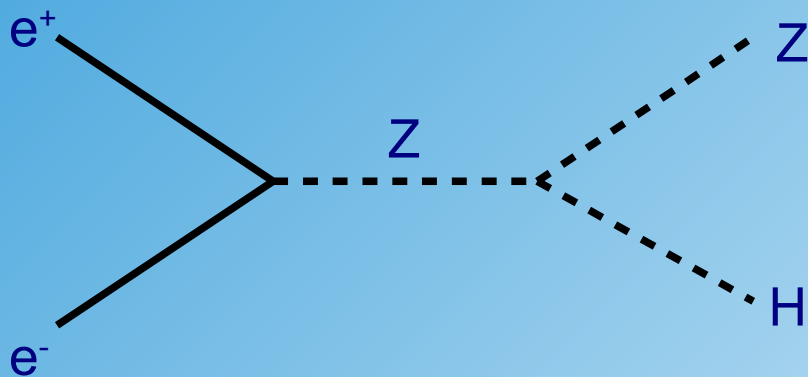
$$L_Y = -g_b \bar{L} \Phi b_R - g_t \bar{L} \tilde{\Phi} t_R \quad \text{with} \quad L = \begin{pmatrix} t \\ b \end{pmatrix}, \quad \Phi = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v+h \end{pmatrix}, \quad \tilde{\Phi} = i\tau_2 \Phi^*$$

Higgs couples to masses:

$$m_{b,t} = \frac{g_{b,t} v}{\sqrt{2}}$$

Electron-Positron annihilation has tiny Higgs coupling!

LEP main production Process (Higgs-Strahlung):  $e^+ e^- \rightarrow ZH$



**ZZH coupling is large!**

# ZH Signature at LEP

Higgs decays dominantly into heaviest fermions:

$$ZH \rightarrow Z \, b\bar{b} \rightarrow \ell\ell \, b\bar{b}, \, jj \, b\bar{b}, \, b\bar{b}b\bar{b}, \, \nu\nu \, b\bar{b}$$

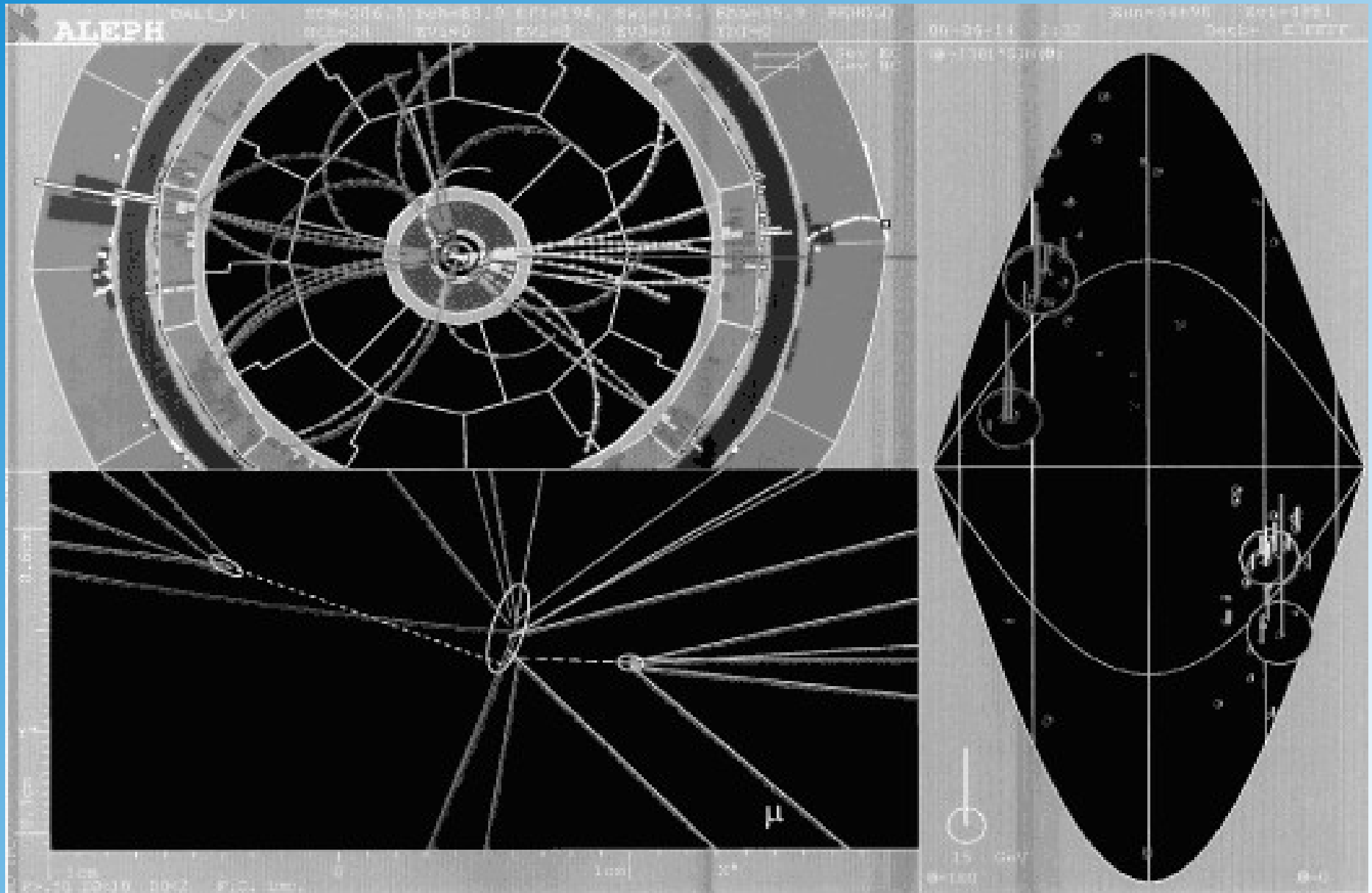
All decay channels require (double) b-tag (lifetime)

LEP2 with  $E_{\text{cms}} = 205 \text{ GeV}$ :

- installation of vertex detectors (high resolution silicon)
- sensitivity up to  $m_H = 114 \text{ GeV}$  ( $E_{\text{cms}} = m_H + m_Z$ )

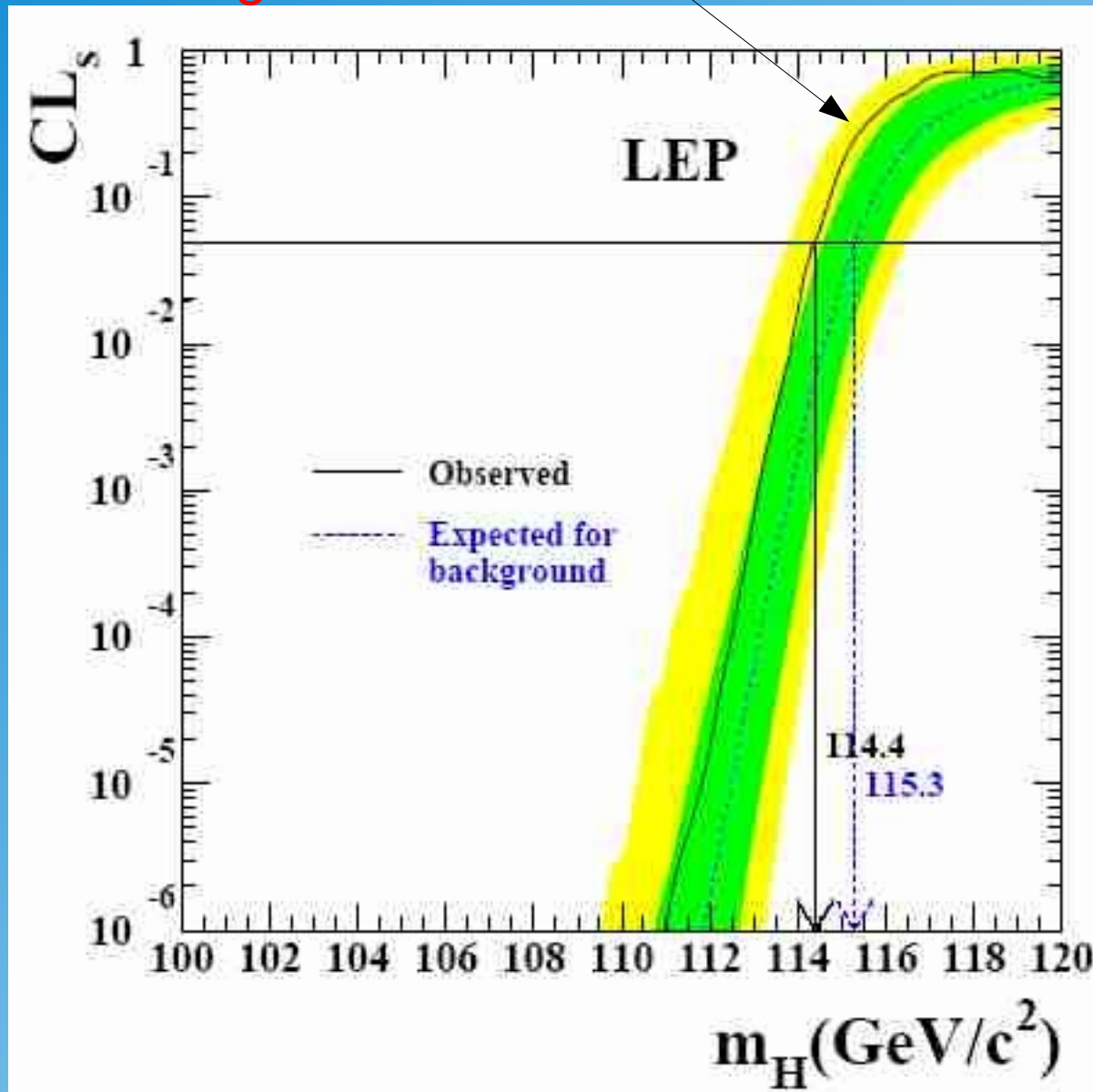


# bbjj-candidate at ALEPH



# Combined LEP2 Higgs Limit

no significant excess



Higgs excluded with:

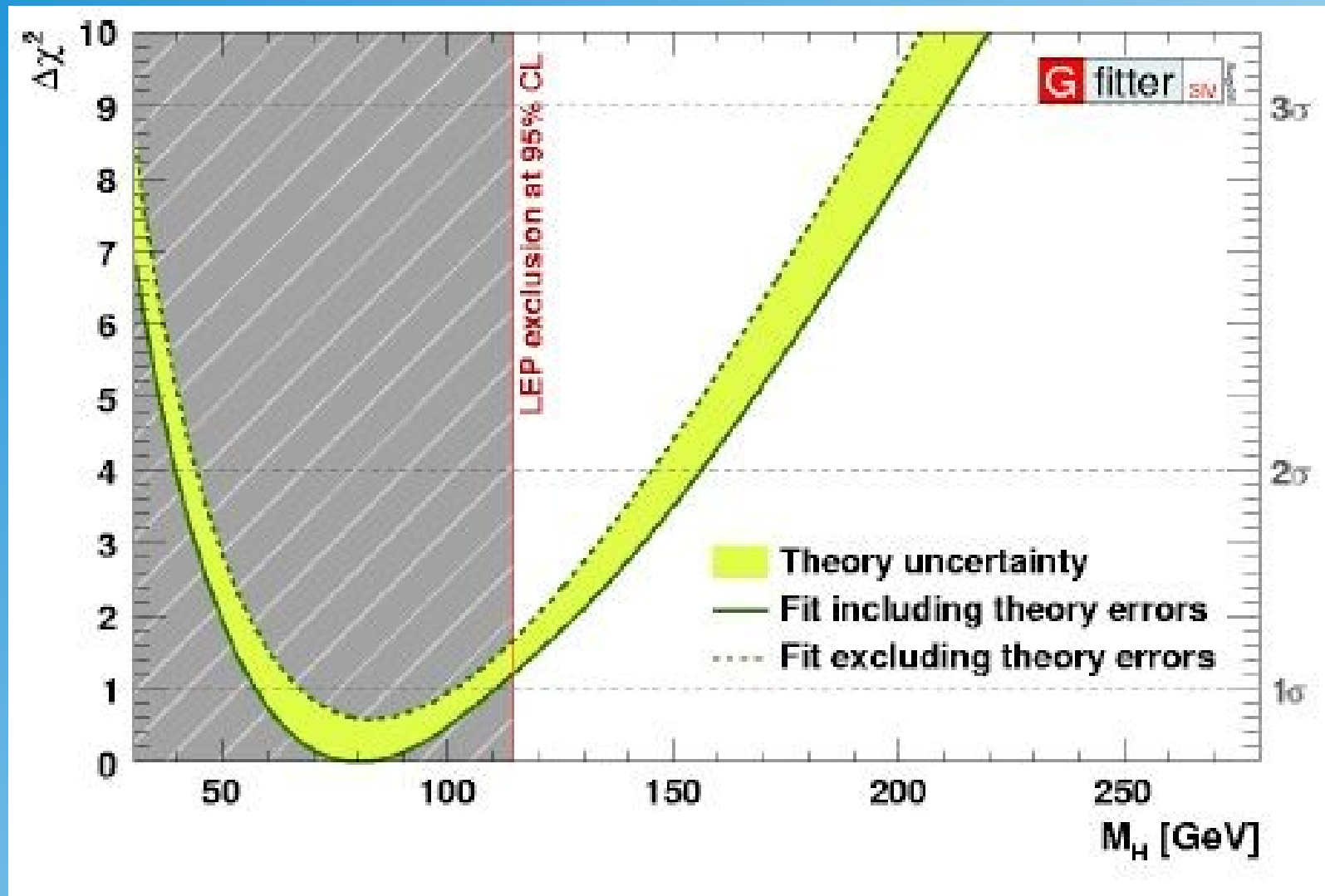
$$m_H < 114.4 \text{ GeV} \\ (\text{expected } 115.3)$$

at 95% CL

The bands shows the  
1 sigma and 2 sigma contours  
of the expected limit

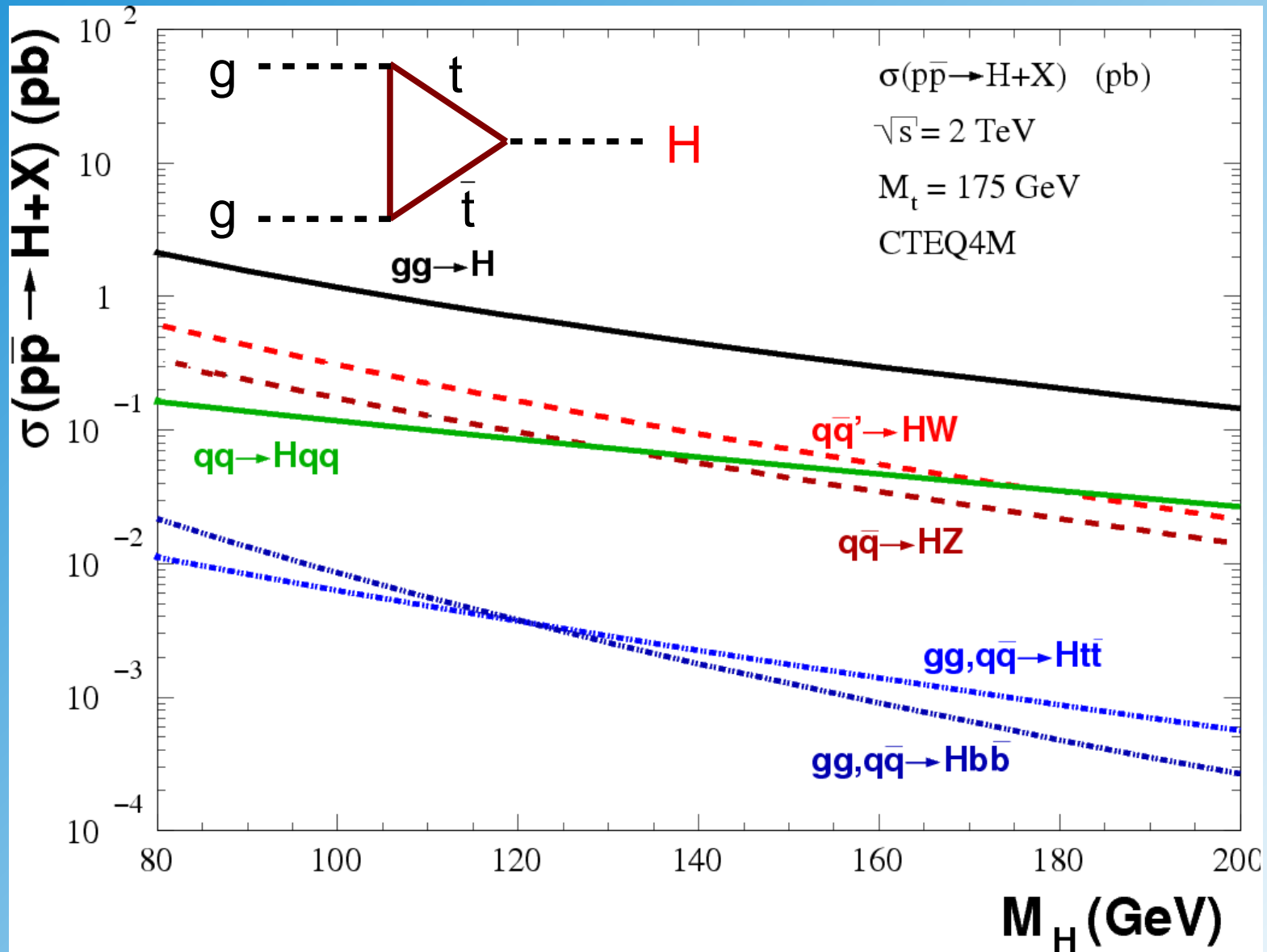
The black line shows the  
expected limit

# Higgs LEP2 Direct Limits



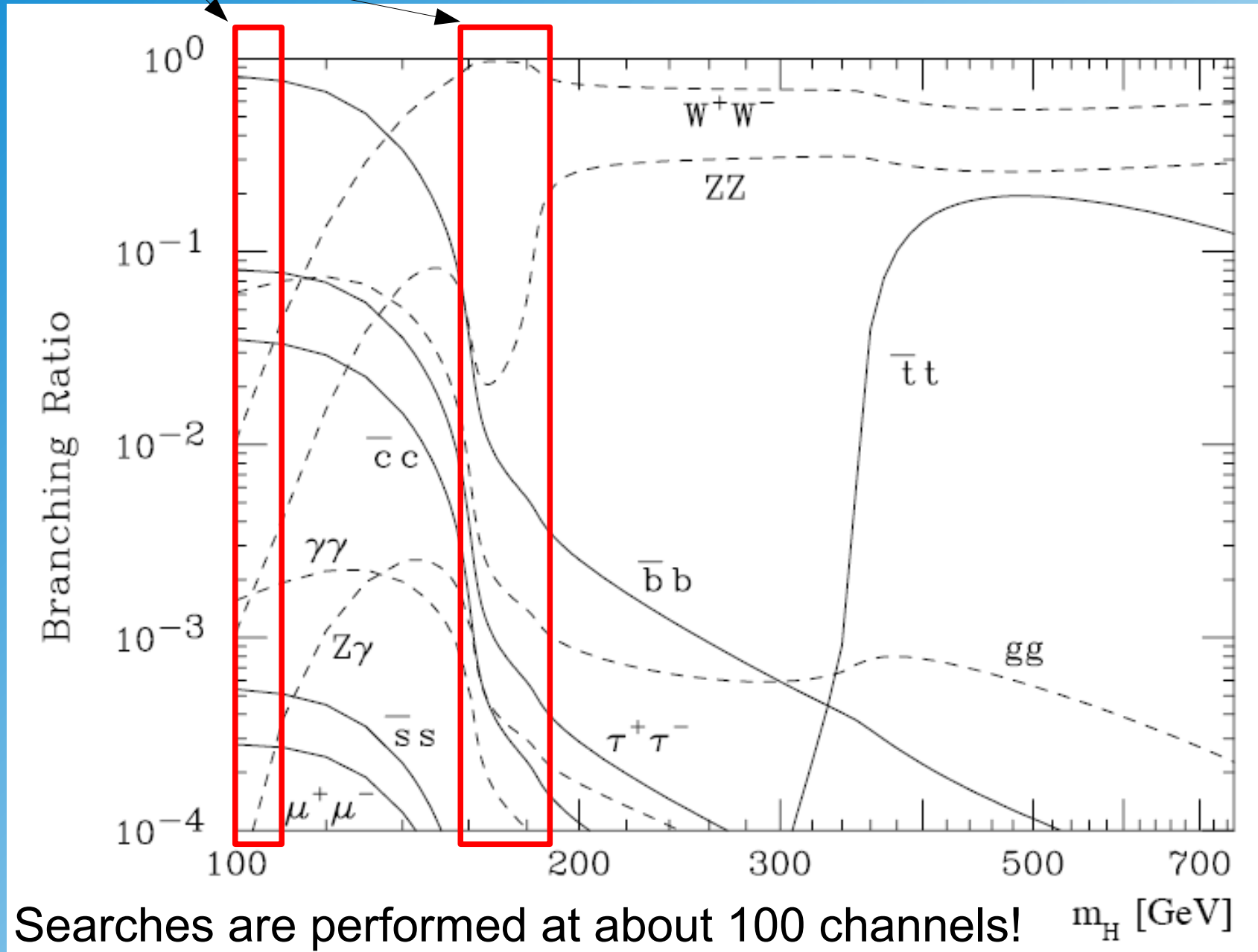
Gfitter performs a SM parameter fit.

# Higgs Production at Tevatron



# Higgs Decays

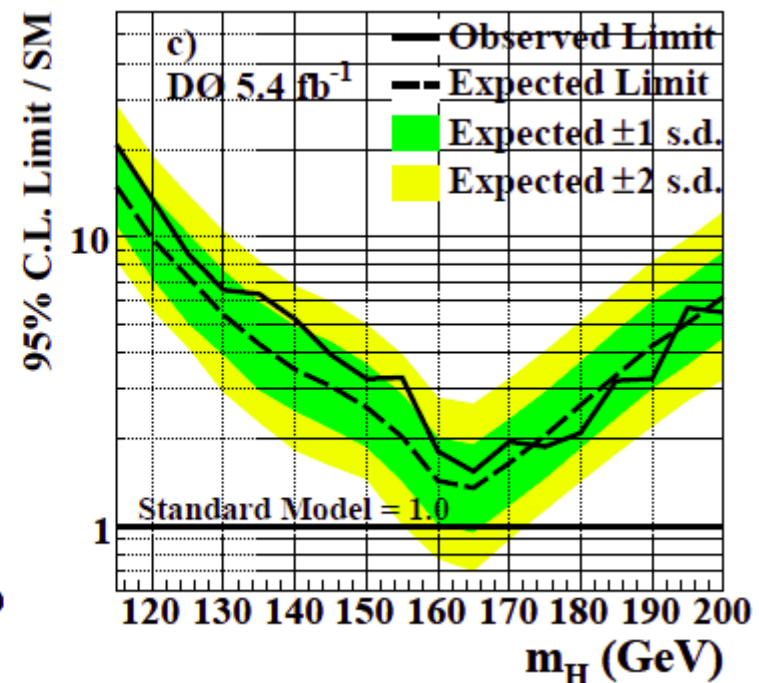
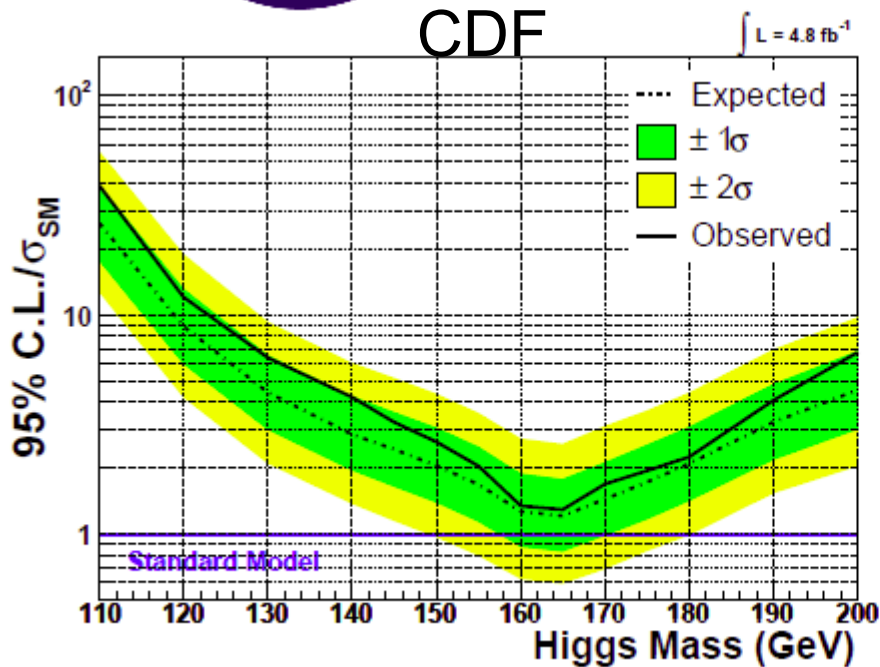
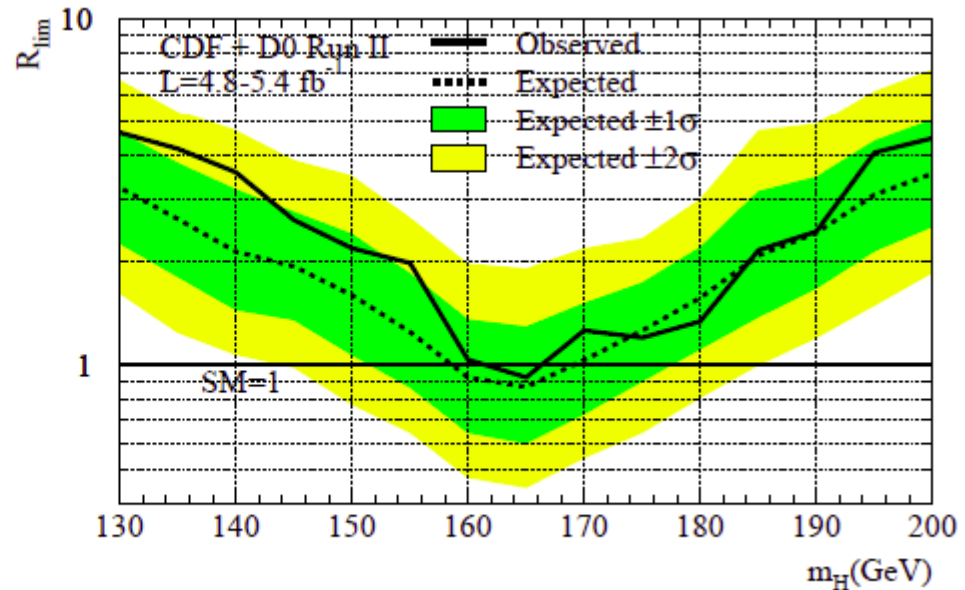
“easy” regions



# Combination Tevatron Searches

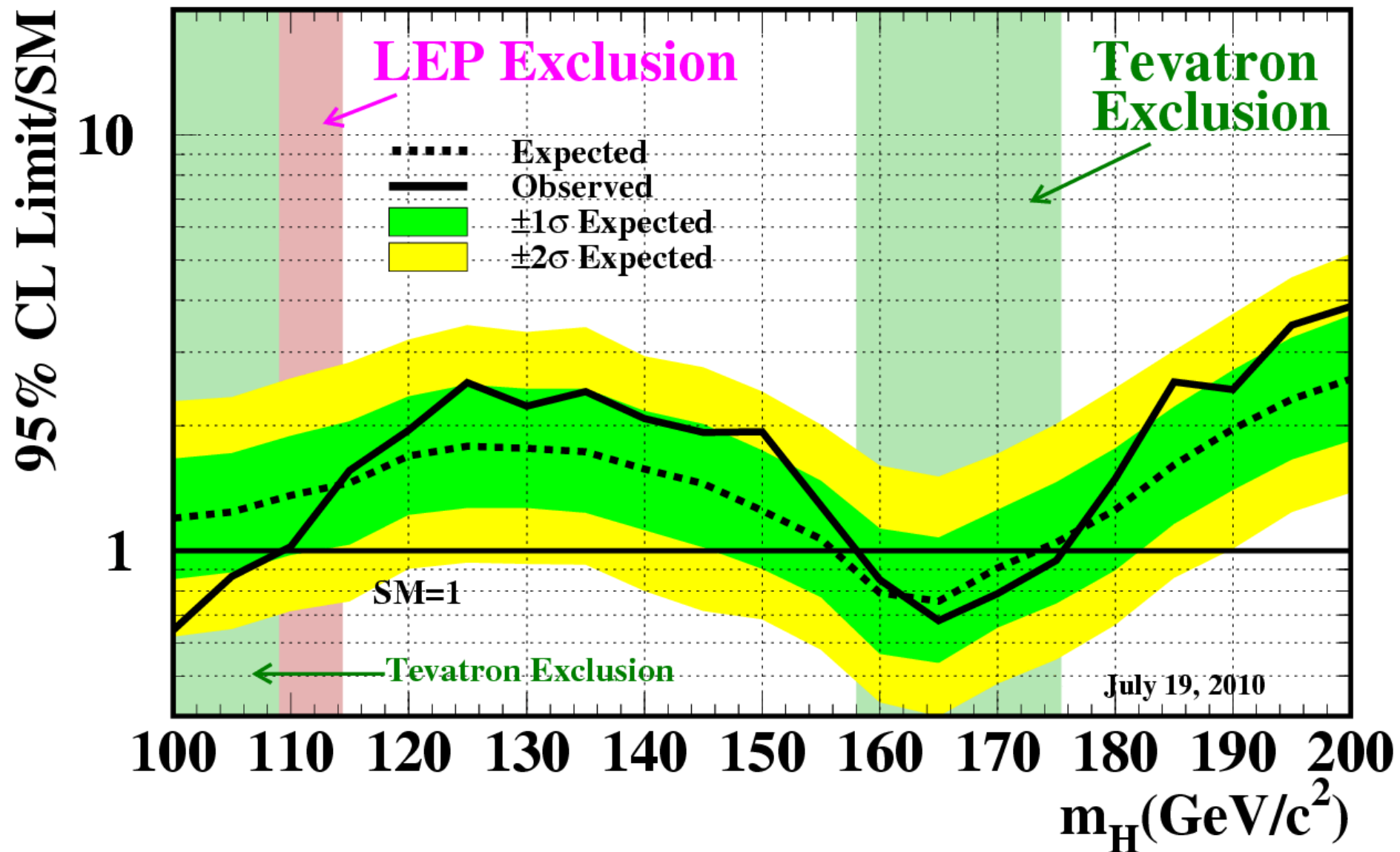


CDF





Tevatron Run II Preliminary,  $\langle L \rangle = 5.9 \text{ fb}^{-1}$





# Situation before LHC

Fits to electro-weak data:

$$m_H = 89^{+35}_{-26} \text{ GeV}$$

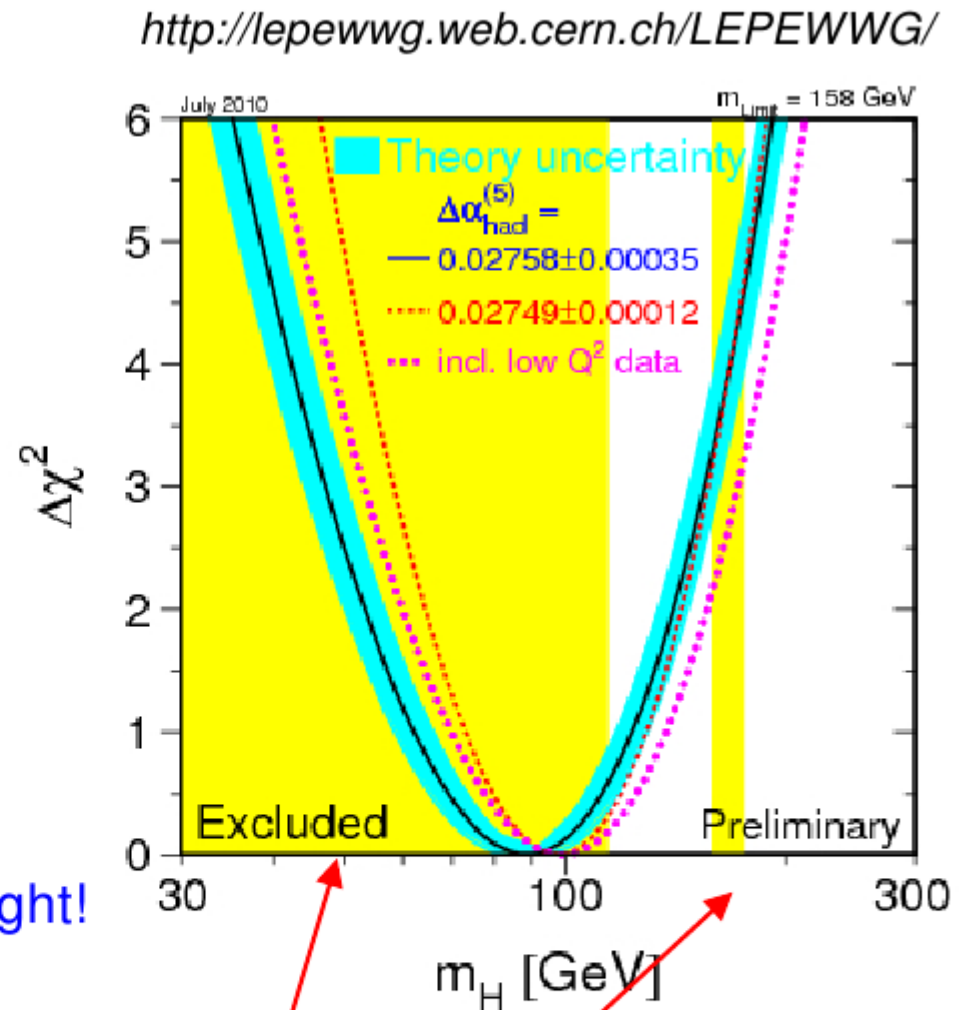
$$m_H < 158 \text{ GeV (95\% CL)}$$

Assumption for fit:

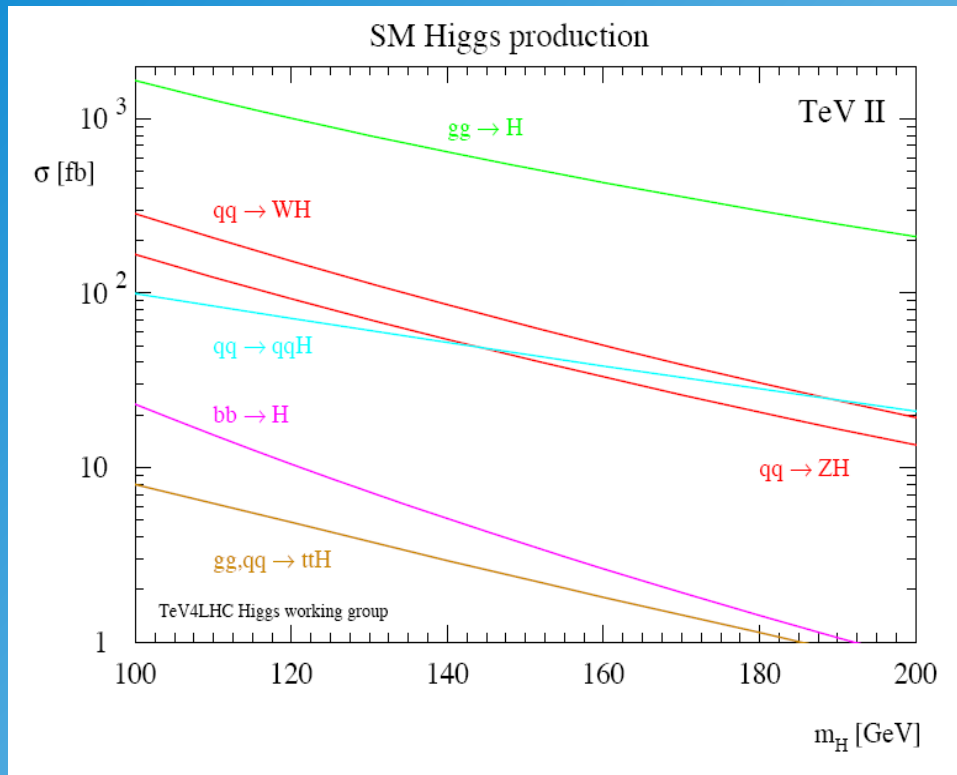
- SM including Higgs

If existing, Higgs seems to be light!

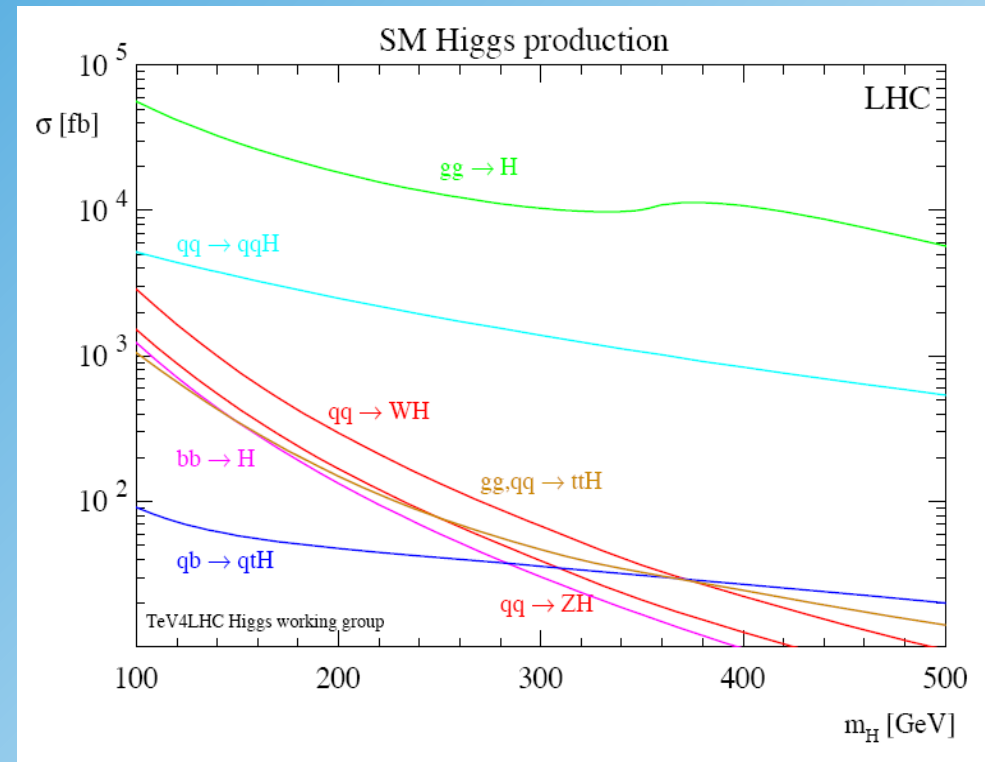
- Direct searches at LEP:  
 $m_H > 114.4 \text{ GeV @ 95\% CL}$
- Direct searches at Tevatron



# Higgs Production at LHC



proton-antiproton at Tevatron  $s^{1/2}=2$  TeV

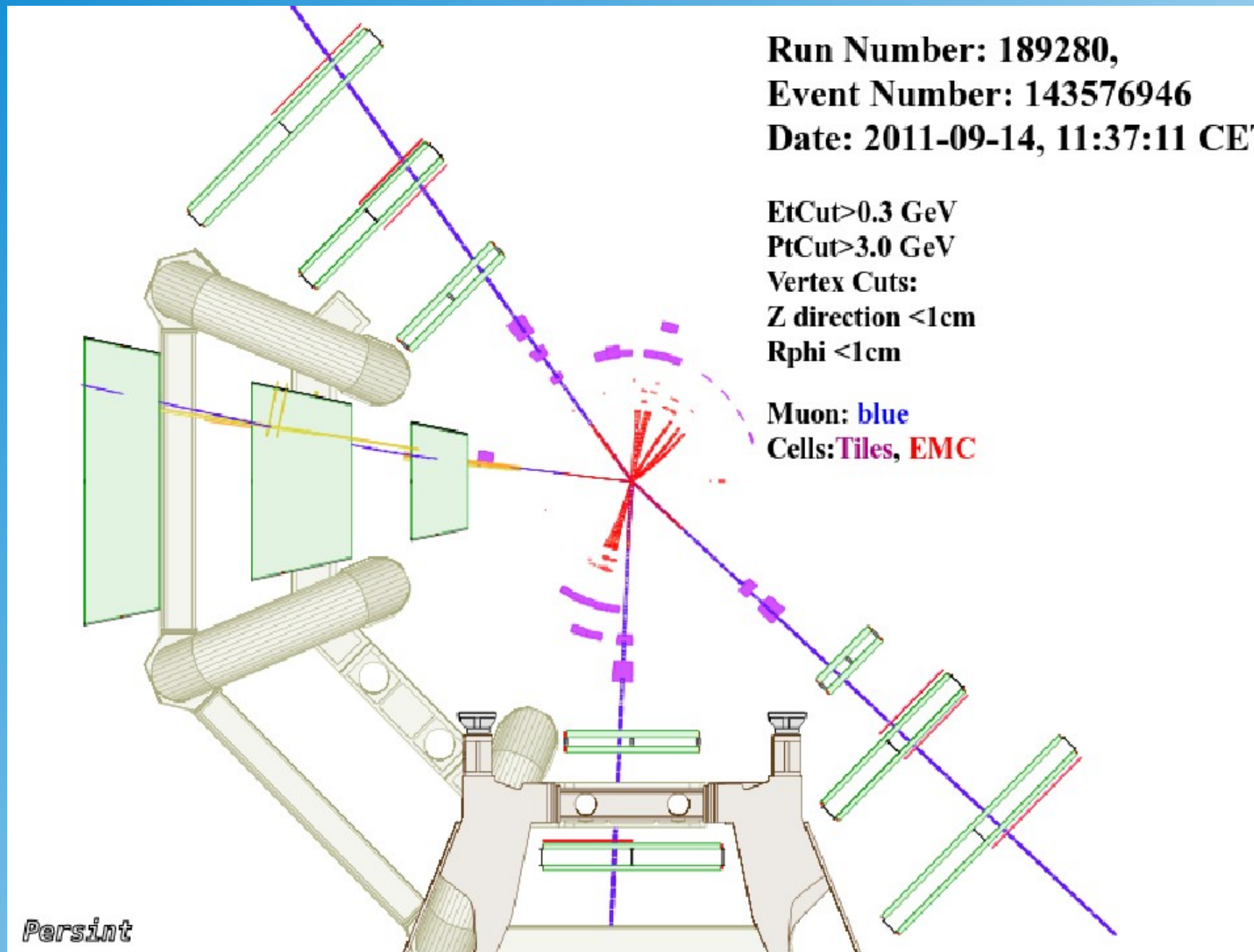


proton-proton at LHC  $s^{1/2}=14$  TeV

much larger cross sections!

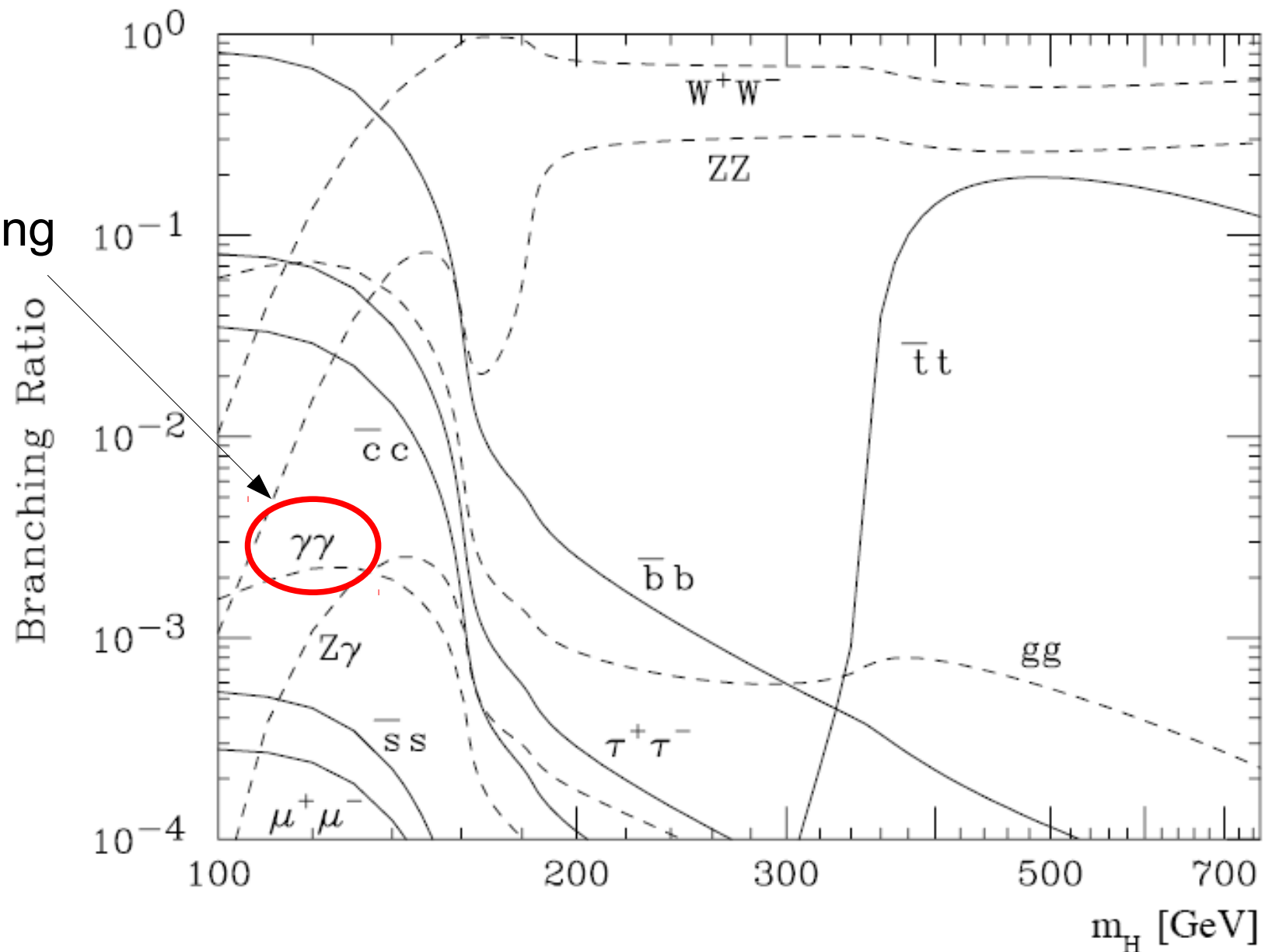
also higher luminosity!

# Candidate $H \rightarrow ZZ \rightarrow \mu\mu\mu\mu$

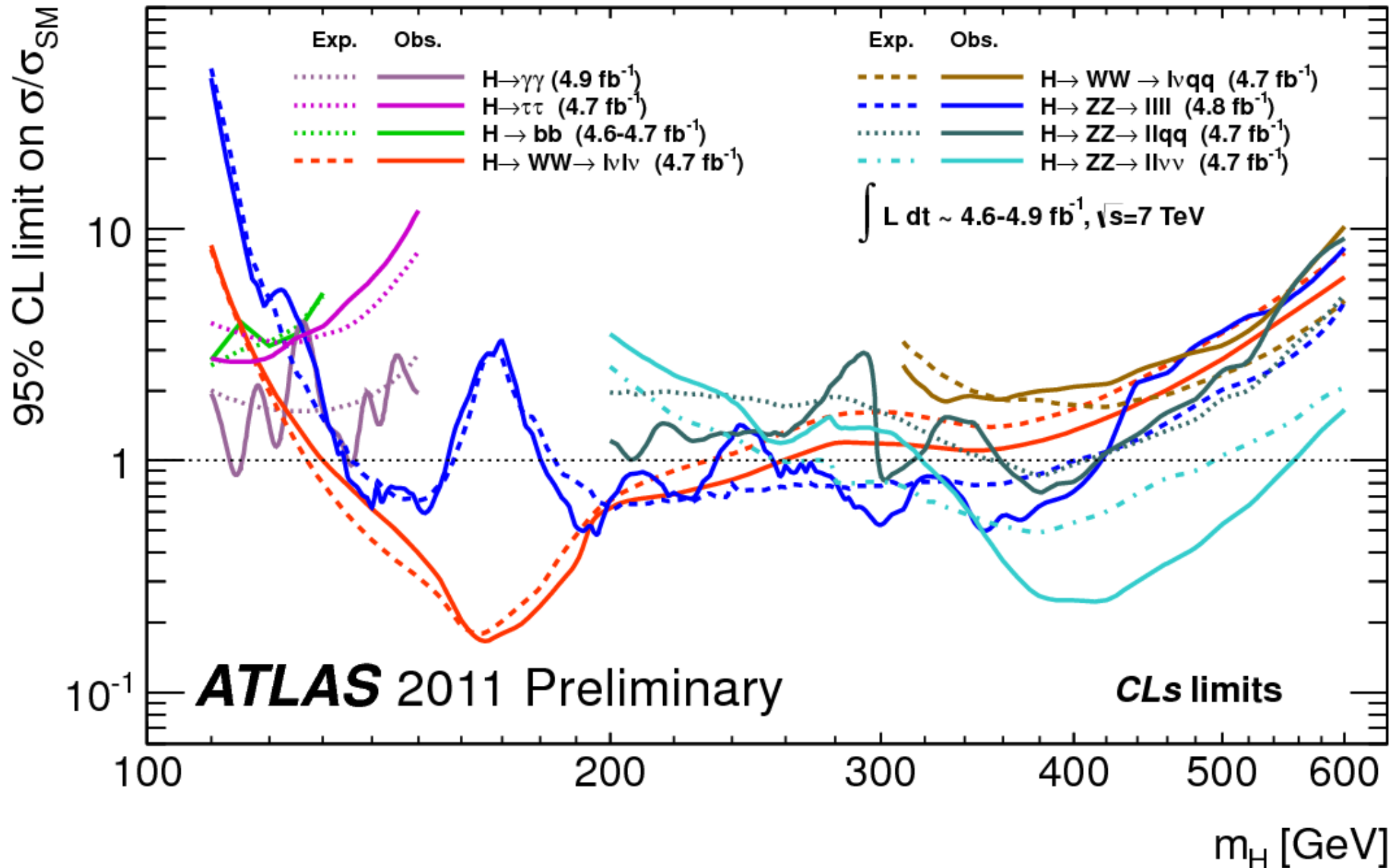


# Special Channel $H \rightarrow \gamma\gamma$ (high inv. mass resolution)

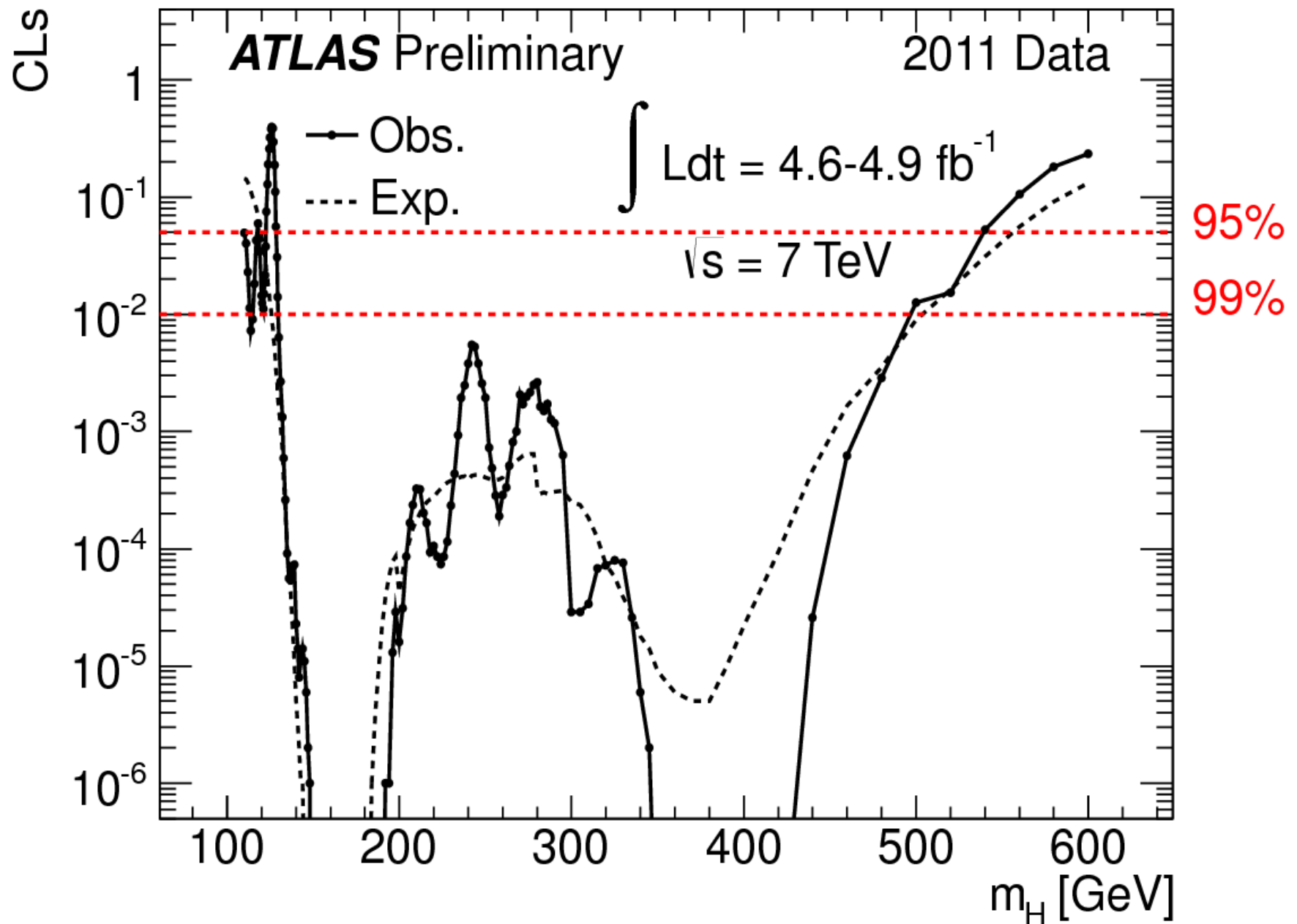
Higgs coupling  
via top loop



# Summary ATLAS Searches

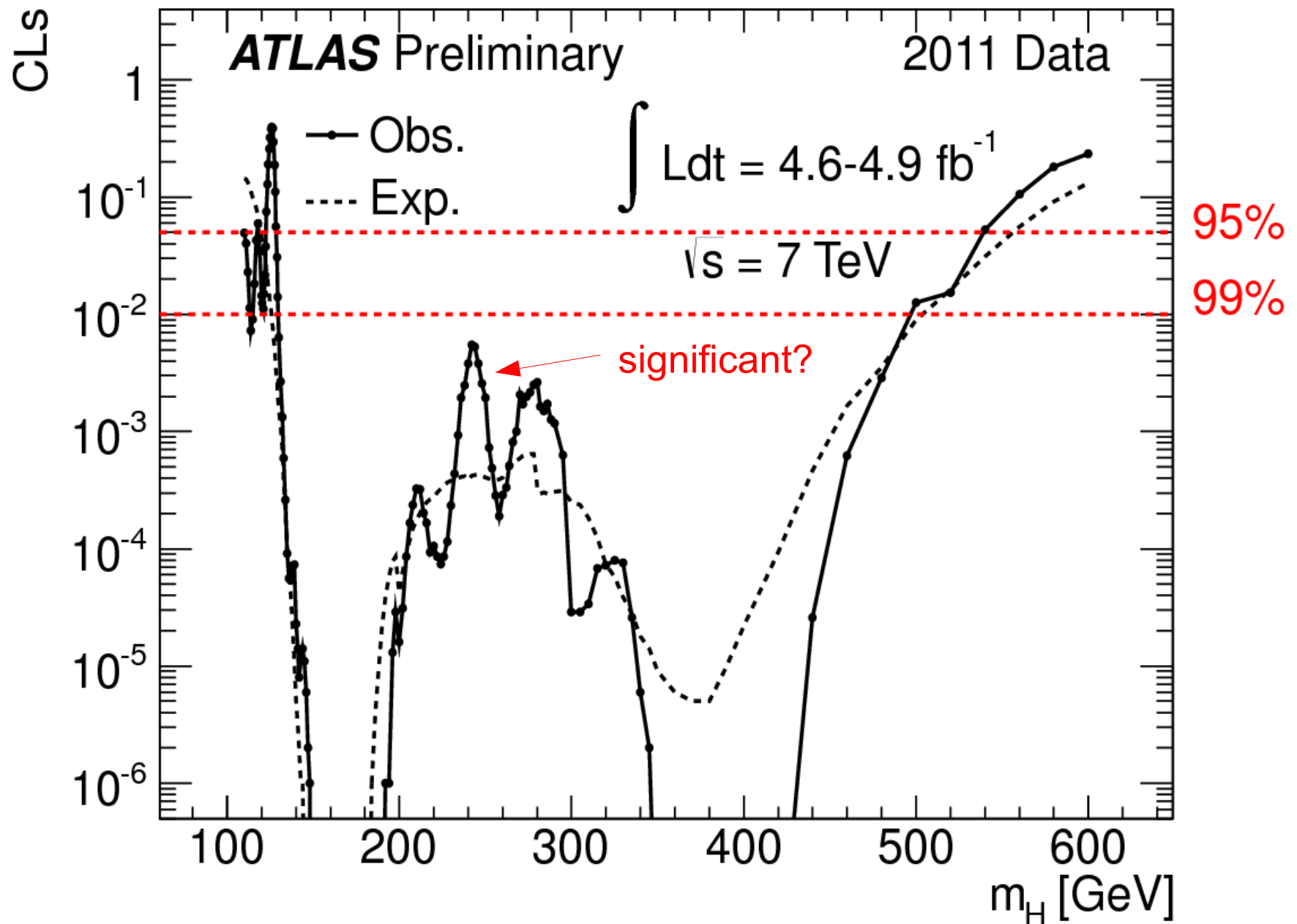


# ATLAS Combined Result



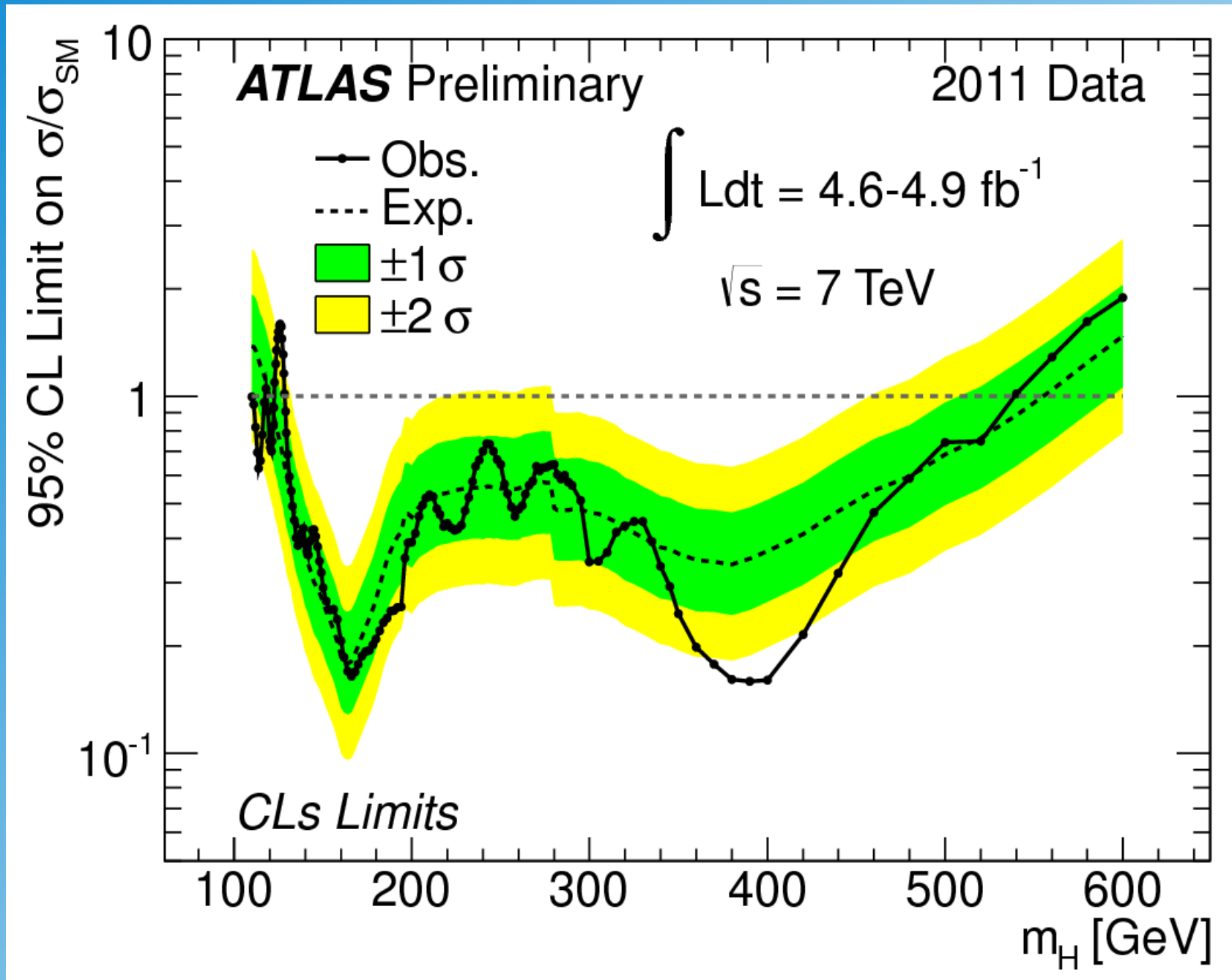


# ATLAS Combined Result

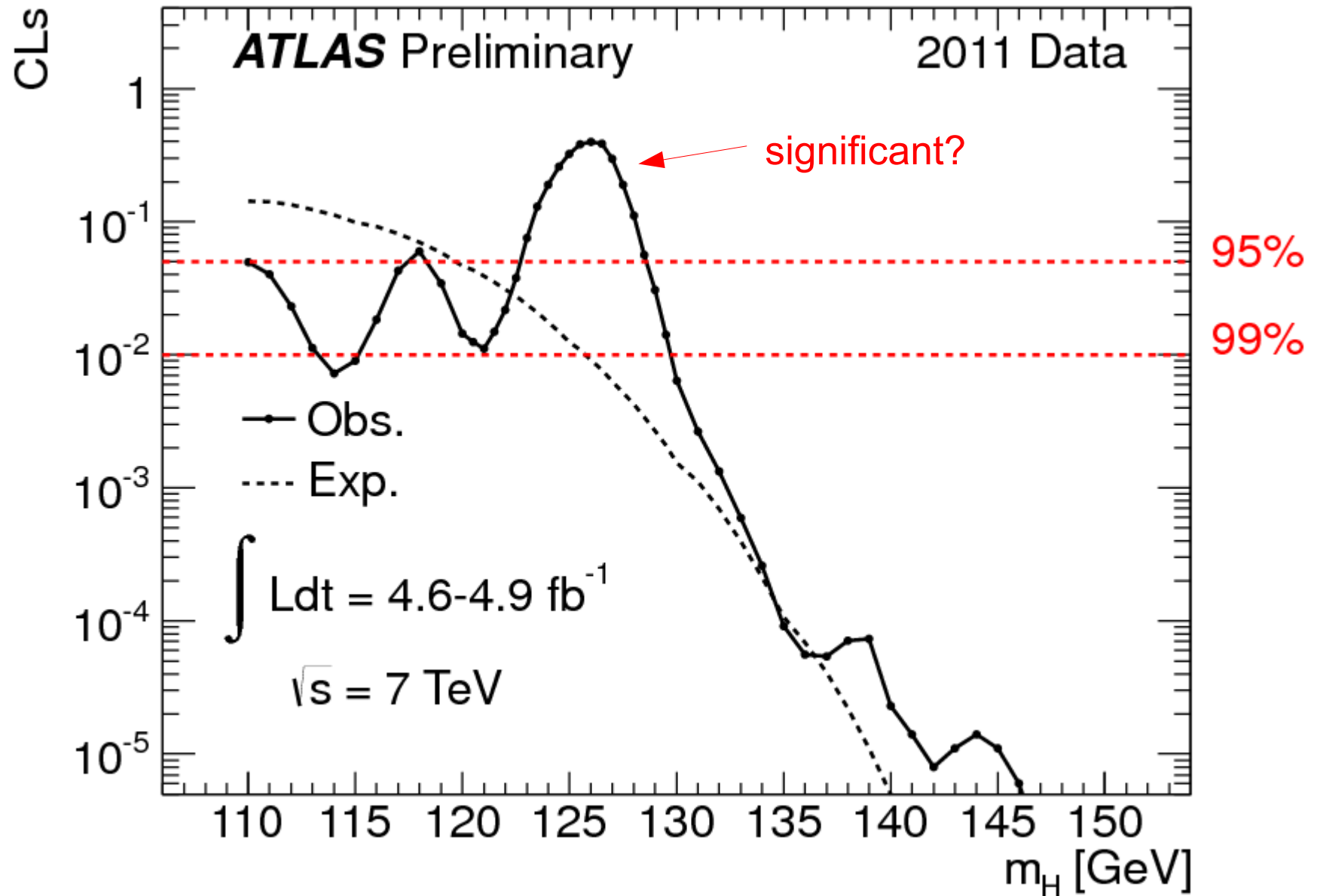




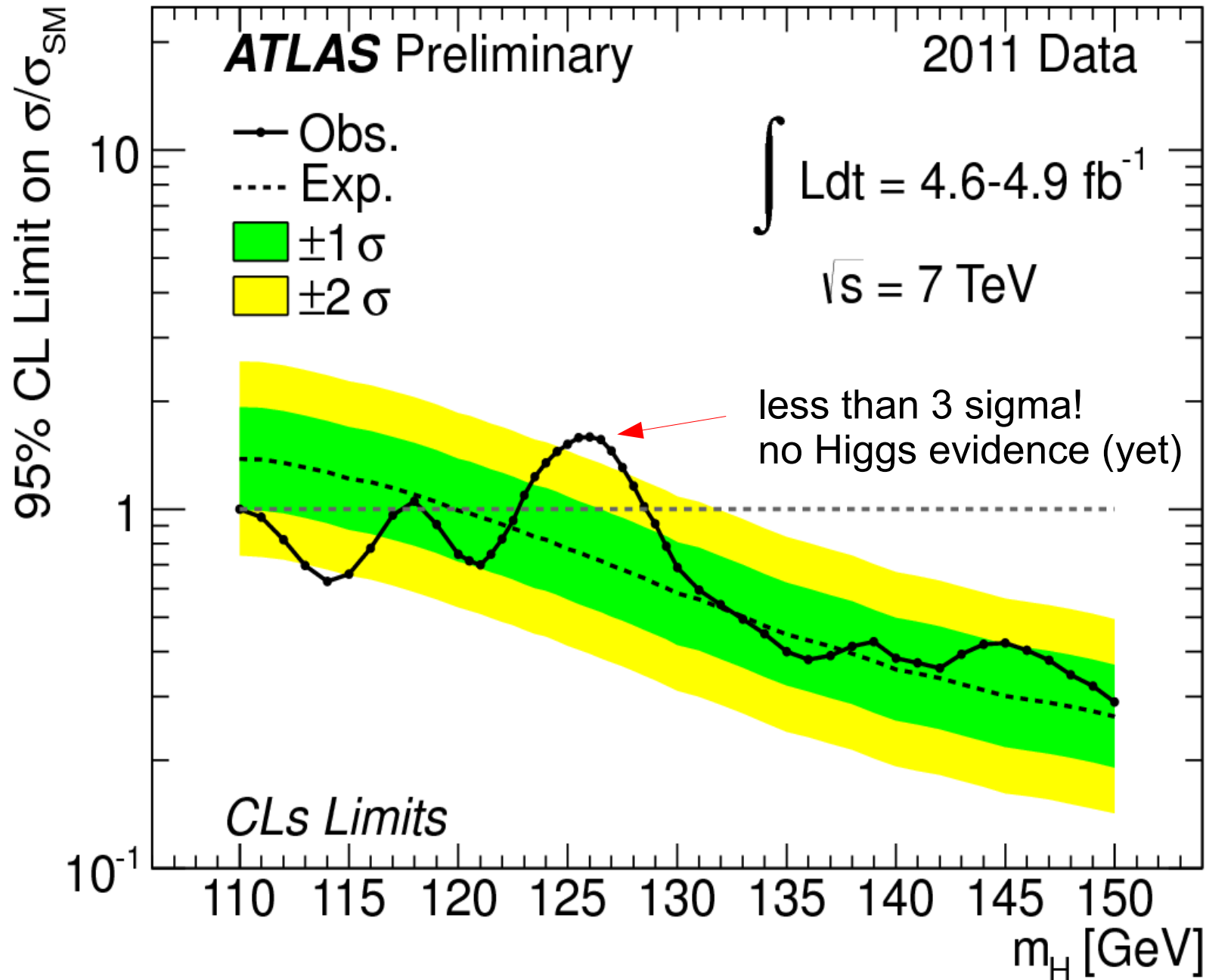
# ATLAS Combined Result



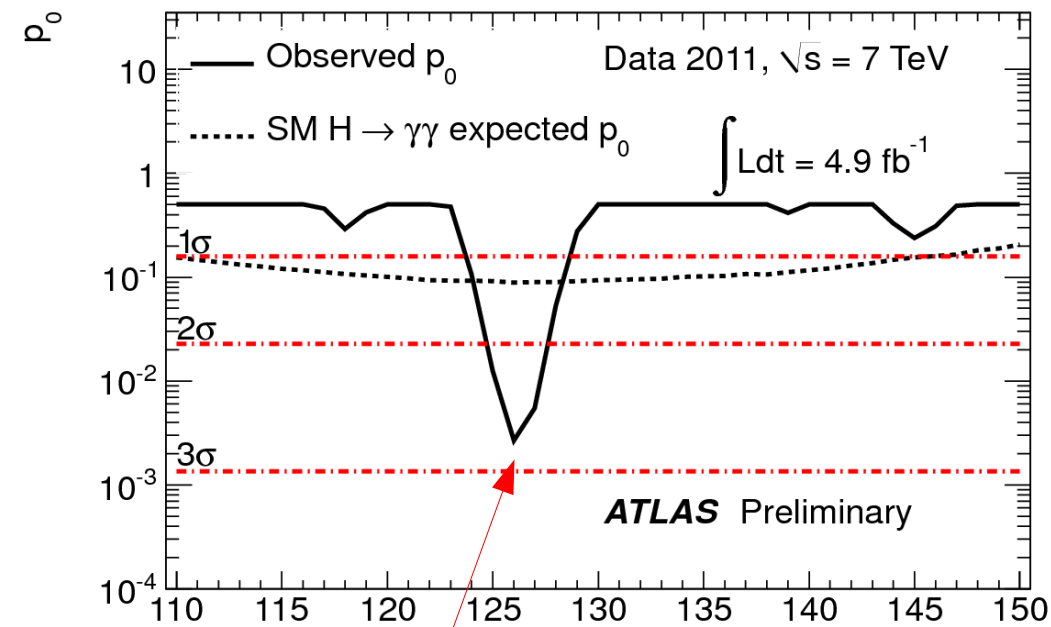
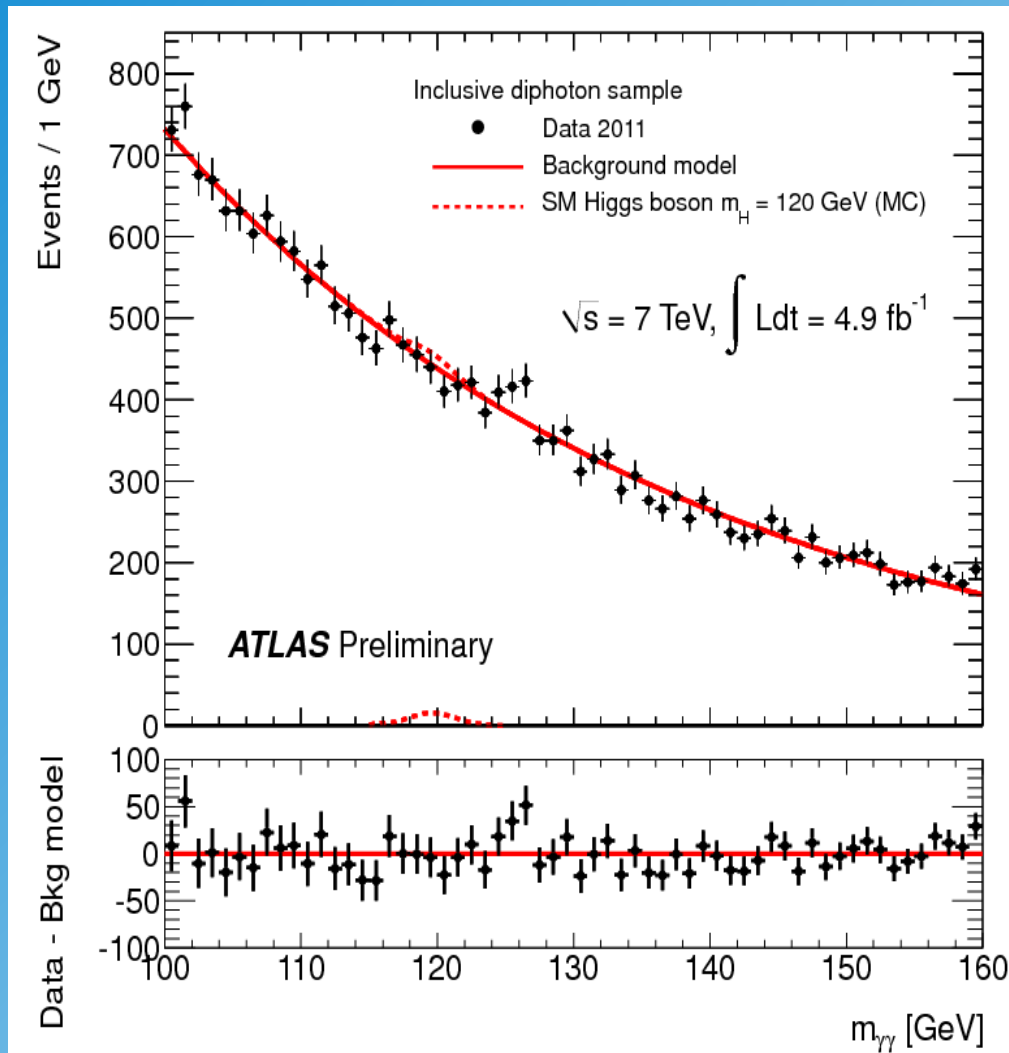
# ATLAS Combined Result



# ATLAS Combined Result

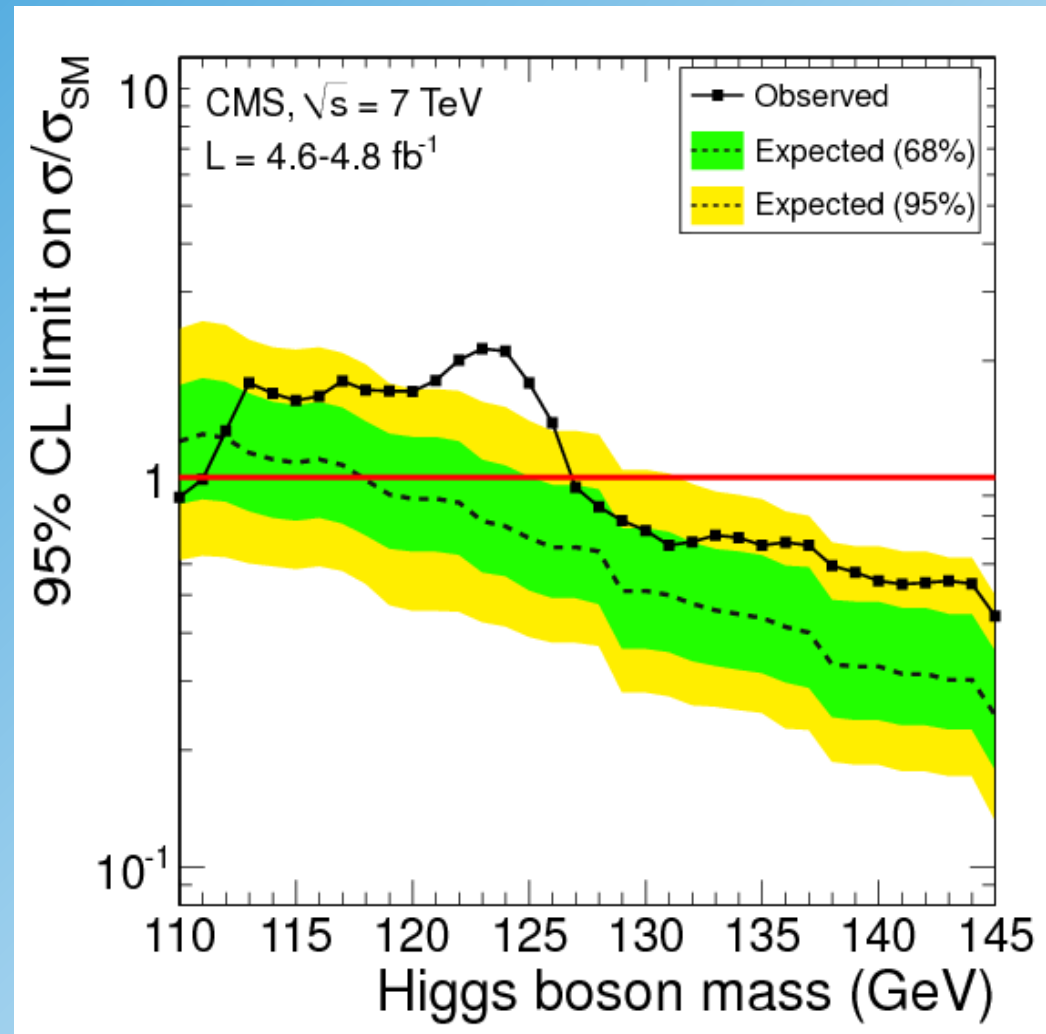
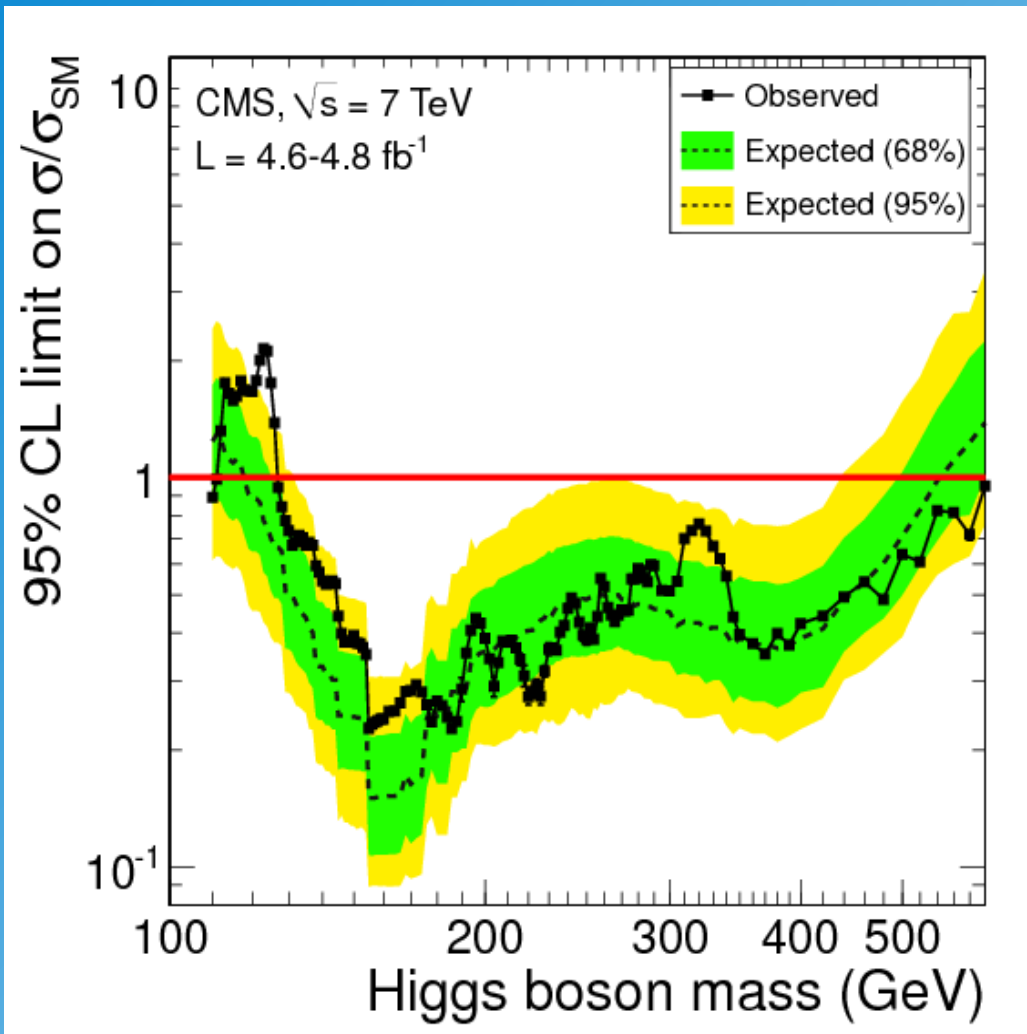


# Higgs $\rightarrow \gamma\gamma$ Search



note, signal too strong for Higgs!

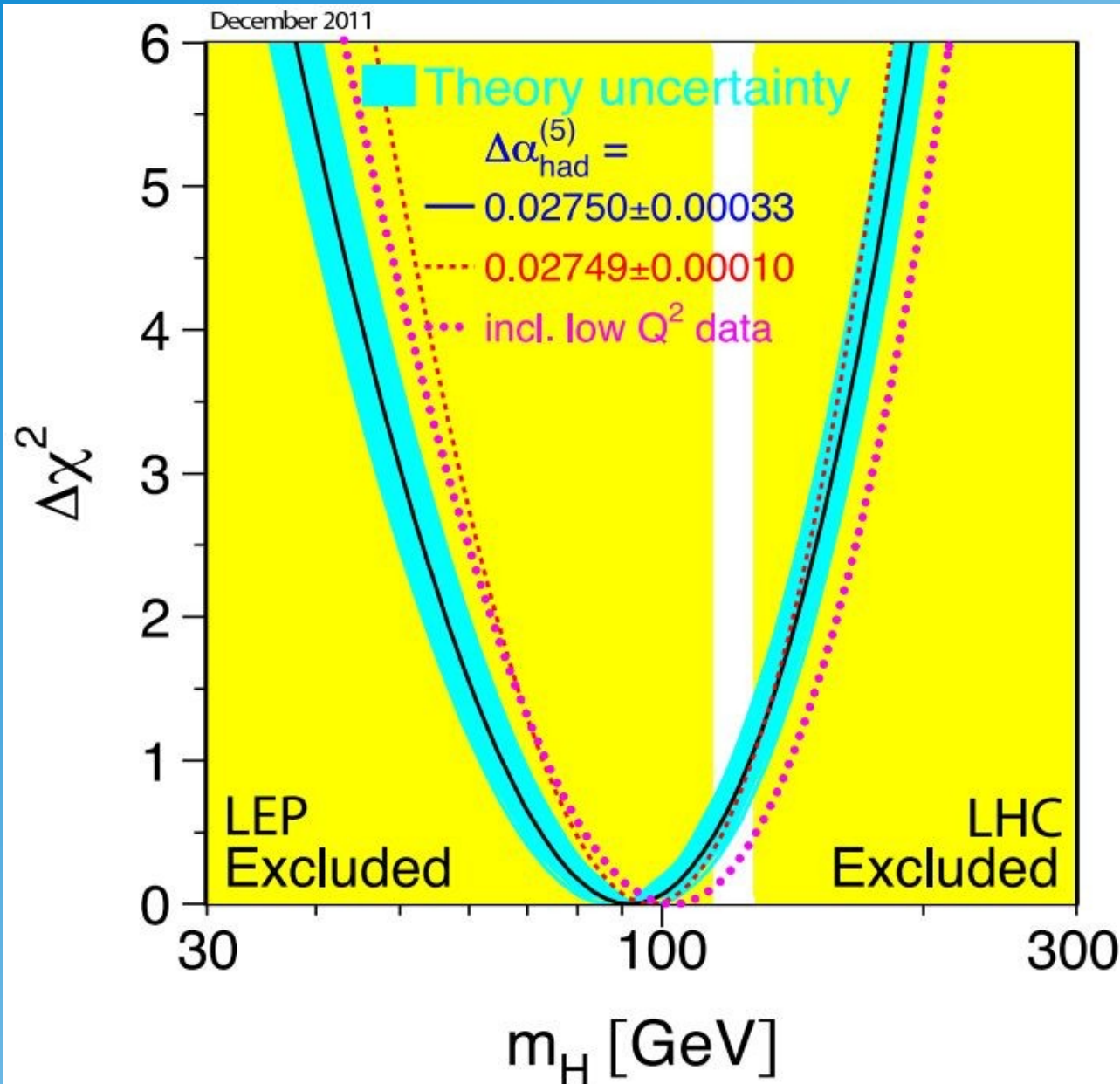
# CMS Higgs Search (2011 data)



very consistent with ATLAS results, Higgs?

# Situation in June 2012

(before summer conferences)



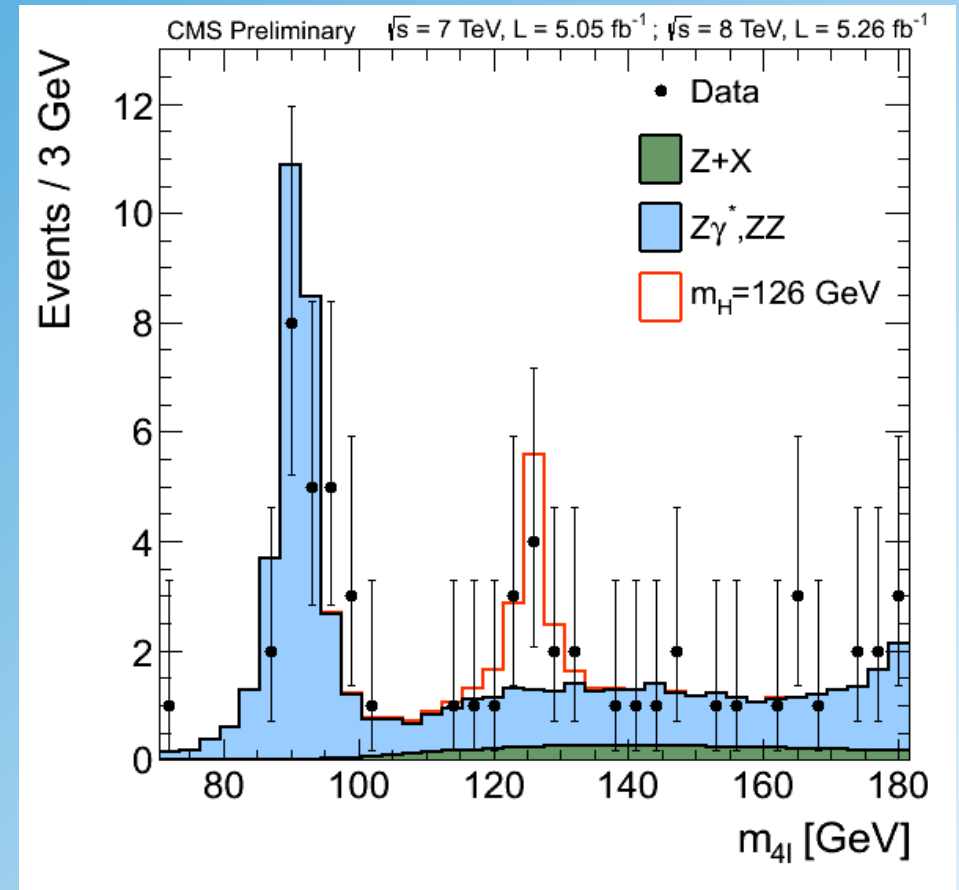
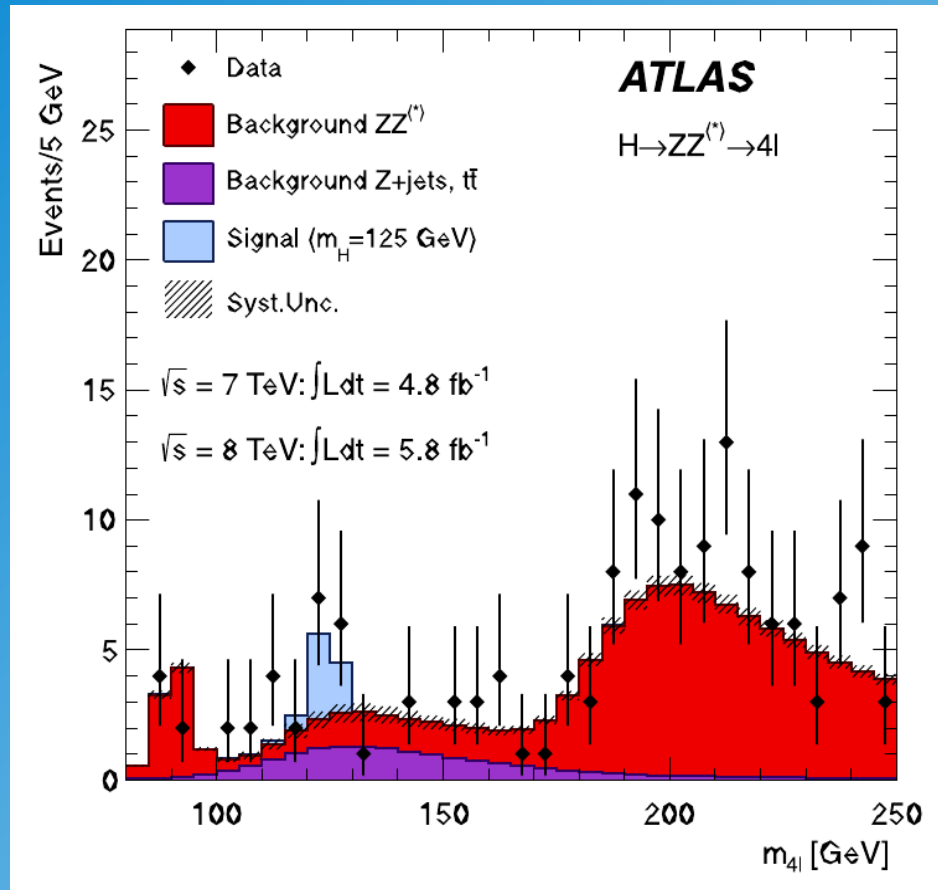
Non-excluded region  
is the most difficult!

small window around  
125 GeV not excluded  
at 95% CL

→ Updates July 4th 2012!



# Results from ATLAS+CMS

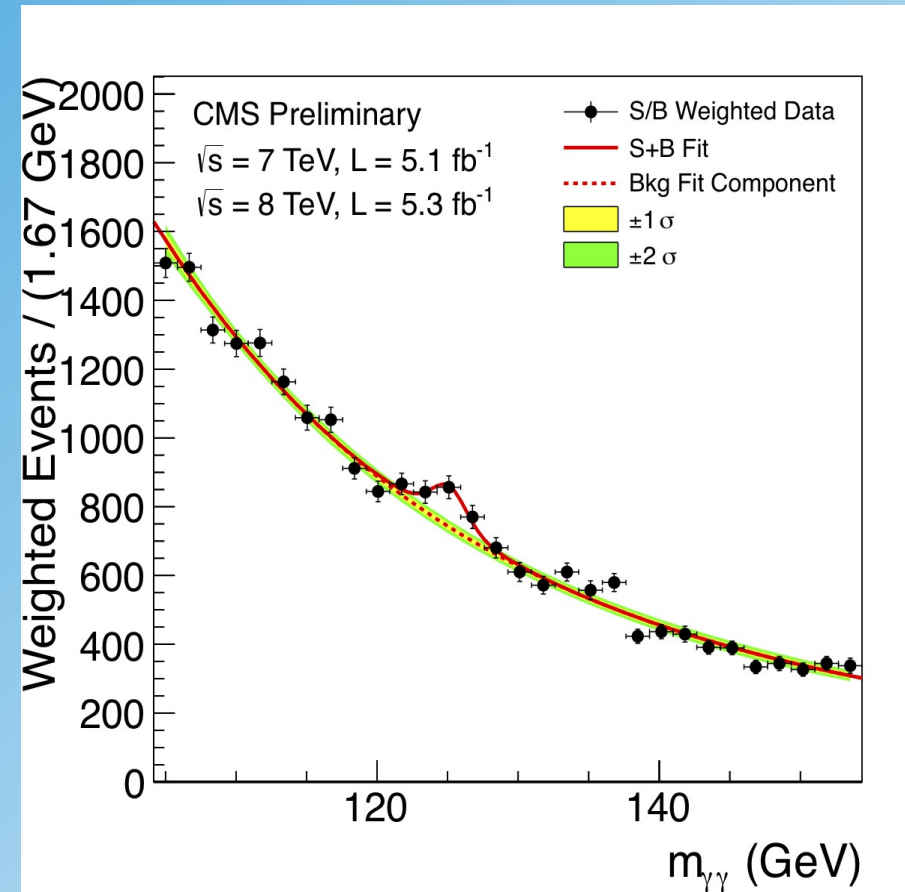
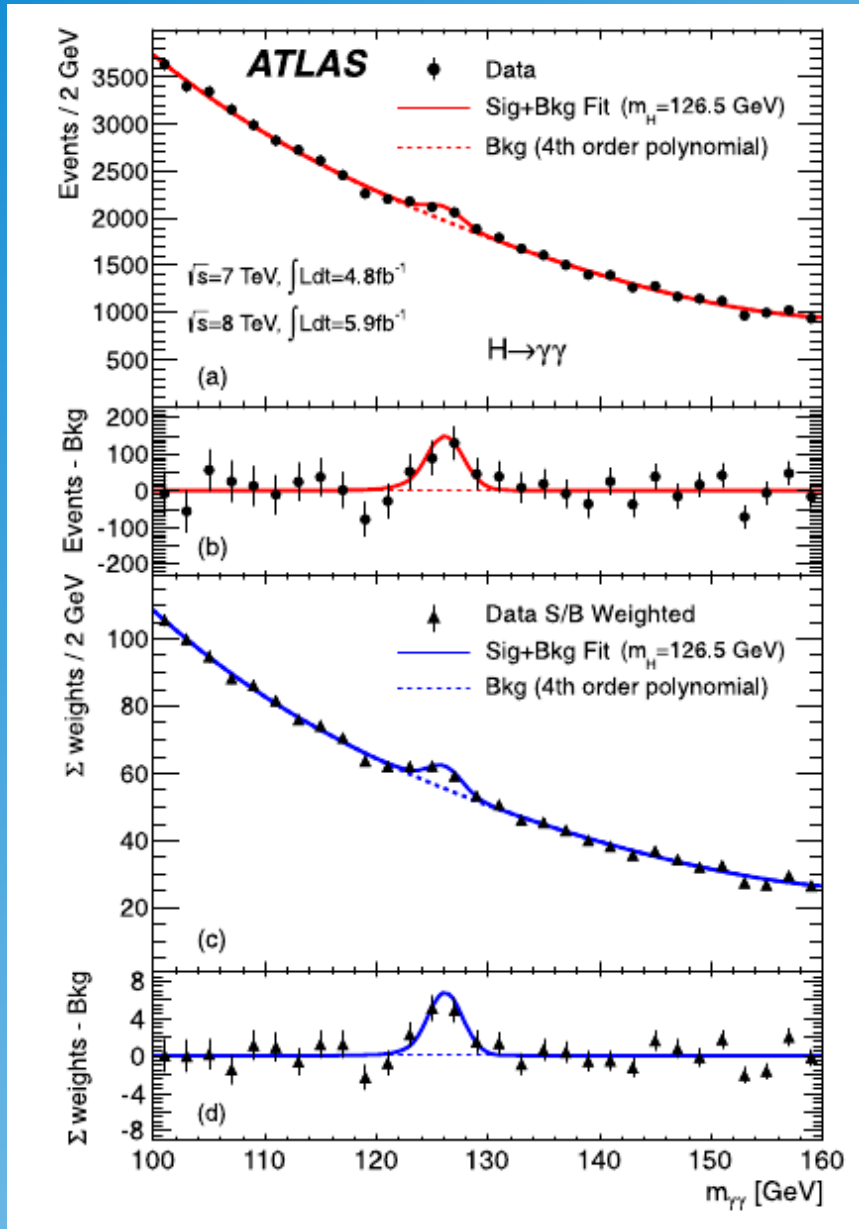


$H \rightarrow ZZ \rightarrow \text{leptons}$

Growing Signal  $H \rightarrow ZZ$

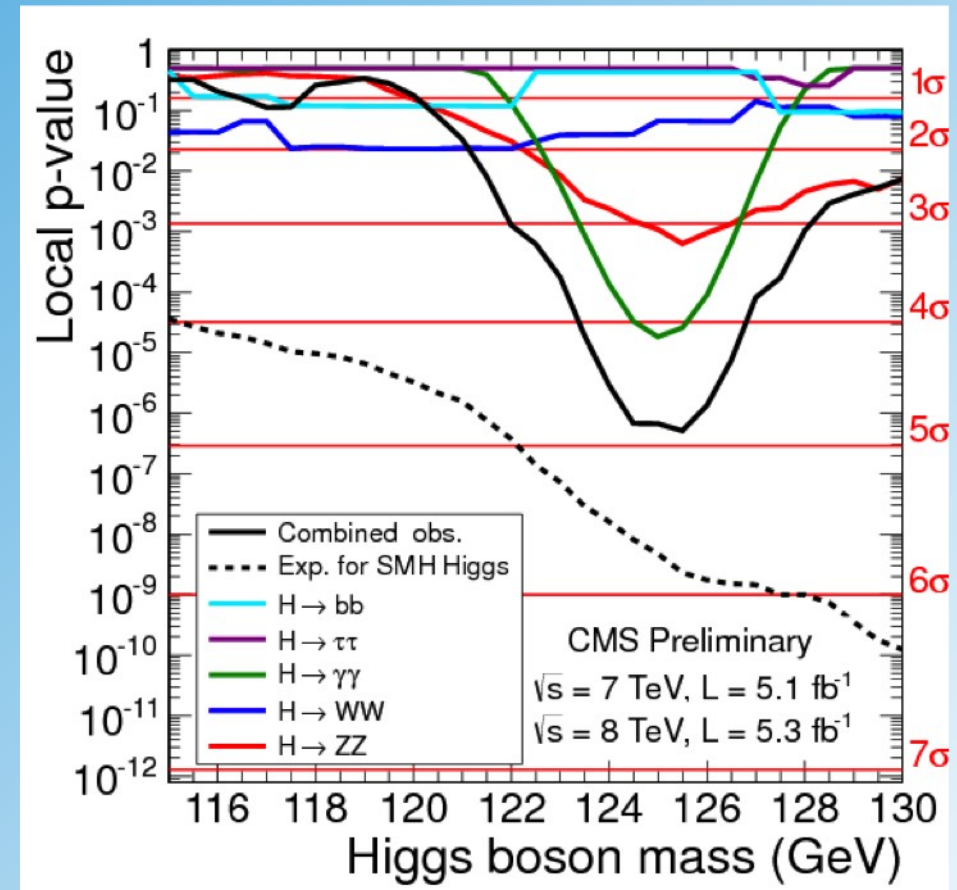
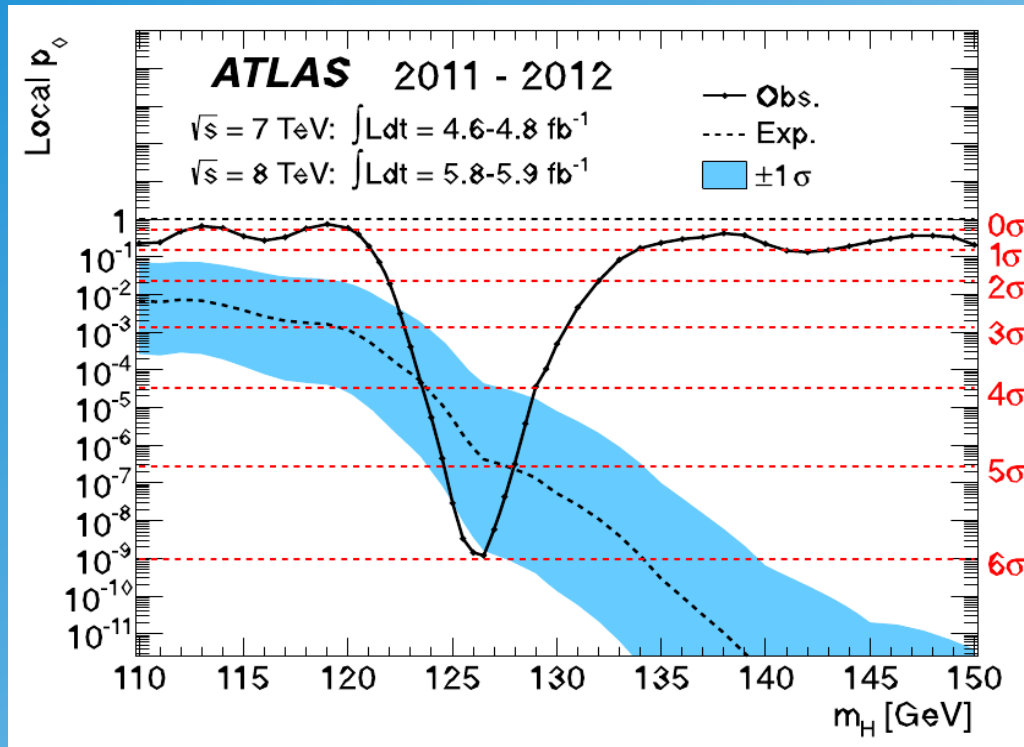
# Results from ATLAS+CMS

$H \rightarrow \text{gamma gamma}$



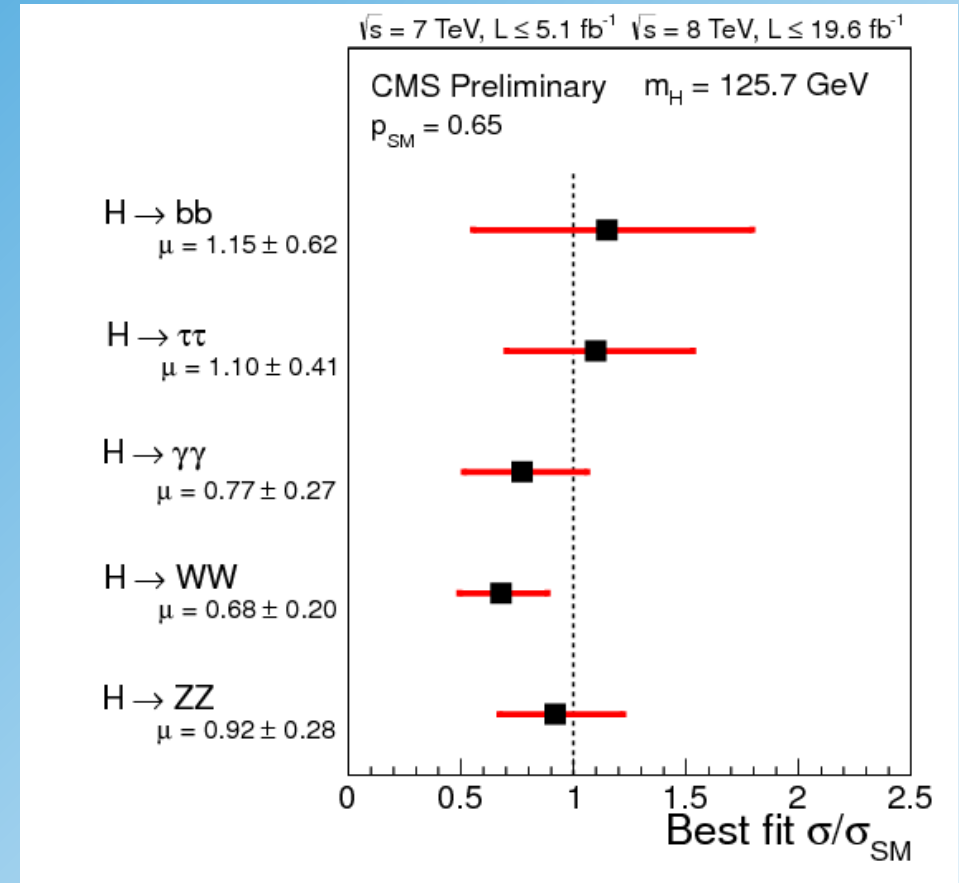
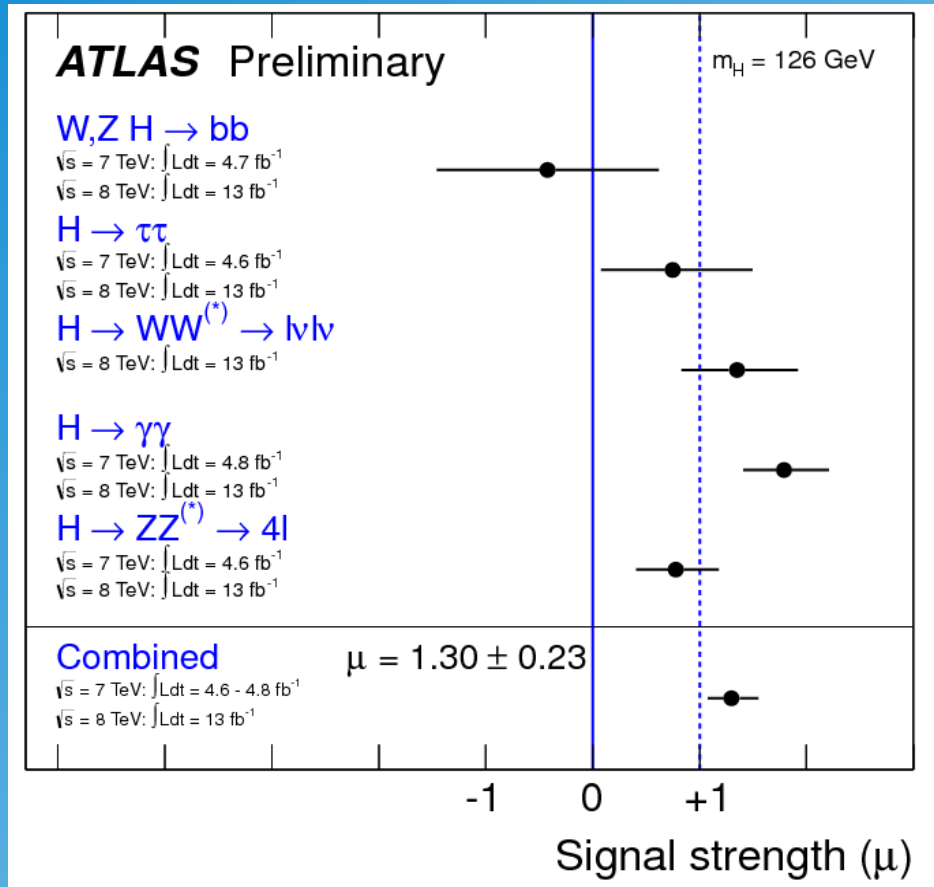
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## Combined local probabilities



Note: difference between local and global probabilities!

# Some First Higgs Decay Results from ATLAS+CMS



Couplings strengths consistent with SM expectation

# Summary

- Indirect constraints predict a light Higgs in the Standard Model
- Experimentally a Higgs with  $m_H \sim 125$  GeV is most difficult to find
- Direct searches exclude (June 2012) Higgs except for a small region around  $m_H \sim 125$  GeV
- Excess of Higgs candidate events at  $m_H \sim 125$  by ATLAS and CMS. More events are seen than expected by SM Higgs model.
- Interestingly,  $m_H \sim 125$ -130 GeV is theoretically also favored by vacuum stability and triviality reasons.  
No new physics required up to very high mass scales!
- At  $m_H \sim 125$  GeV Higgs almost all decay channels can be measured
  - allows for many experimental tests
  - nature is very kind to particle physicists!



