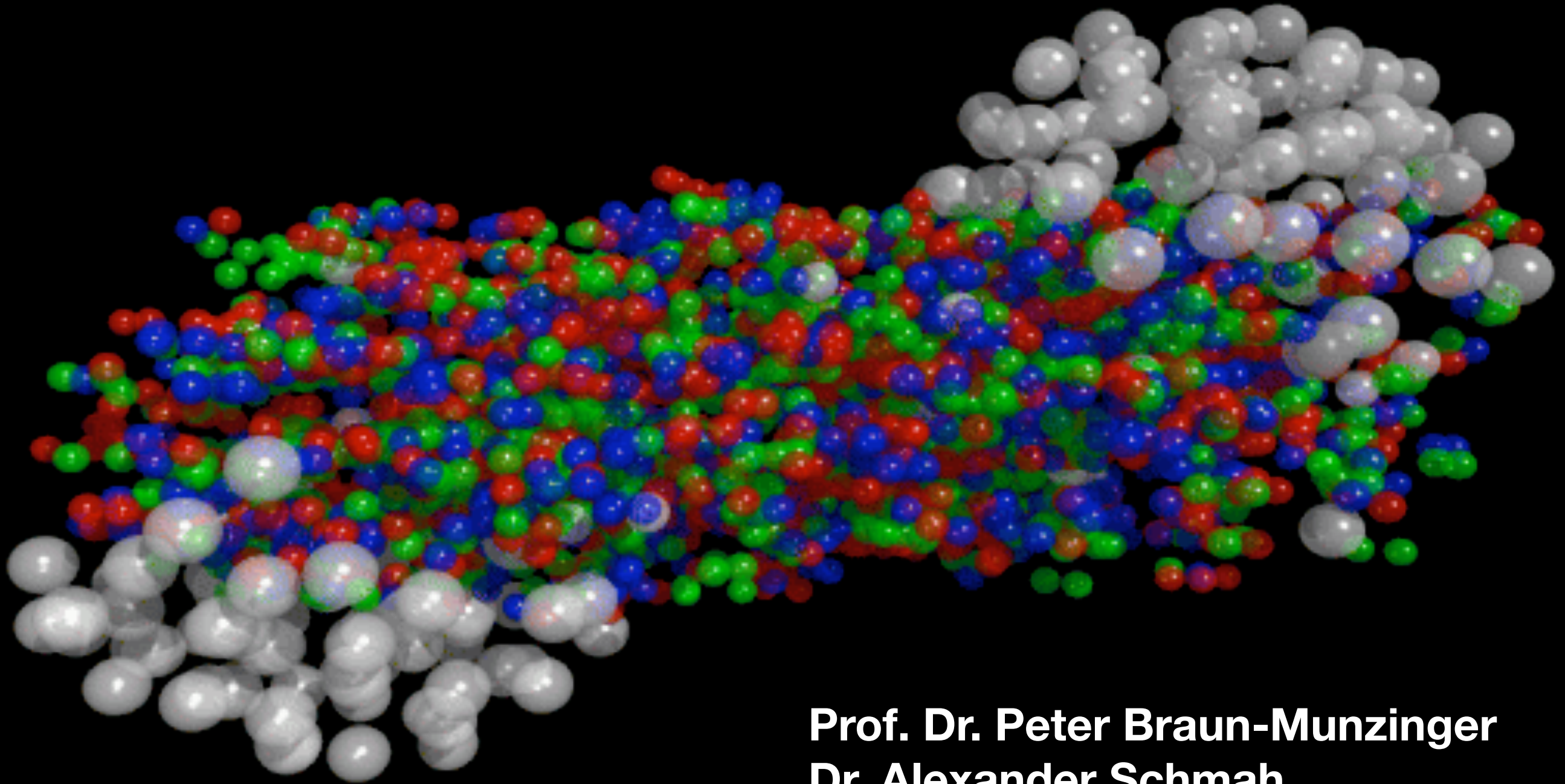


# Quark-Gluon Plasma Physics

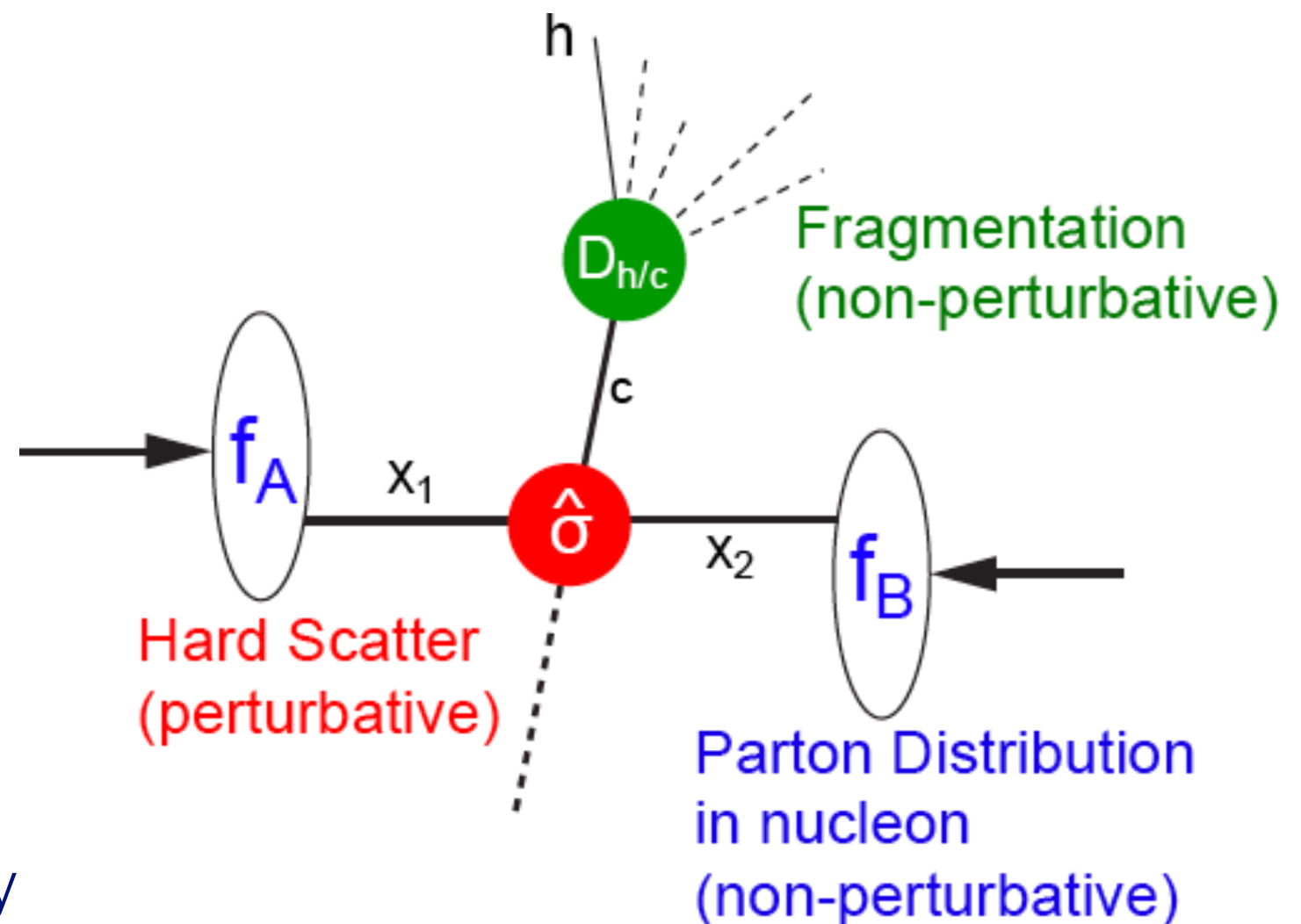
## 9. Hard scattering and Jets



**Prof. Dr. Peter Braun-Munzinger**  
**Dr. Alexander Schmah**  
**Heidelberg University**  
**SS 2021**

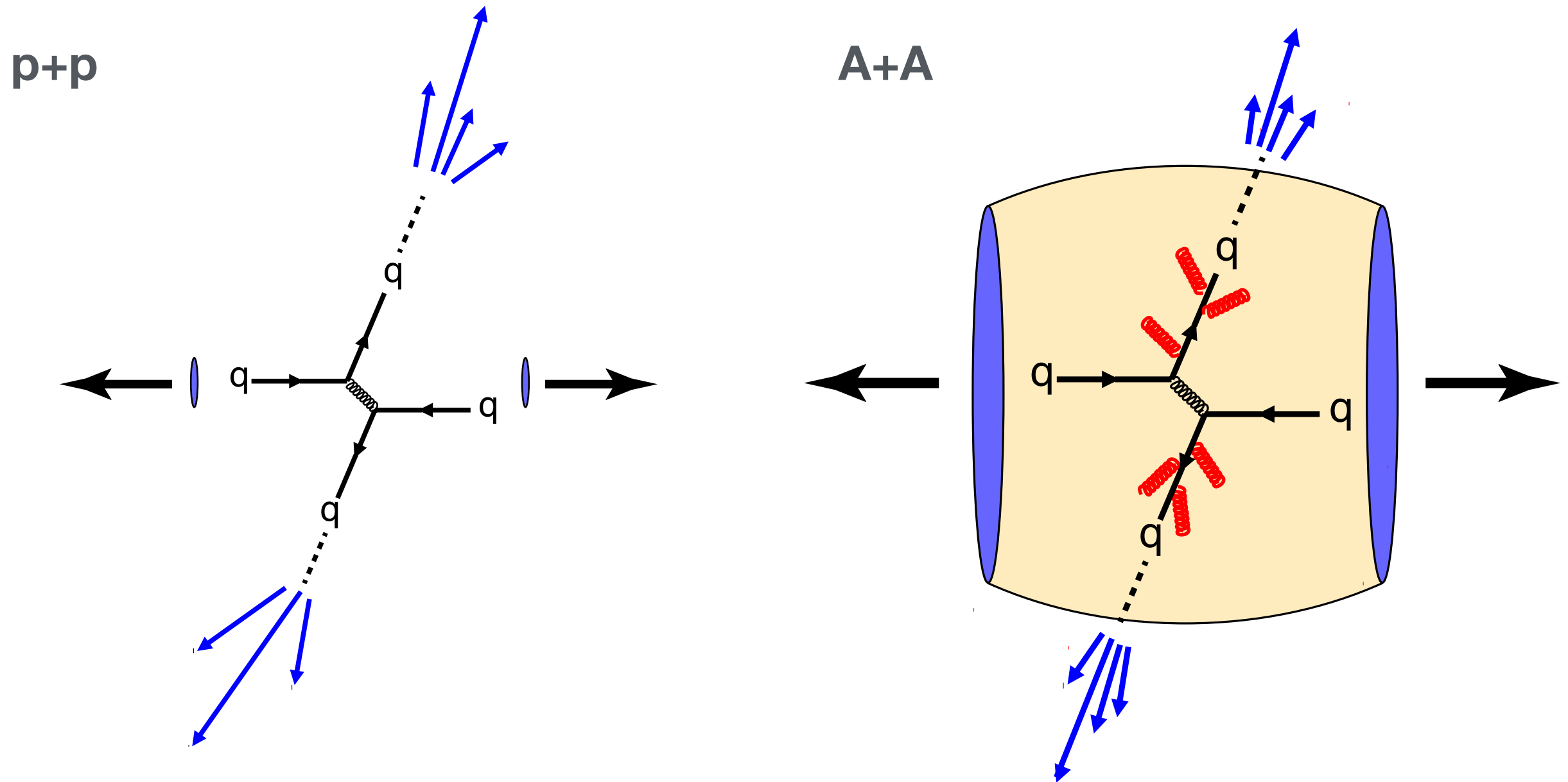
# Theoretical description of High- $p_T$ particle production: Perturbative QCD

- Scattering of pointlike partons described by QCD perturbation theory (pQCD)
- Soft processes described by universal, phenomenological functions
  - ▶ Parton distribution function from deep inelastic scattering
  - ▶ Fragmentation functions from  $e^+e^-$  collisions
- Particle production dominated by hard scattering for  $p_T \gtrsim 3 \text{ GeV}/c$ 
  - ▶ However, 99% or so of all particle from soft processes



$$d\sigma = \sum_{a,b,c} f_a \otimes f_b \otimes d\hat{\sigma}_{ab}^c \otimes D_c^{\text{Hadron}}$$

# Jet quenching in heavy-ion collisions



A-A collision: shower evolution in the medium, energy loss of the leading parton

# Jet quenching history

Energy Loss of Energetic Partons in Quark-Gluon Plasma:  
Possible Extinction of High  $p_T$  Jets in Hadron-Hadron Collisions.

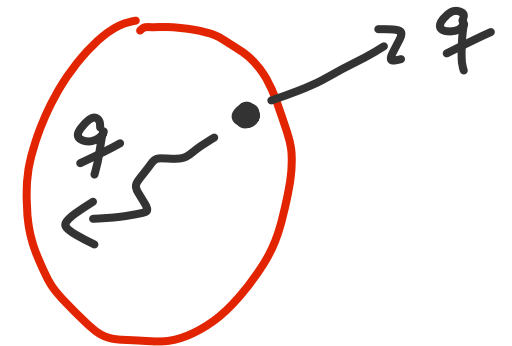
FERMILAB-Pub-82/59-THY  
August, 1982

J. D. BJORKEN  
Fermi National Accelerator Laboratory  
P.O. Box 500, Batavia, Illinois 60510

## Abstract

High energy quarks and gluons propagating through quark-gluon plasma suffer differential energy loss via elastic scattering from quanta in the plasma. This mechanism is very similar in structure to ionization loss of charged particles in ordinary matter. The  $dE/dx$  is roughly proportional to the square of the plasma temperature. For this effect. An interesting signature may be events in which the hard collision occurs near the edge of the overlap region, with one jet escaping without absorption and the other fully absorbed.

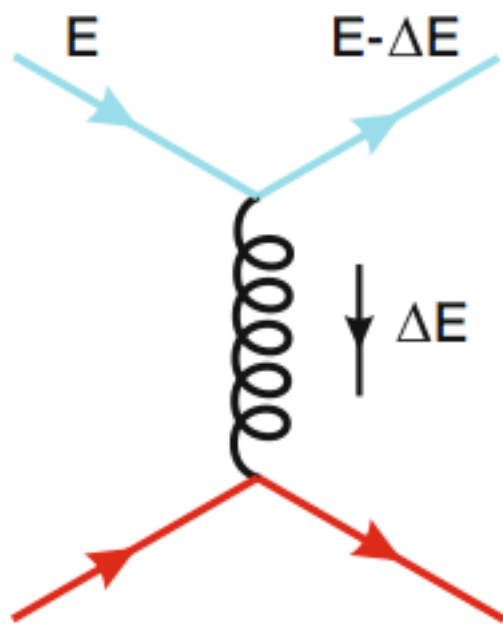
It is now believed that radiative energy loss (gluon bremsstrahlung) is more important than elastic scattering





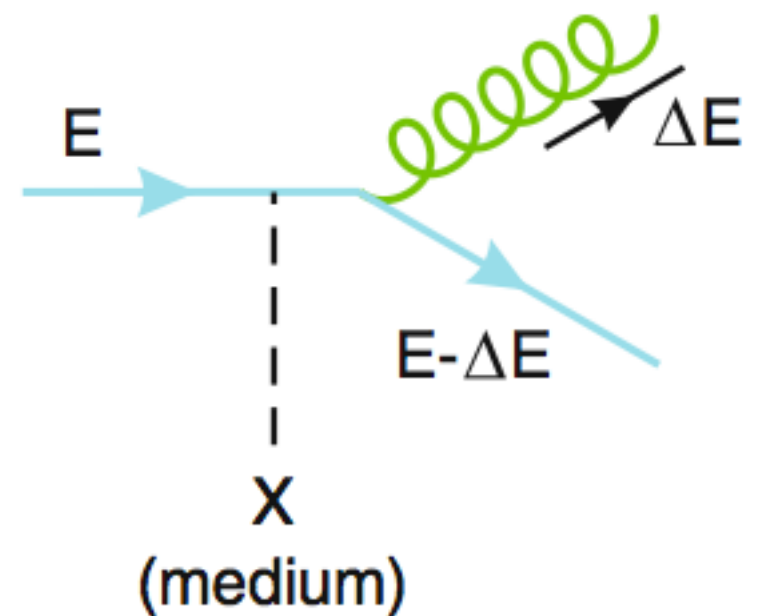
# Collisional vs. radiative parton energy loss

Collisional energy loss:



- Elastic scatterings with medium constituents
- Dominates at low parton momenta

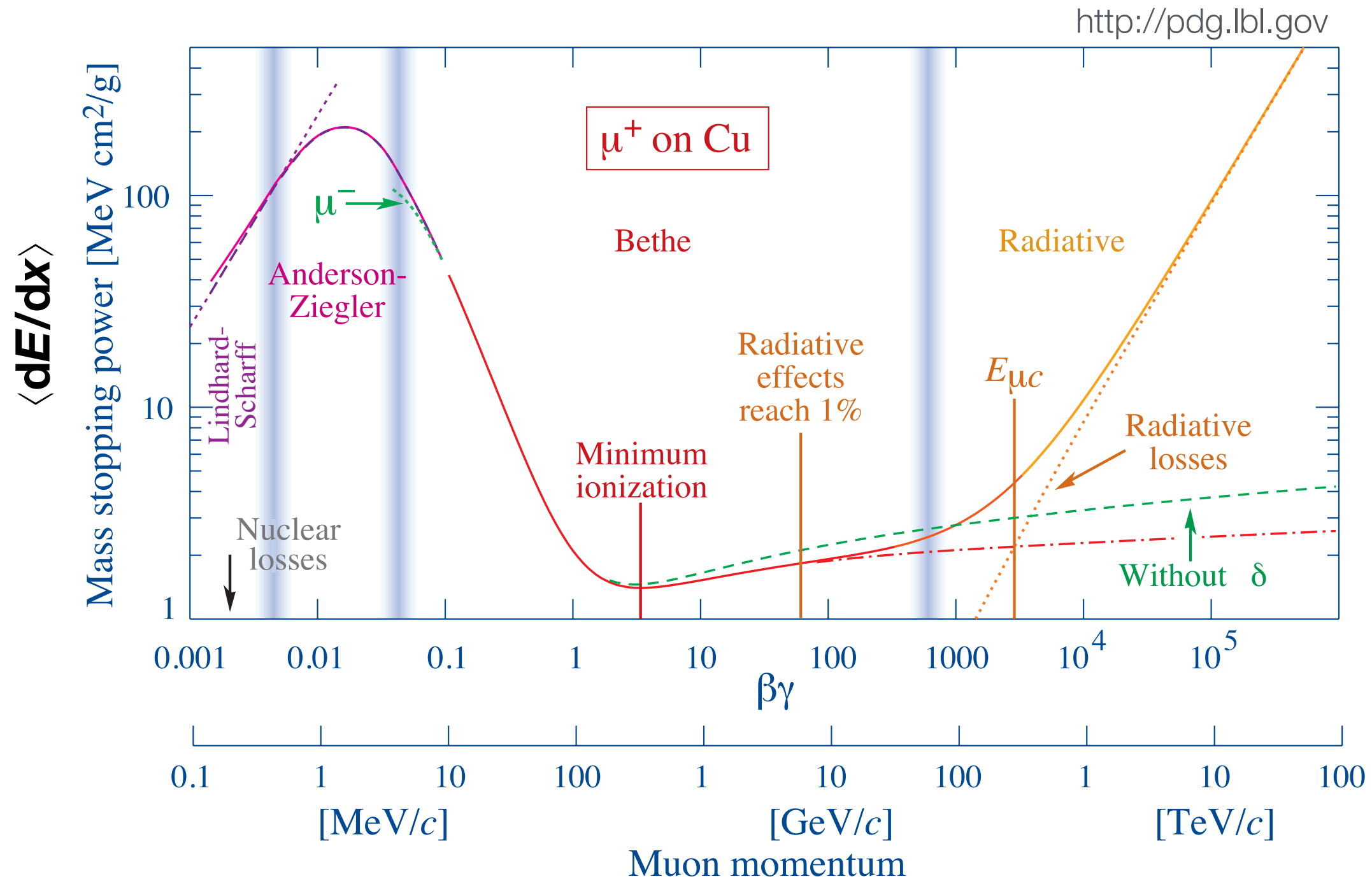
Radiative energy loss:



- Inelastic scatterings within the medium
- Dominates at higher momenta

# Analogy:

## Energy loss of charged particles in ordinary matter



$\mu^+$  on Cu: Radiational energy loss ("bremsstrahlung") starts to dominate over collisional energy loss ("Bethe-Bloch") for  $p \gg 100$  GeV/c

# Basics of radiative parton energy loss (1)

- Energy loss  $E$  in a static medium of length  $L$  for a parton energy  $E \rightarrow \infty$ :

$$\Delta E \propto \alpha_s C_F \hat{q} L^2 \quad \hat{q} = \frac{\mu^2}{\lambda} \quad C_F = \begin{cases} 3 & \text{for gluon jets} \\ 4/3 & \text{for quark jets} \end{cases} \quad \begin{matrix} \text{Colour} \\ \text{factors} \end{matrix}$$

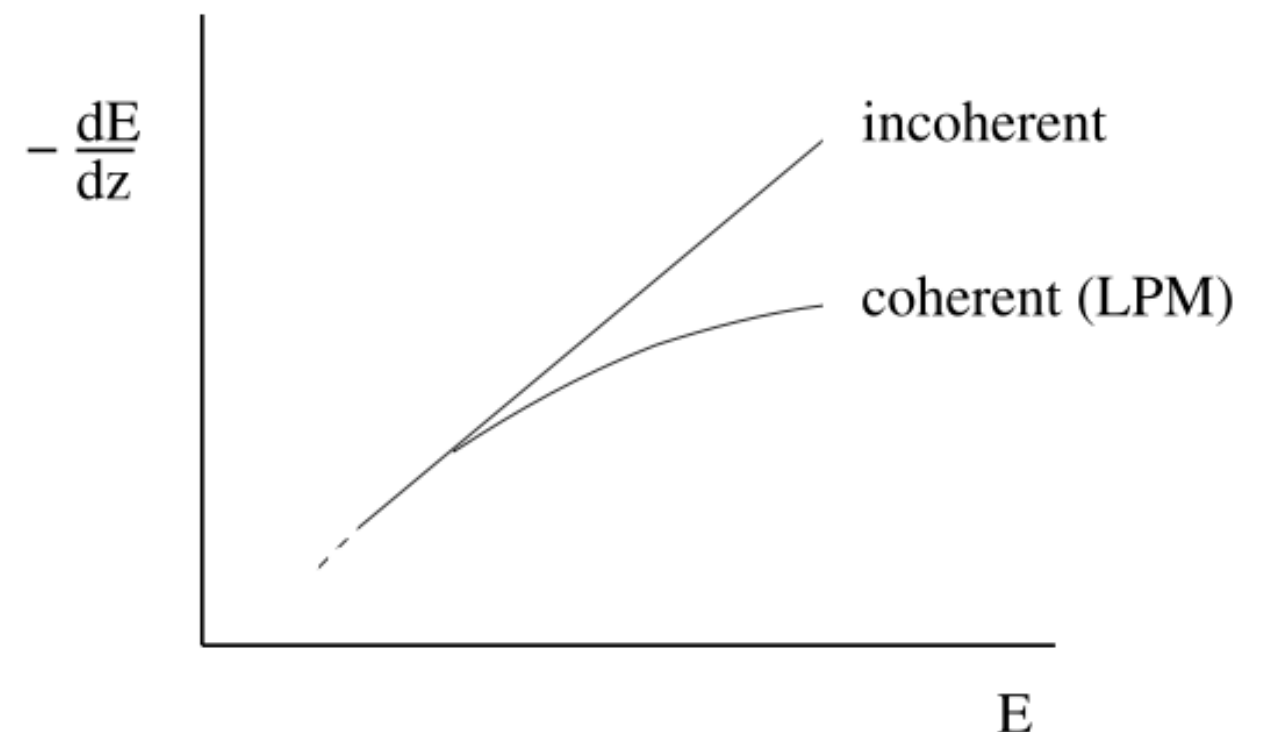
$\mu^2$  : typical momentum transfer from medium to parton per collision

$\lambda$  : mean free path length in the medium

BDMPS result, Nucl. Phys. B 483, 291, 1997

- Landau-Pomeranchuk-Migdal (LPM) effect

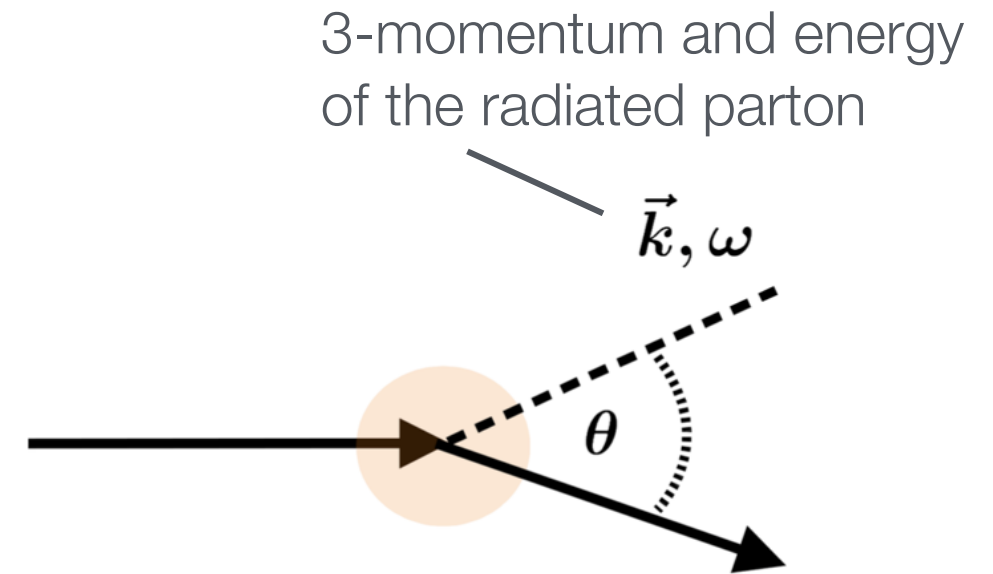
- ▶ Parton scatters coherently off many medium constituents: destructive interference
- ▶ Reduces radiative energy loss



# Basics of radiative parton energy loss (2)

Formation time (or length) of a radiated gluon:  
("time for the fast parton to get rid of its virtuality")

$$z_{\text{coh}} = t_{\text{coh}} \simeq \frac{\omega}{k_T^2} \simeq \frac{1}{\omega \theta^2}$$



The gluon acquires additional transverse momentum if it scatters with medium constituents within its formation time (or formation length  $z_c$ ):

$$k_T^2 \simeq \hat{q} z_{\text{coh}} = \frac{\mu^2}{\lambda} z_{\text{coh}}$$

This results in a medium-modified formation length:

$$z_{\text{coh}} \simeq \frac{\omega}{k_T^2} \simeq \frac{\omega}{\hat{q} z_{\text{coh}}} \rightsquigarrow z_{\text{coh}} \simeq \sqrt{\frac{\omega}{\hat{q}}}$$

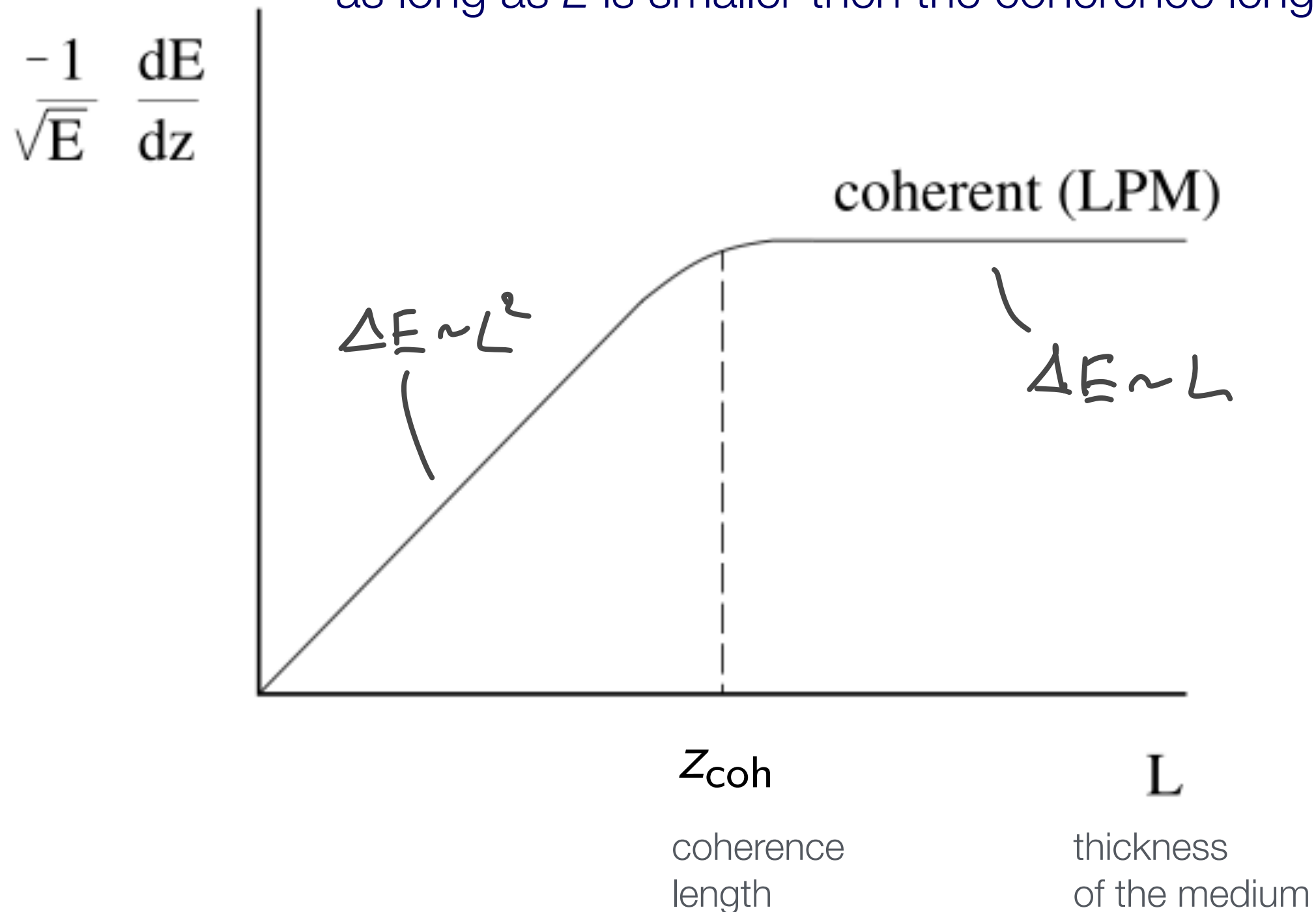
$\lambda < z_{\text{coh}}$  : Coherent scattering with destructive interference

$\lambda > z_{\text{coh}}$  : incoherent multiple scattering

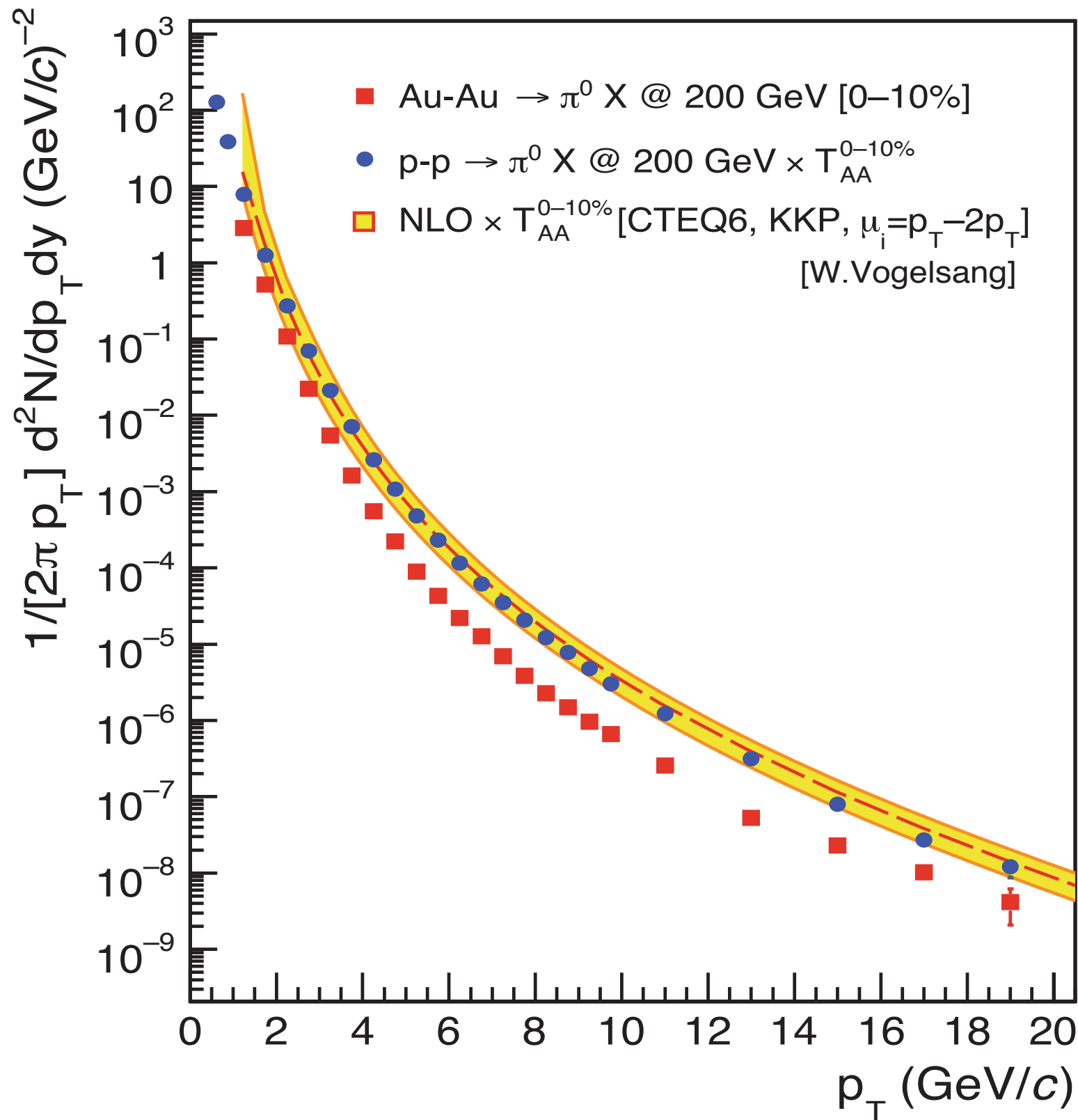


# Basics of radiative parton energy loss (3)

$dE/dz$  increases linearly with medium thickness  $L$  as long as  $L$  is smaller than the coherence length  $z_{\text{coh}}$



# $N_{\text{coll}}$ -scaled $\pi^0$ yields in pp compared to Au-Au



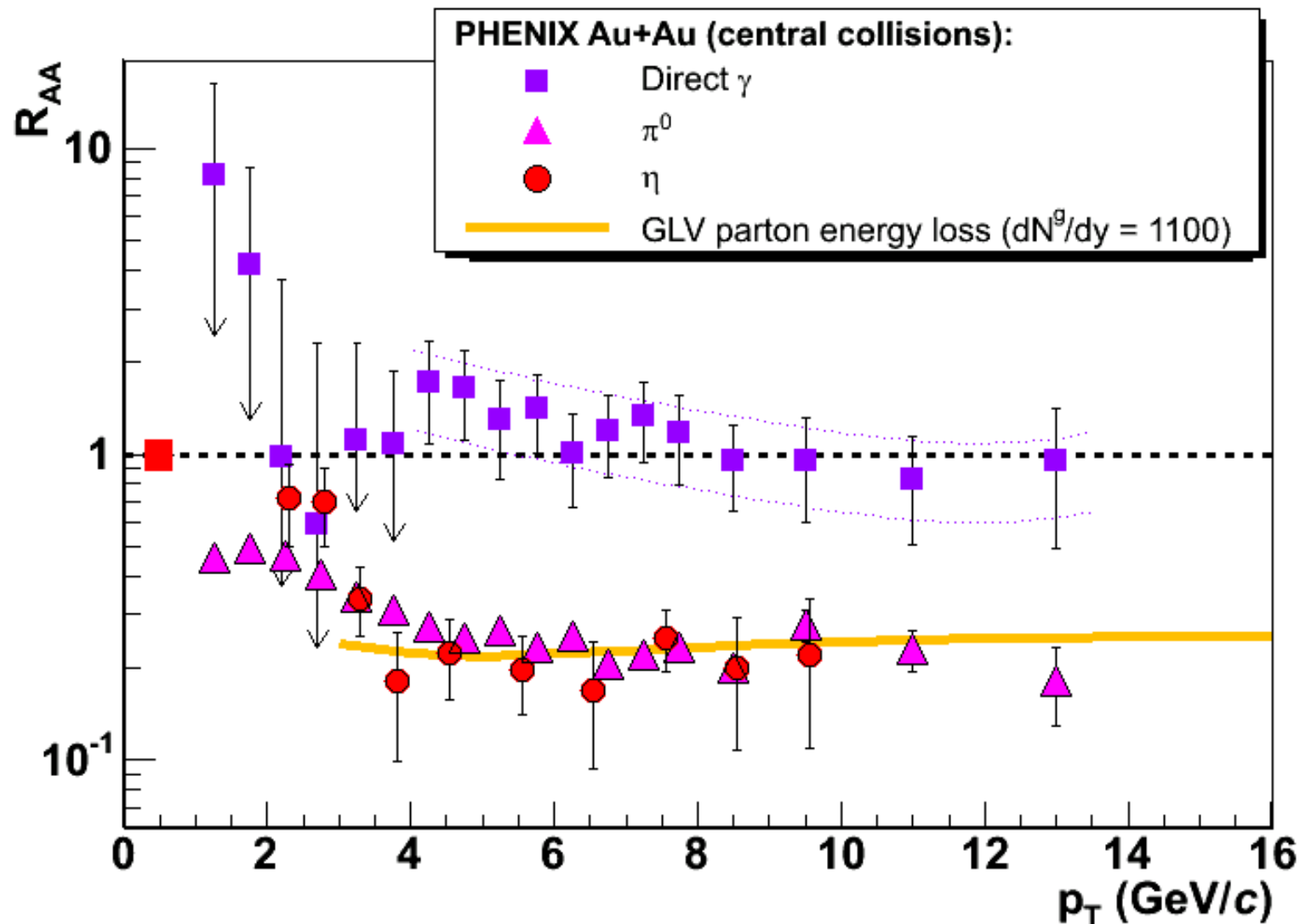
$$T_{AA} = \langle N_{\text{coll}} \rangle / \sigma_{\text{inel}}^{\text{NN}}$$

“increase in parton luminosity”  
per collision when going from  
pp to AA”

Without a medium,  
hadron yields for  
 $p_T \gtrsim 2\text{-}3$  GeV are  
expected to scale  
with  $N_{\text{coll}}$

Observation:  
Clear suppression w.r.t.  
 $N_{\text{coll}}$  scaling

# Discovery of Jet Quenching at RHIC (ca. 2000–2003)



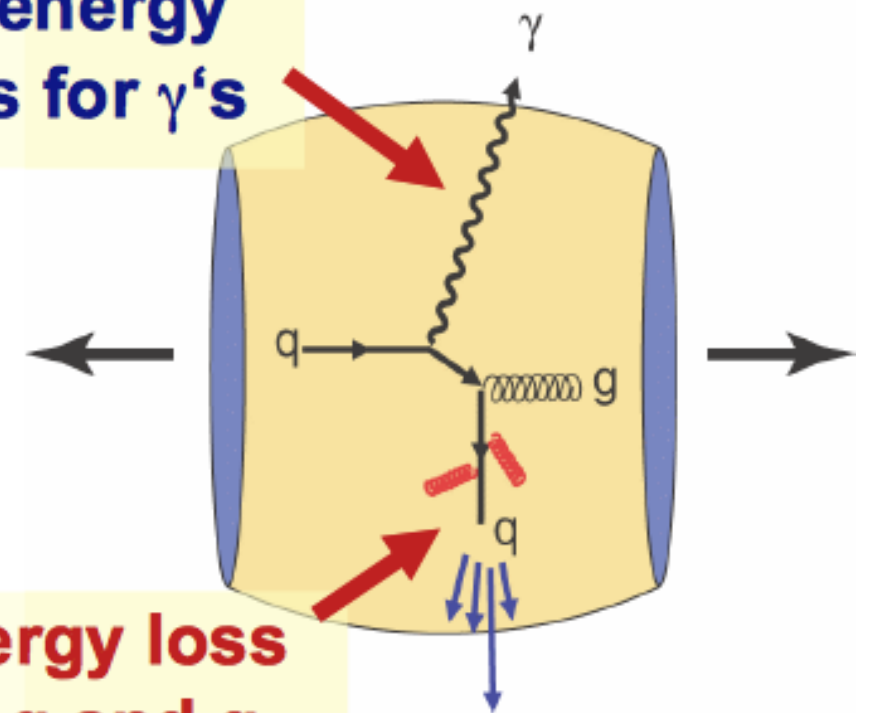
- Hadrons are suppressed, direct photons are not
- Evidence for parton energy loss

$$R_{AB} = \frac{dN/dp_T|_{A+B}}{\langle T_{AB} \rangle \times d\sigma_{inv}/dp_T|_{p+p}},$$

where  $\langle T_{AB} \rangle = \langle N_{coll} \rangle / \sigma_{inel}^{NN}$

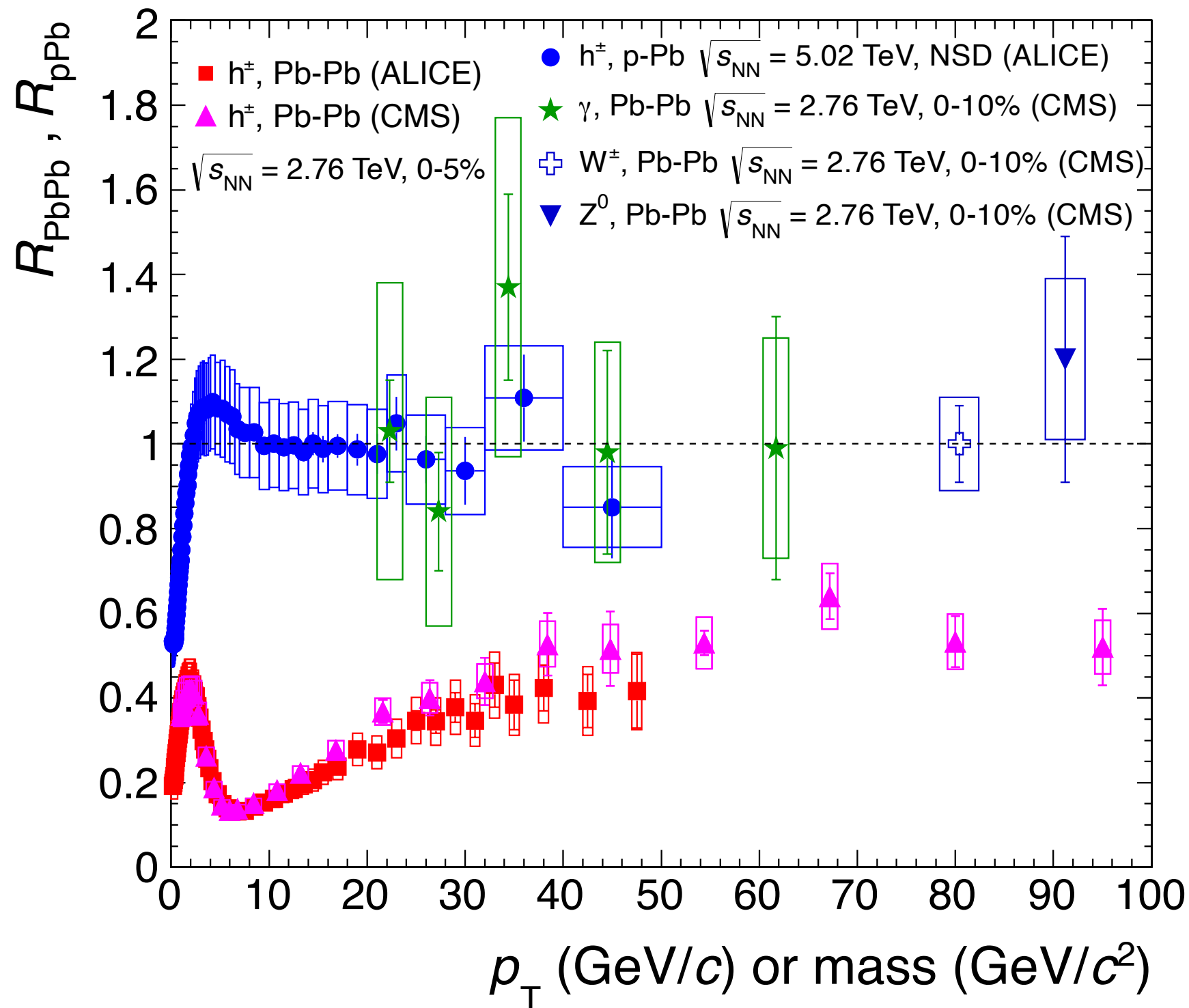
**No energy loss for  $\gamma$ 's**

**energy loss for q and g**



# Single-particle $R_{AA}$ in Pb-Pb at the LHC:

## Qualitatively similar observation as for RHIC energies



- No suppression for  $\gamma$ ,  $W^\pm$ ,  $Z^0$  in Pb-Pb
- No suppression of hadrons in p-Pb
- Strong suppression of hadrons in Pb-Pb

# Medium properties from charged hadron $R_{AA}(p_T)$

- Fit of various models to  $R_{AA}(p_T)$  at RHIC and the LHC
- Jet transport parameter for radiative energy loss at the highest temperatures reached (for  $E_{\text{parton}} = 10$  GeV, QGP thermalization at  $\tau_0 = 0.6$  fm/c):

$$\frac{\hat{q}}{T^3} \approx \begin{cases} 4.6 \pm 1.2 & \text{at RHIC,} \\ 3.7 \pm 1.4 & \text{at LHC,} \end{cases}$$

$$\hat{q} \approx \begin{cases} 1.2 \pm 0.3 \\ 1.9 \pm 0.7 \end{cases} \text{ GeV}^2/\text{fm} \text{ at } \begin{cases} T=370 \text{ MeV,} \\ T=470 \text{ MeV,} \end{cases}$$

Jet Coll., Phys.Rev. C90 (2014) 014909

- Result relies on standard hydro description of the medium evolution
- Conjectured relation to  $\eta/s$ :

$$\frac{T^3}{\hat{q}} = K \frac{\eta}{s}$$

weakly-coupled QGP:  $K \approx 1$

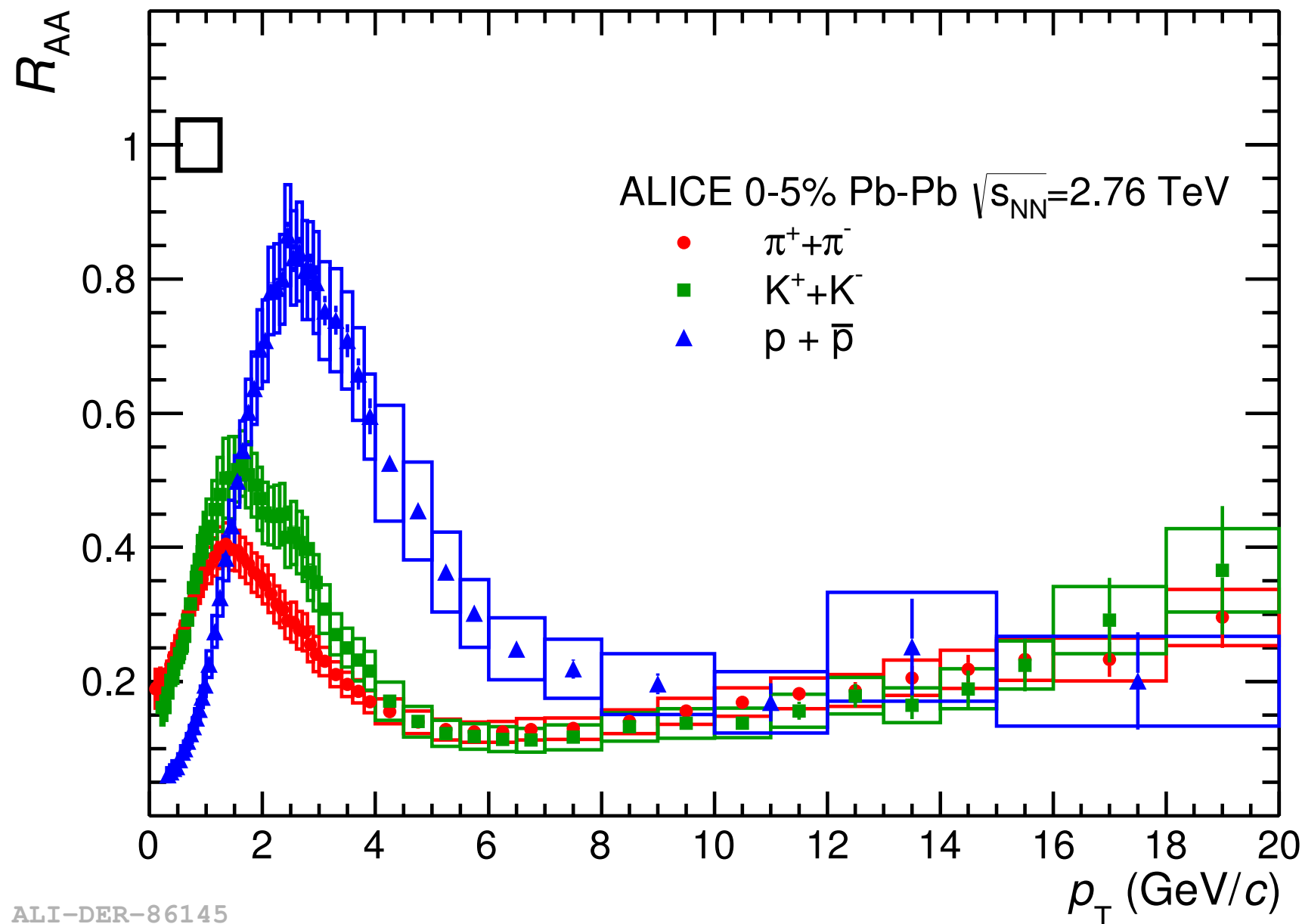
strongly-coupled QGP:  $K \ll 1$

Majumder, Müller, Wang,  
PRL, 99 (2007) 192301



# $\pi, K, p$ $R_{AA}$ : Suppression independent of hadron species for $p_T \gtrsim 8$ GeV/c

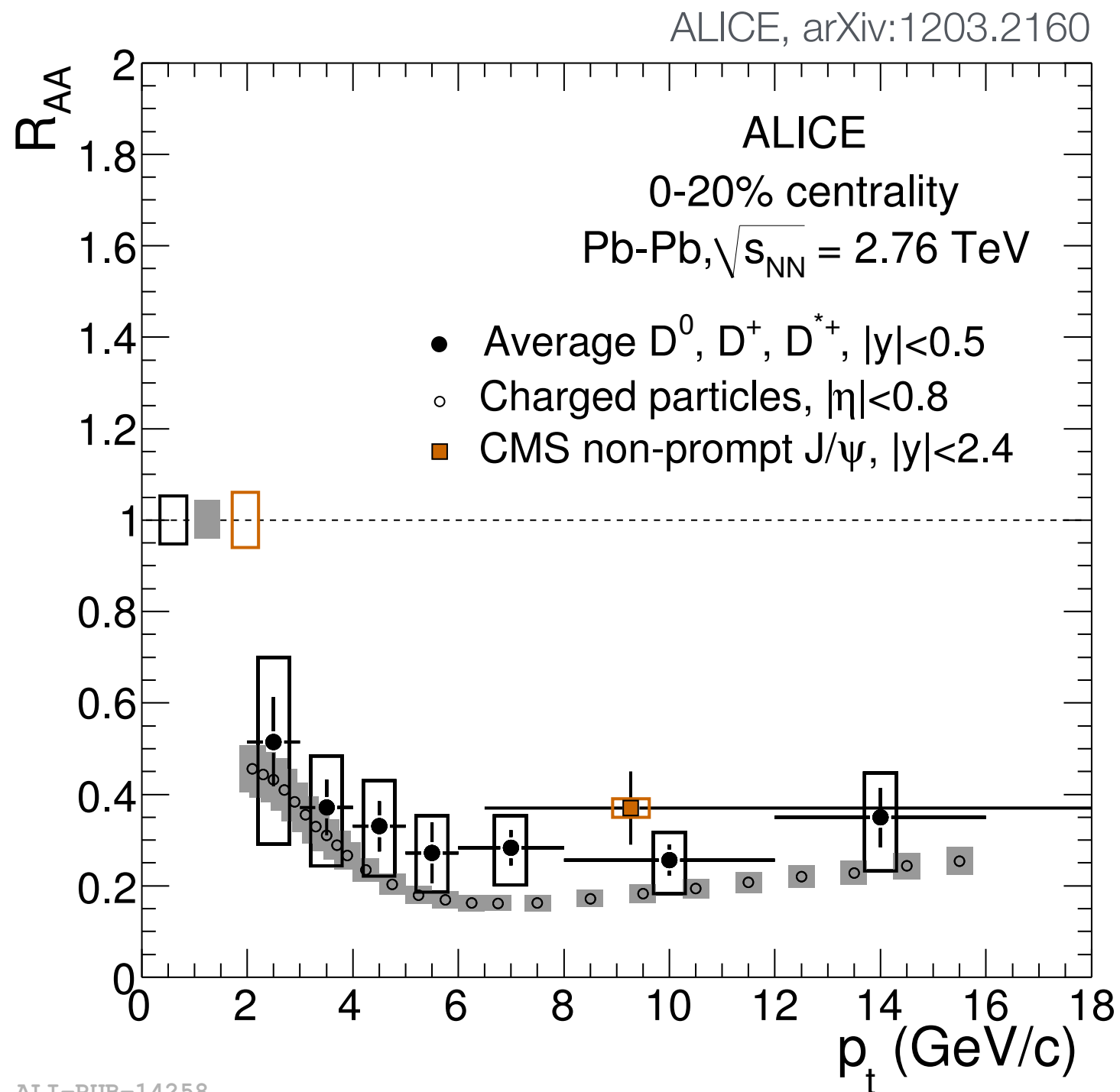
ALICE, arXiv:1506.07287



Leading-parton energy loss followed by fragmentation in QCD vacuum (as in pp) for  $p_{T, \text{hadron}} > 8$  GeV/c?

- $R_{AA}(p) > R_{AA}(K) \approx R_{AA}(\pi)$  for  $3 < p_T < 8$  GeV/c
- Similar  $p, K$  and  $\pi$   $R_{AA}$  for  $p_T > 8$  GeV/c

# D meson $R_{AA}$ : Charm quark energy loss similar to quark and gluon energy loss



ALI-PUB-14258

$$\Delta E_g > \Delta E_{u,d,s} > \Delta E_c > \Delta E_b$$

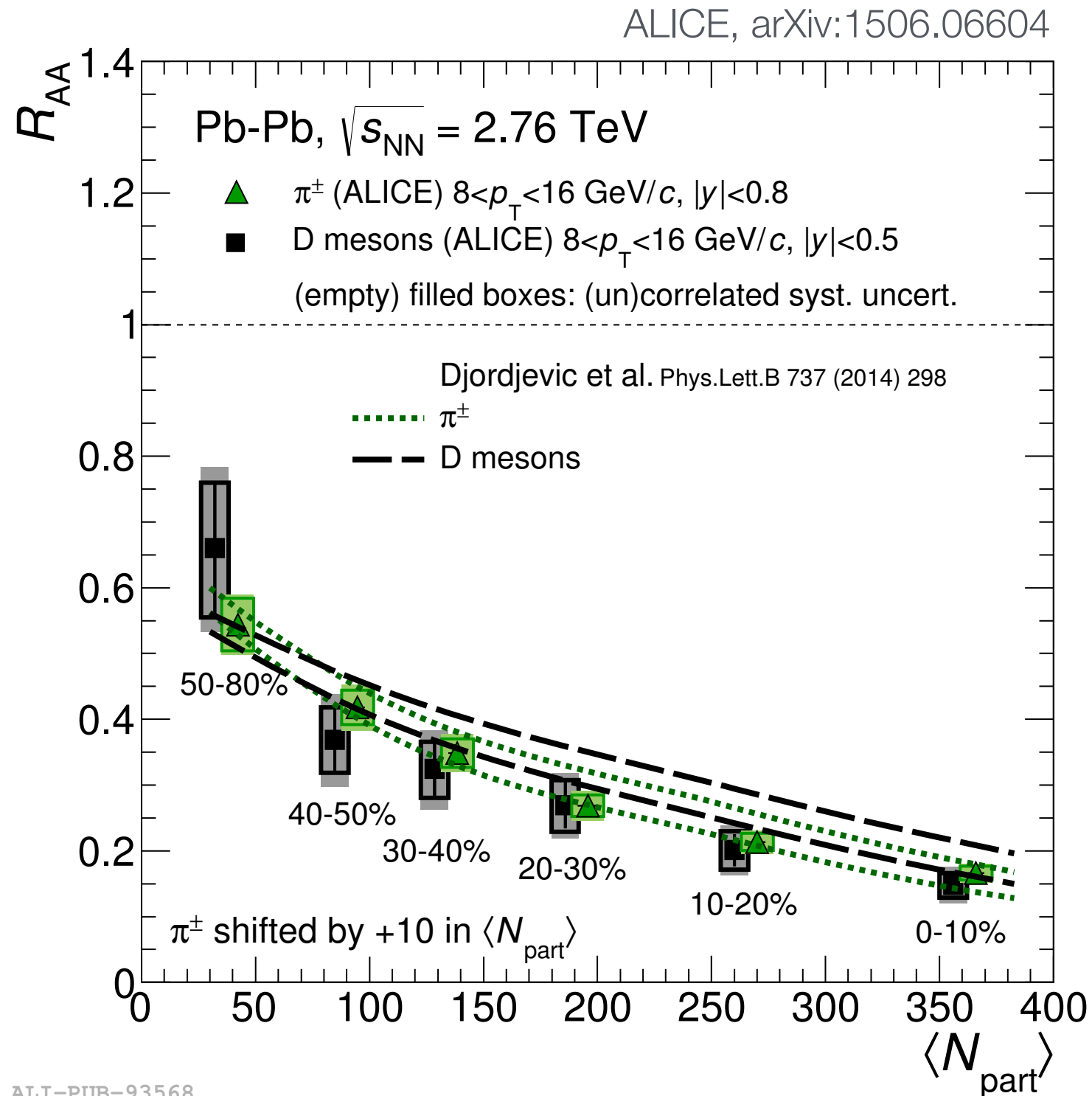
color factor

dead cone effect

(gluon emission suppressed at forward angles for slow quarks)

- Strong suppression also for D mesons (which cannot be explained by shadowing)
- Suppression of D mesons and pions (surprisingly?) similar
  - ▶ Pions mainly from gluons
  - ▶ Dead cone effect for c and b
- Still hint for expected hierarchy?
  - ▶ However, need to carefully consider also the steepness of the initial parton spectra

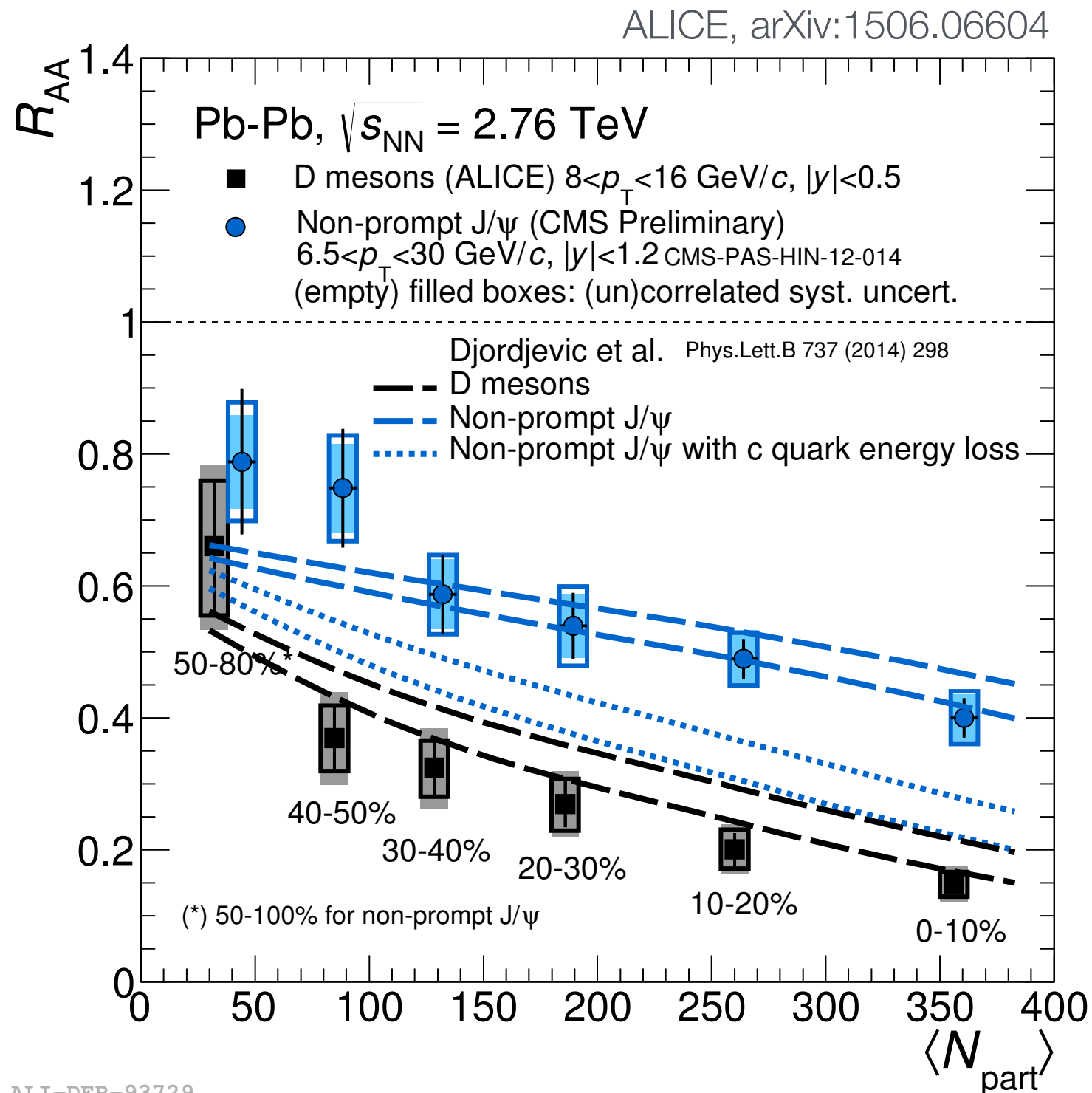
# Centrality dependence of $R_{AA}$ for pions and D mesons



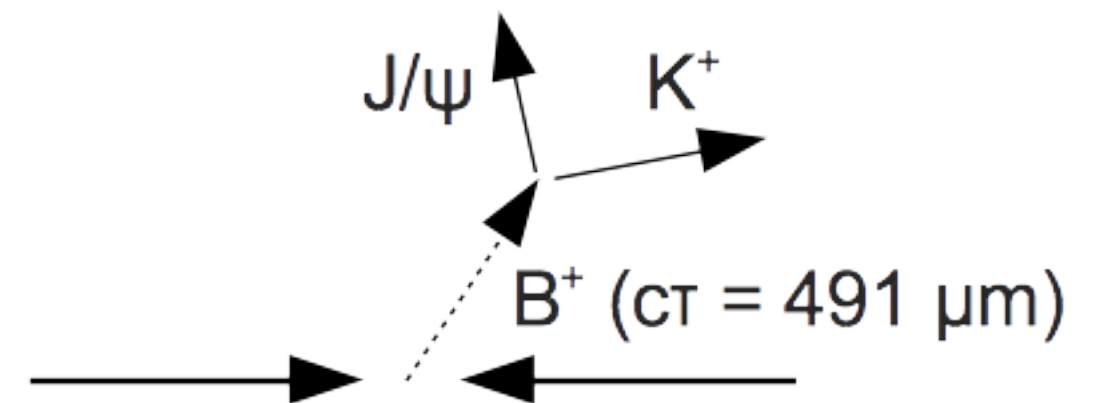
ALI-PUB-93568

- $N_{part}$  dependence similar for pions and D mesons
- Actually expected in parton energy loss calculation

# Evidence for smaller energy loss for b quarks than for c quarks

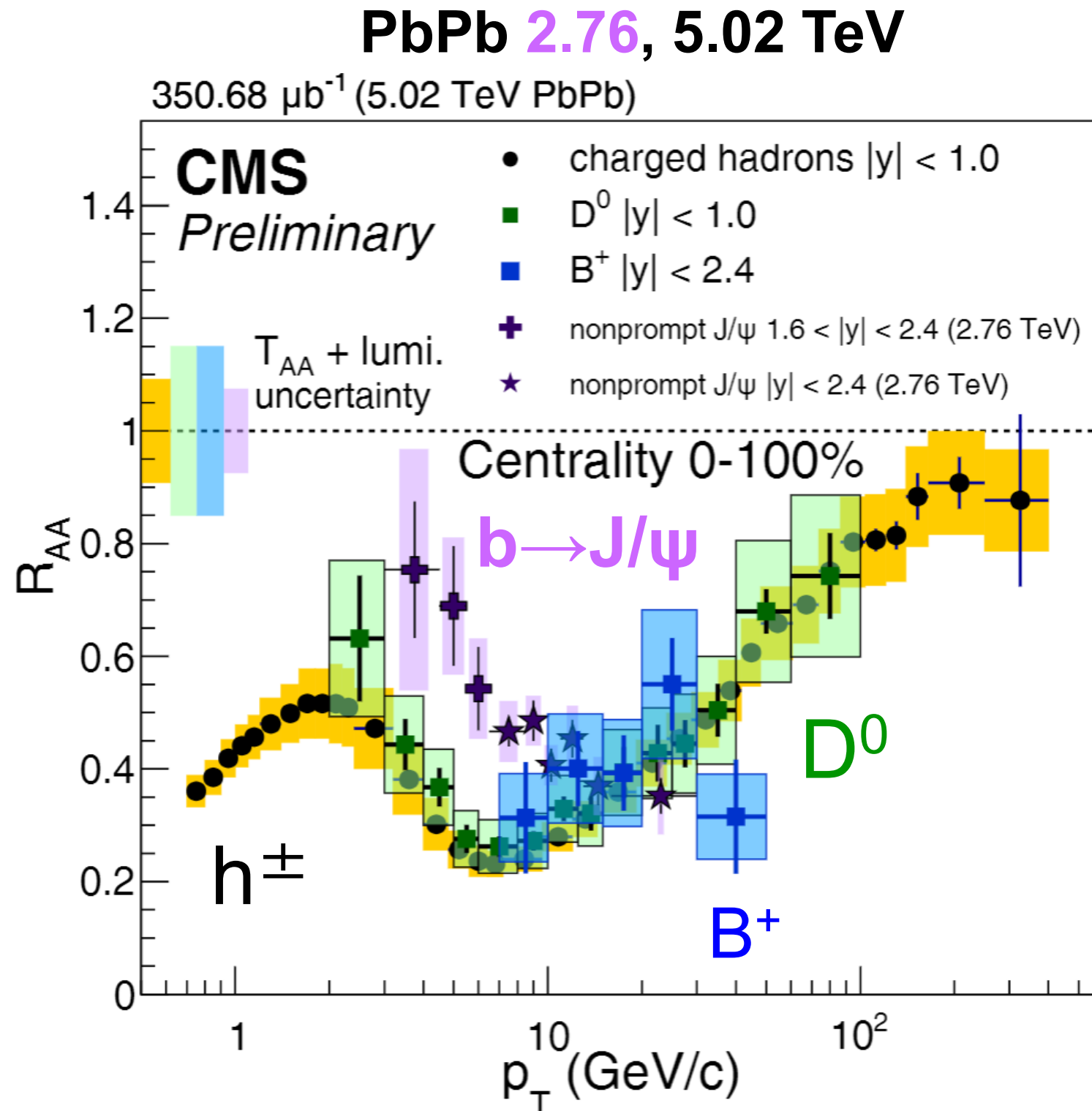


b quark energy loss via non-prompt J/ψ:



Evidence for the expected energy loss hierarchy between c and b quarks

# B-Meson $R_{AA}(p_T)$



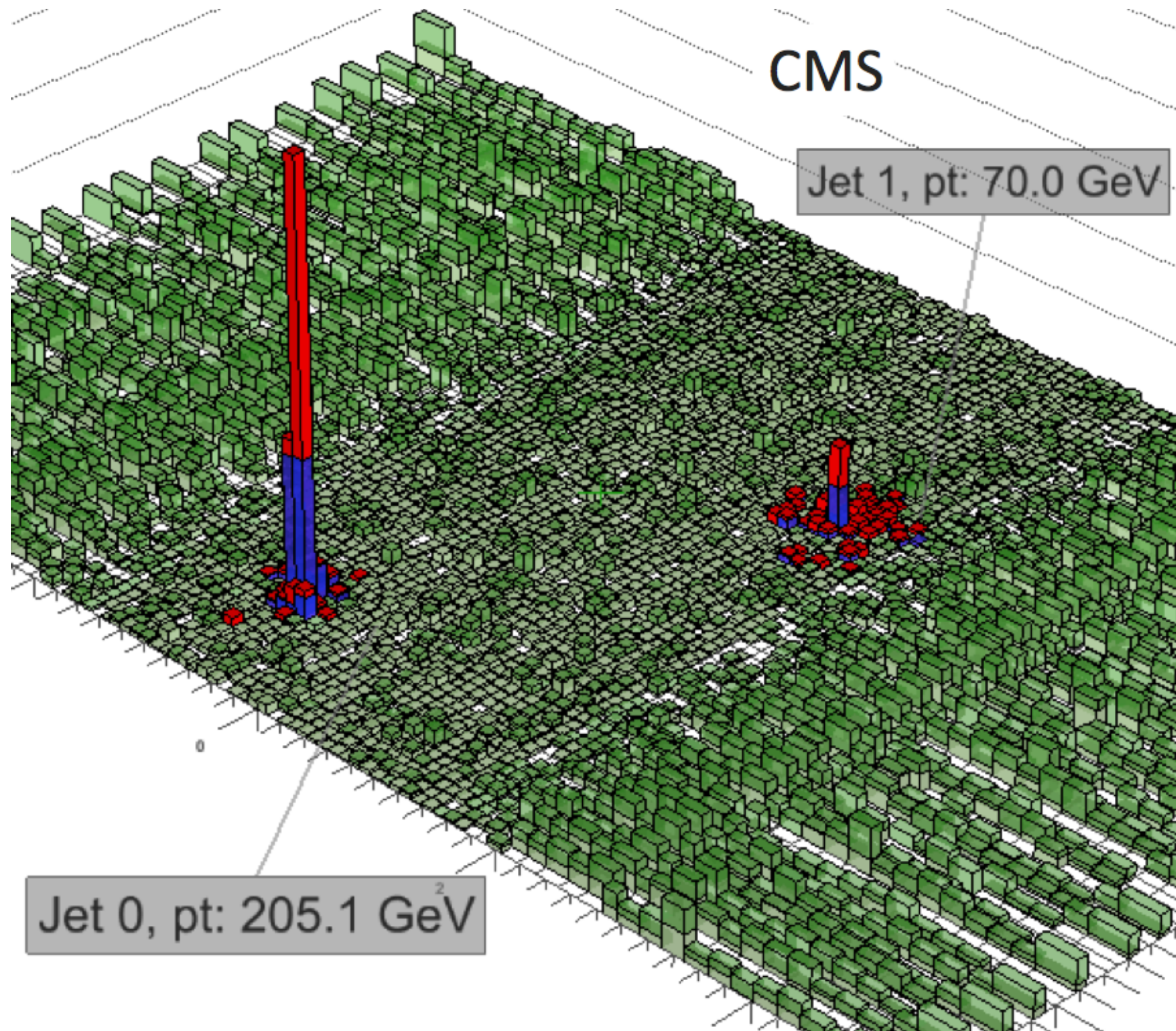
No quark flavor dependence of  $R_{AA}(p_T)$  for  $p_T \gtrsim 10$  GeV/c

Indication for quark flavor dependence at lower  $p_T$  from  $B \rightarrow J/\psi + X$

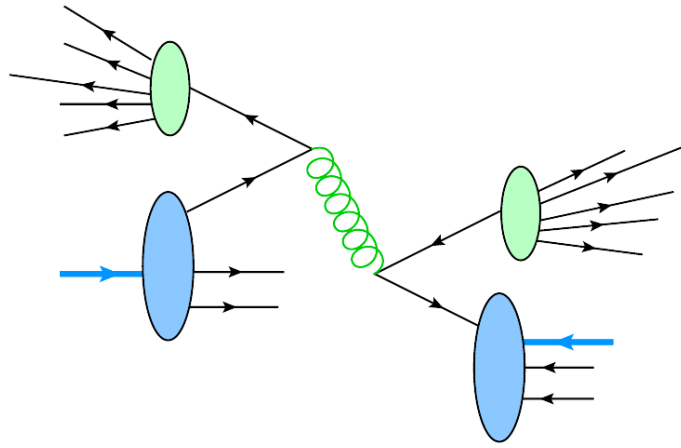
CMS, 1705.04727  
also arXiv:1810.11102



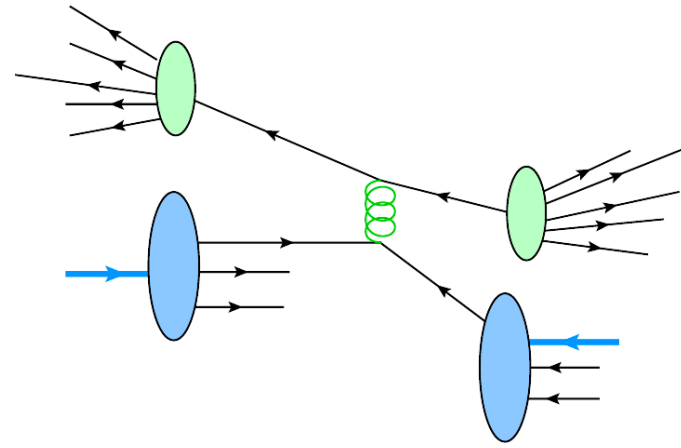
# Studying jet quenching with jets: Large dijet energy asymmetries in Pb-Pb



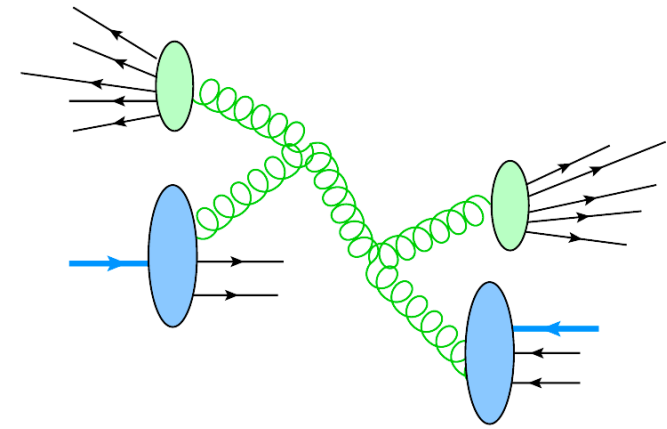
# Jet Production



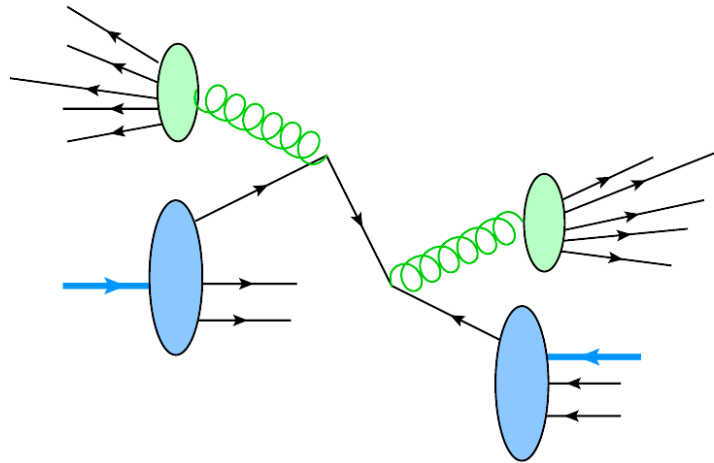
$$q\bar{q} \rightarrow q\bar{q}$$



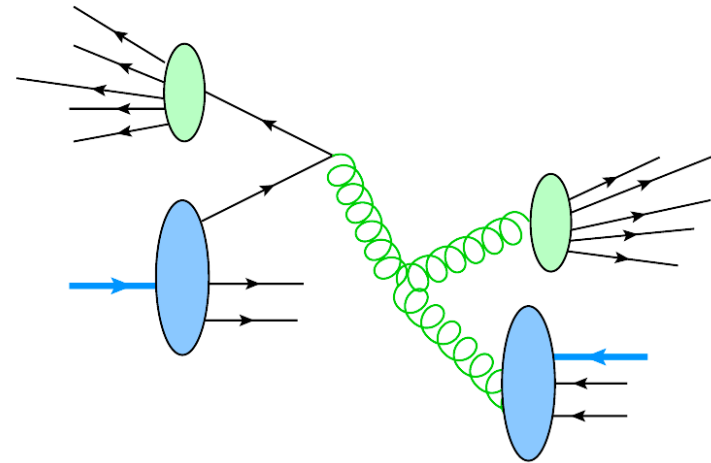
$$q\bar{q} \rightarrow q\bar{q}$$



$$gg \rightarrow gg$$



$$q\bar{q} \rightarrow gg$$



$$qg \rightarrow qg$$

Jet production via high momenta parton-parton collisions, quark-quark up to gluon-gluon



# Quark and Gluon Jets

**Light Quark jet**



$$C_F = \frac{4}{3}$$

**Gluon Jet**

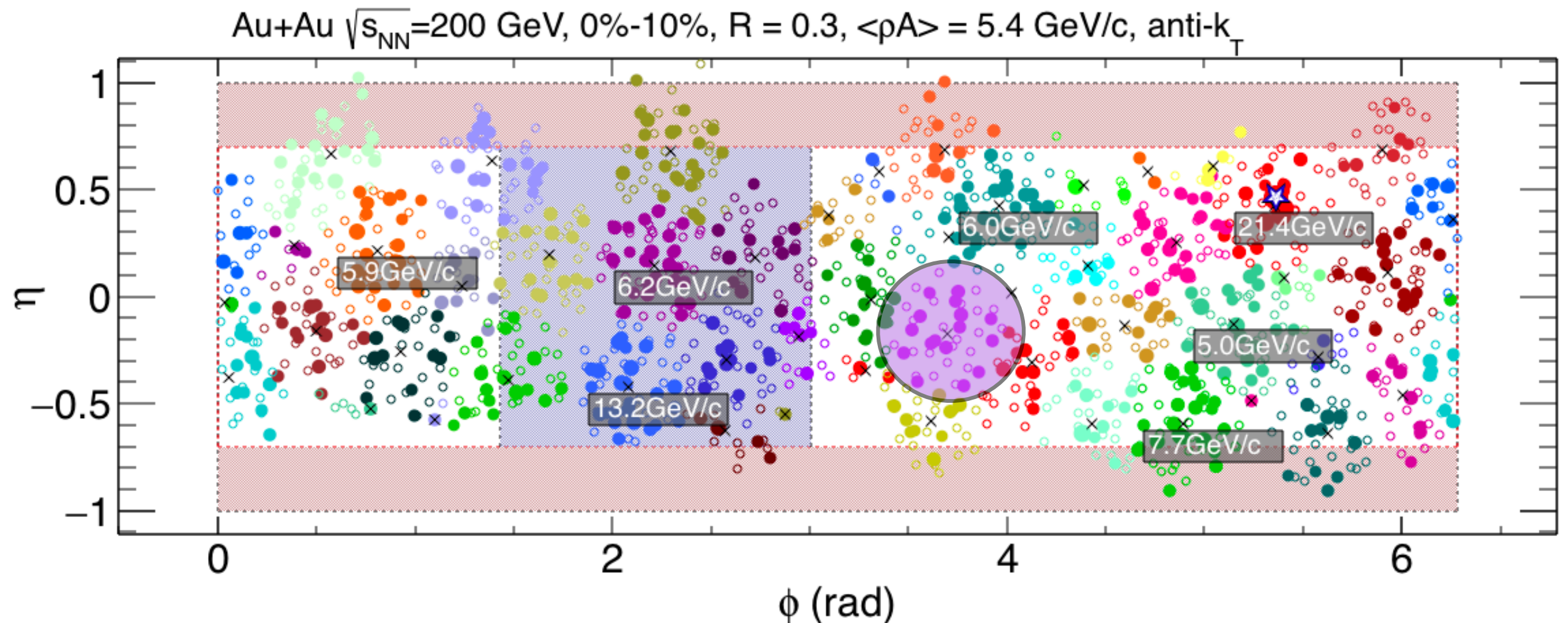


$$C_F = 3$$

Color factor based on differences between quark flavors + colors vs. gluons

Gluon jets have more constituent, are wider and the constituent energies are more uniform

# How to reconstruct jets

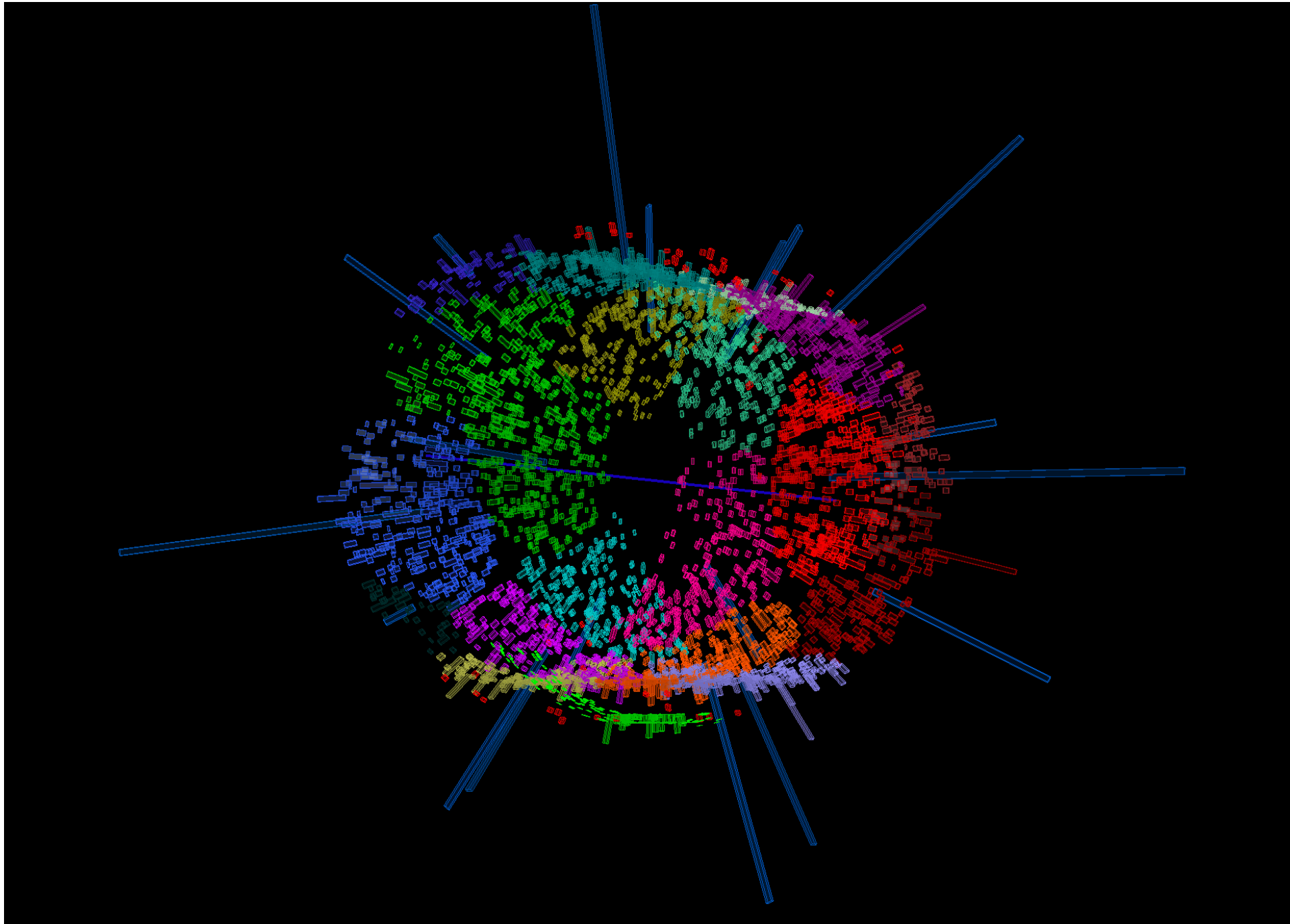


- Anti- $k_T$  algorithm is clustering track around a high  $p_T$  seed particle
- Jet resolution parameter  $R$  defines approximately a radius of the jet:

$$R = \sqrt{\Delta\phi^2 + \Delta\eta^2}$$

- By far not all reconstructed jets are physical jets, most are background

# Jet reconstruction



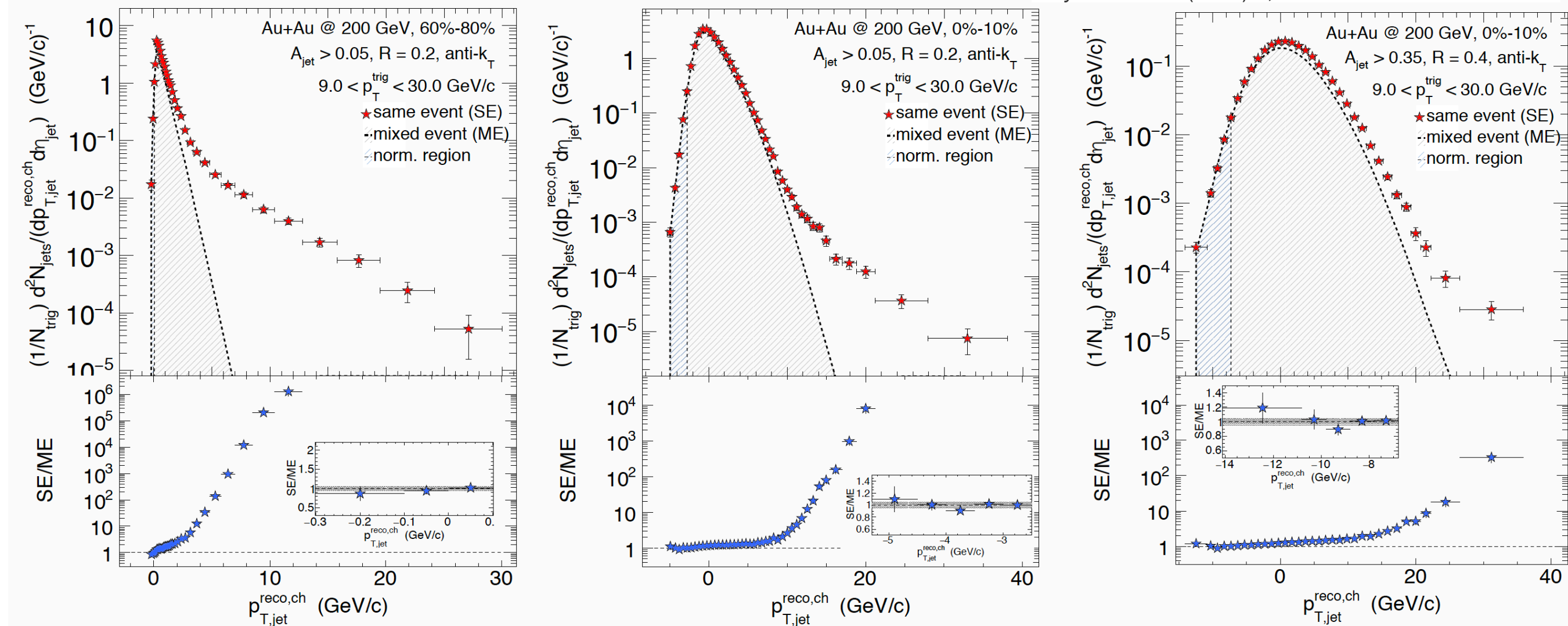
Some important aspects:

- Use of only charged particles in tracking detectors or calorimetry?
- Fiducial selection (acceptance)
- Trigger, e.g. select high  $p_T$  particle in jet



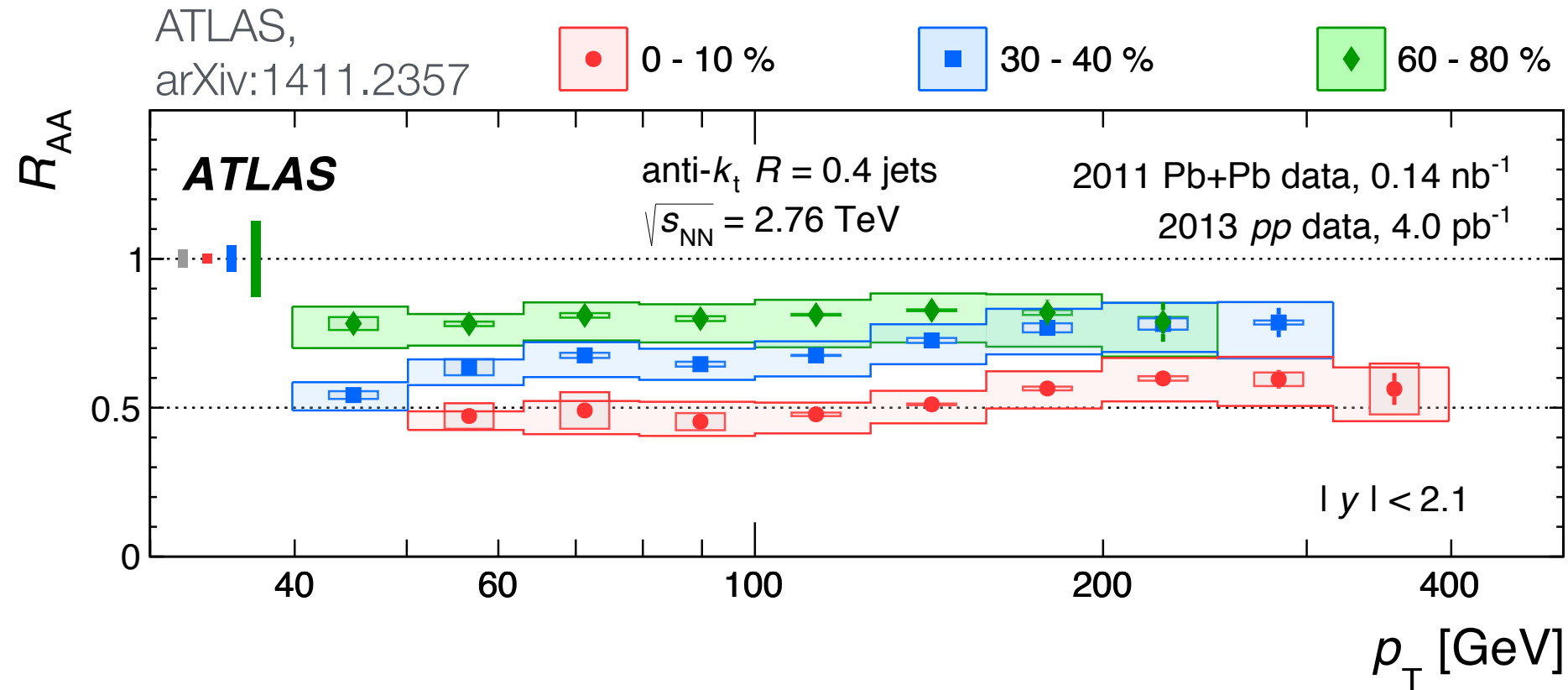
# Jets and background for different R

Phys.Rev.C 96 (2017) 2, 024905 • e-Print: [1702.01108](https://arxiv.org/abs/1702.01108)

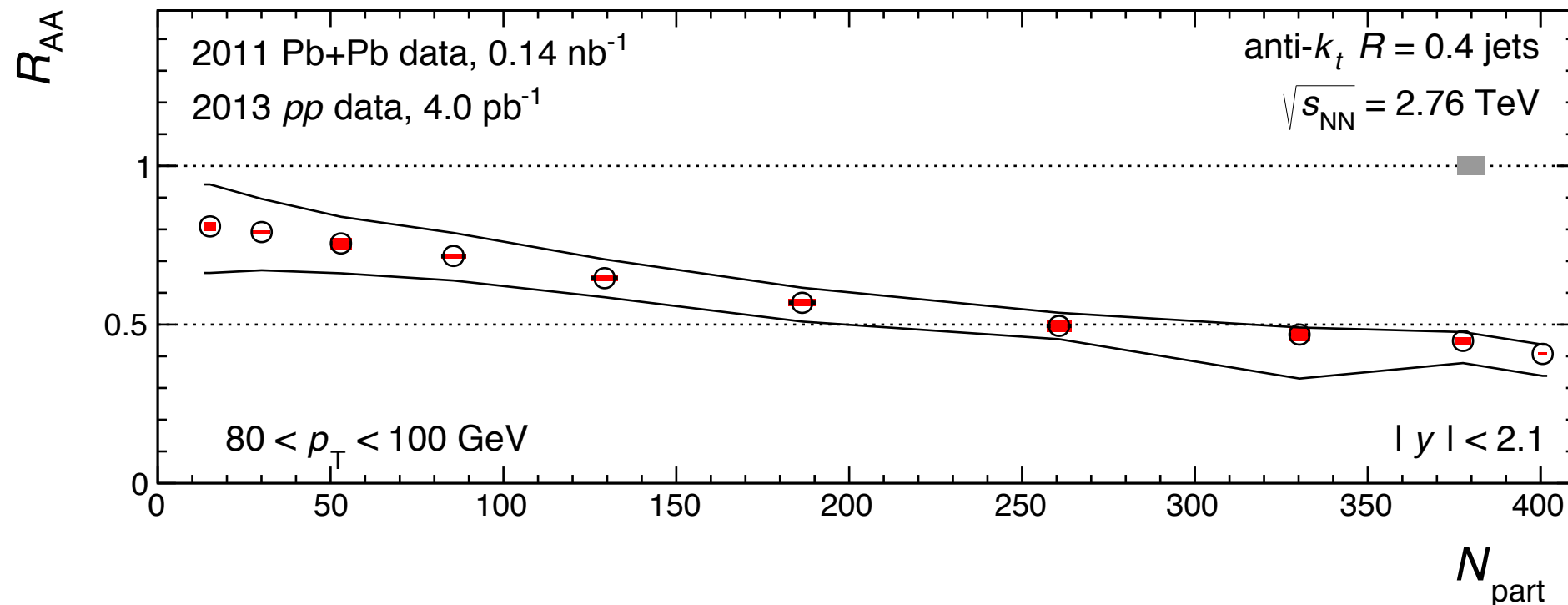


- Combinatorial “fake” jets are very dominant at lower  $p_T$
- Width of the distribution increases with  $R$ , additional particles are picked up

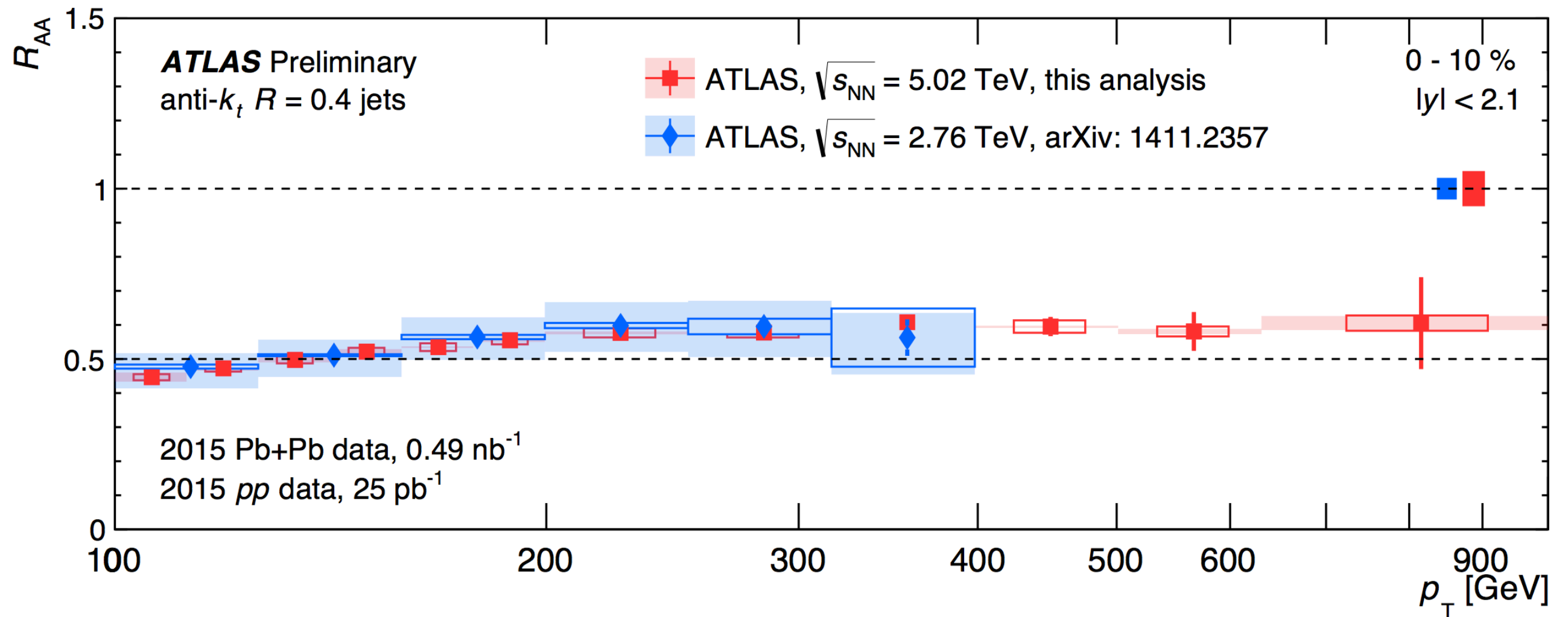
# Jet suppression in Pb-Pb at $\sqrt{s_{NN}} = 2.76$ TeV: $R_{AA} \approx 0.5$ in central collisions



Interestingly, there is not much of a  $p_T$  dependence of the jet suppression



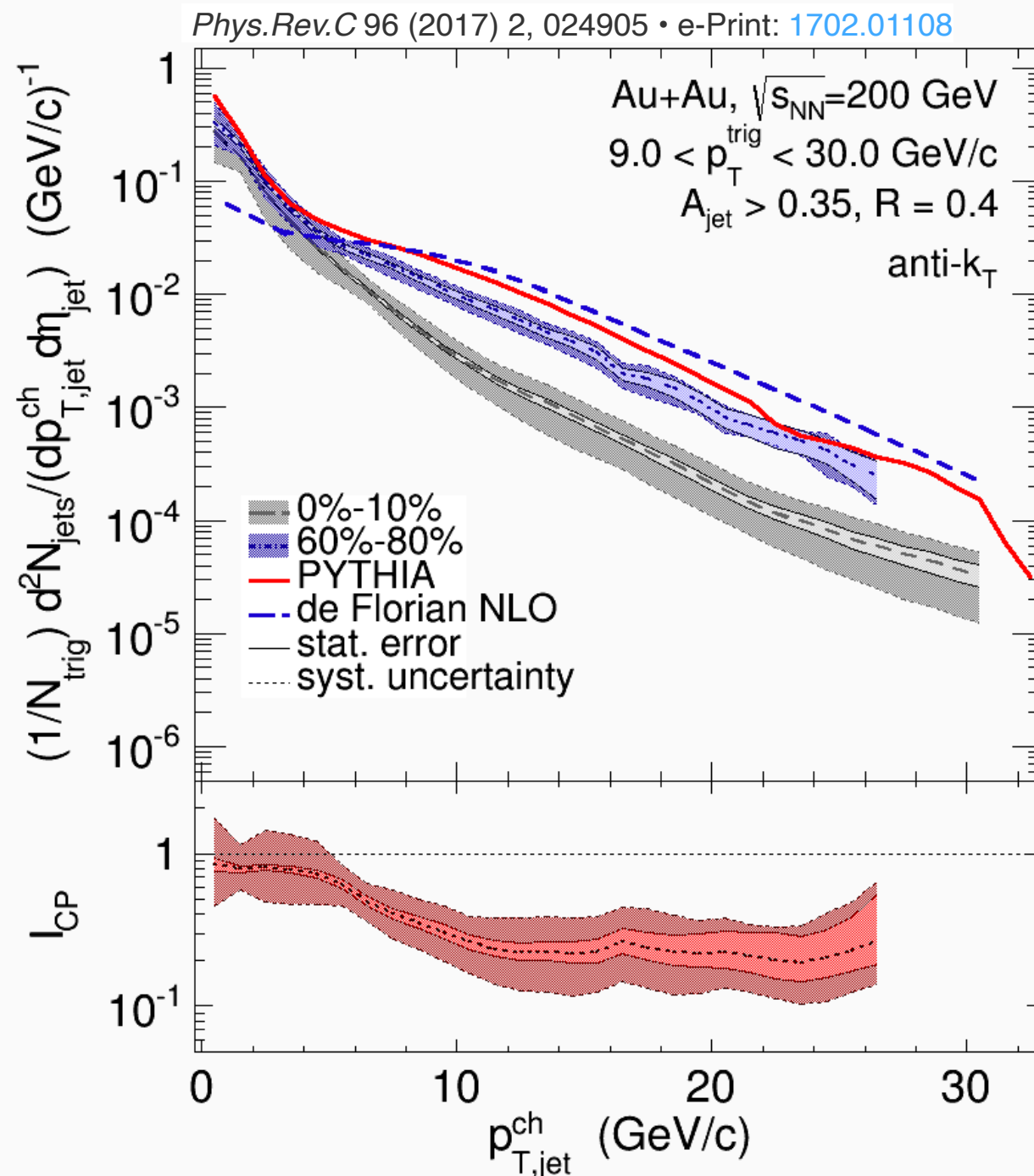
# Pb-Pb at $\sqrt{s_{NN}} = 5.02$ TeV: Jet $R_{AA}$ up to $p_T = 1$ TeV



ATLAS-CONF-2017-009, <http://cds.cern.ch/record/2244820>

Same suppression at the two different  
energies at very high  $p_T$

# Jet suppression down to low $p_T$



For low  $p_T$  the  $I_{CP}$  is close to one at 200 GeV

Next stage is to measure almost unbiased jets to very low  $p_T$

# Summary/questions jet quenching

- Collisional and radiative parton energy loss
  - ▶ Radiative energy loss expected to be dominant for light quarks
- Evidence for expected quark mass dependence of the energy loss
- QCD inspired models are capable of reproducing many features seen in the data
  - ▶ Medium properties can be constrained
- What's next?
  - ▶ First generation of models focussed on leading-particle energy loss (“medium-modified fragmentation function”)
  - ▶ Need to describe full parton shower evolution in the medium
  - ▶ Can one eventually describe parton energy loss based on first principles?
  - ▶ Can one connect heavy-quark energy loss to string theory via the gauge/gravity duality?