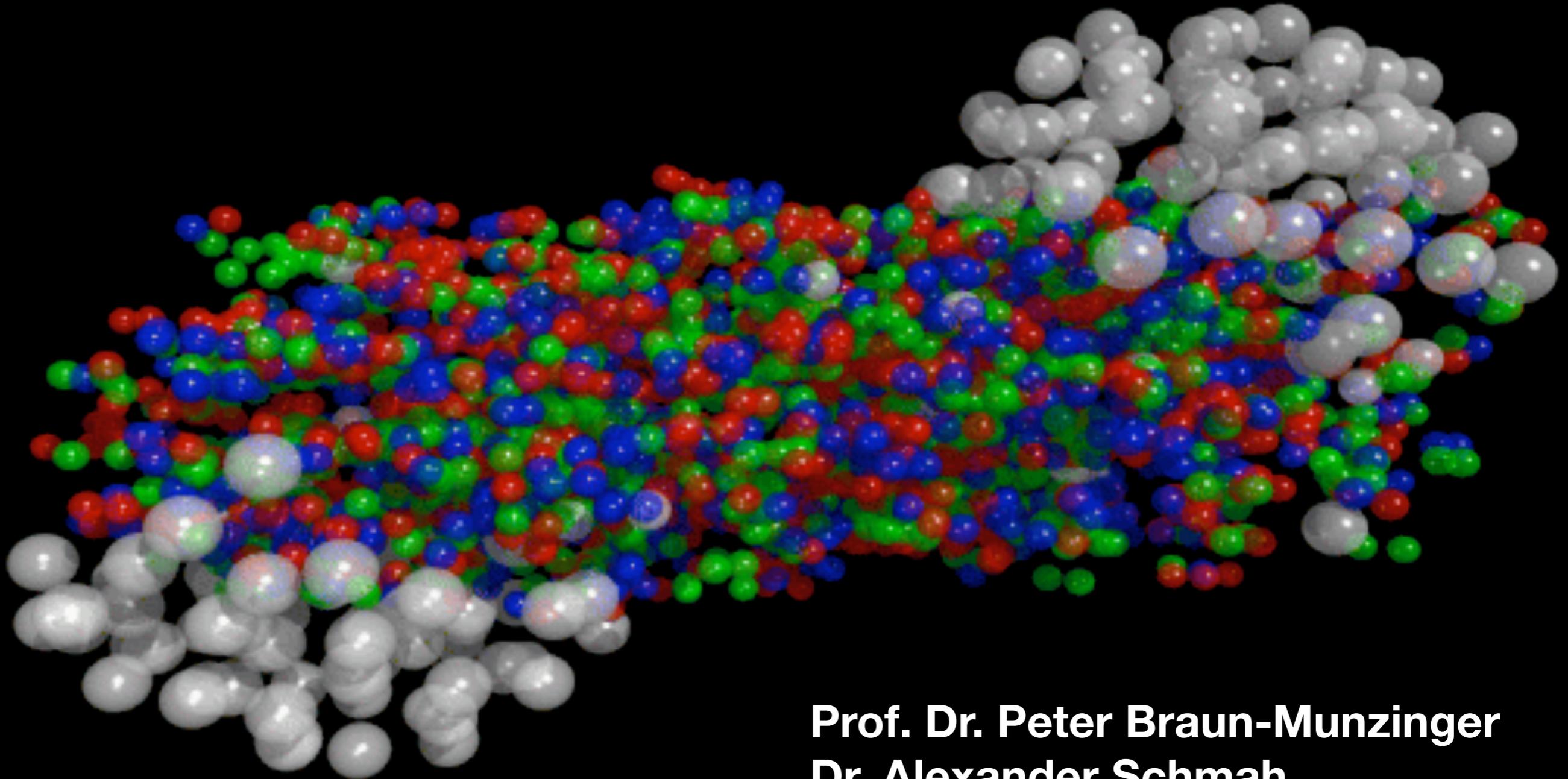


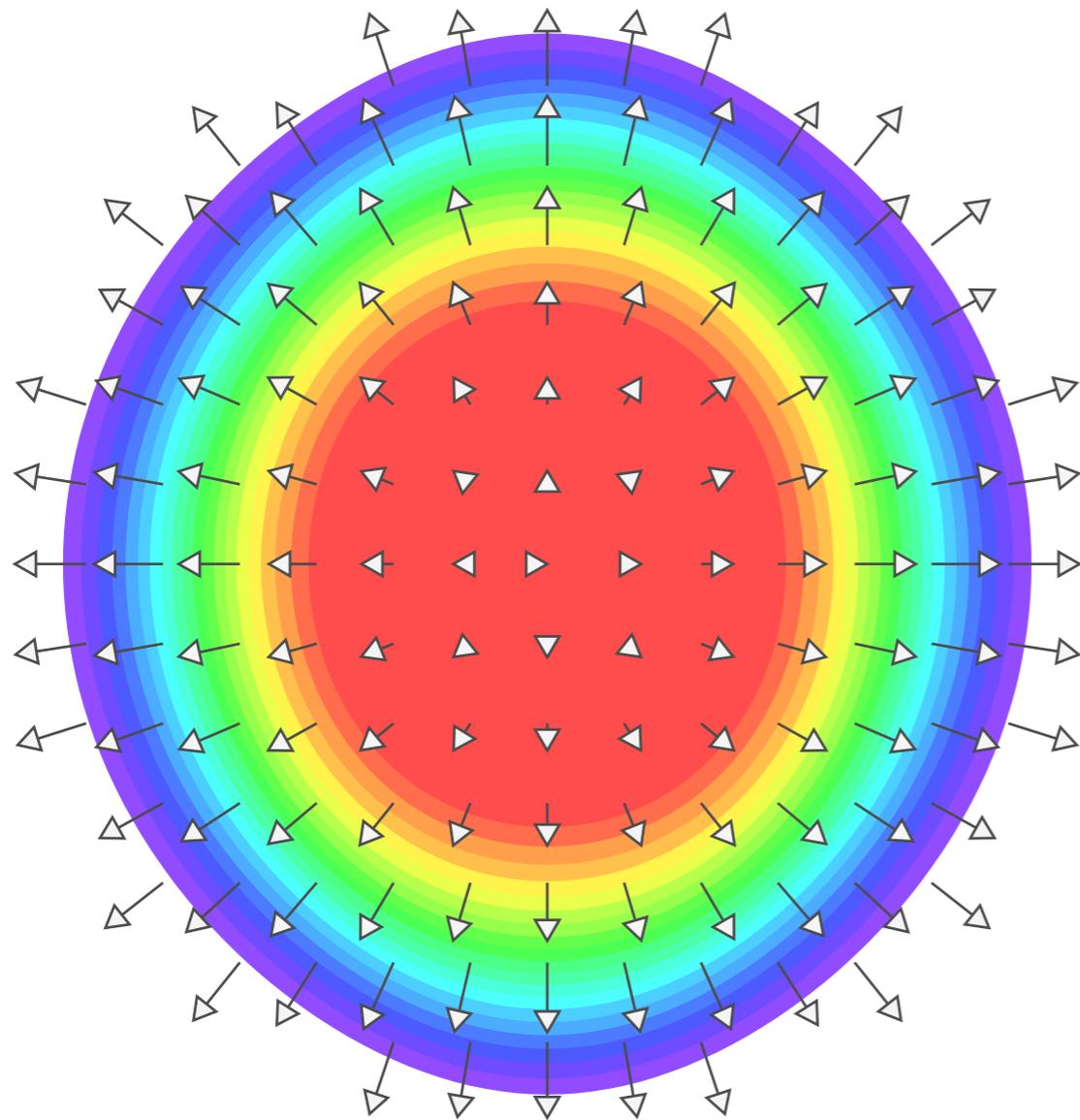
Quark-Gluon Plasma Physics

7. Space-time evolution of the QGP



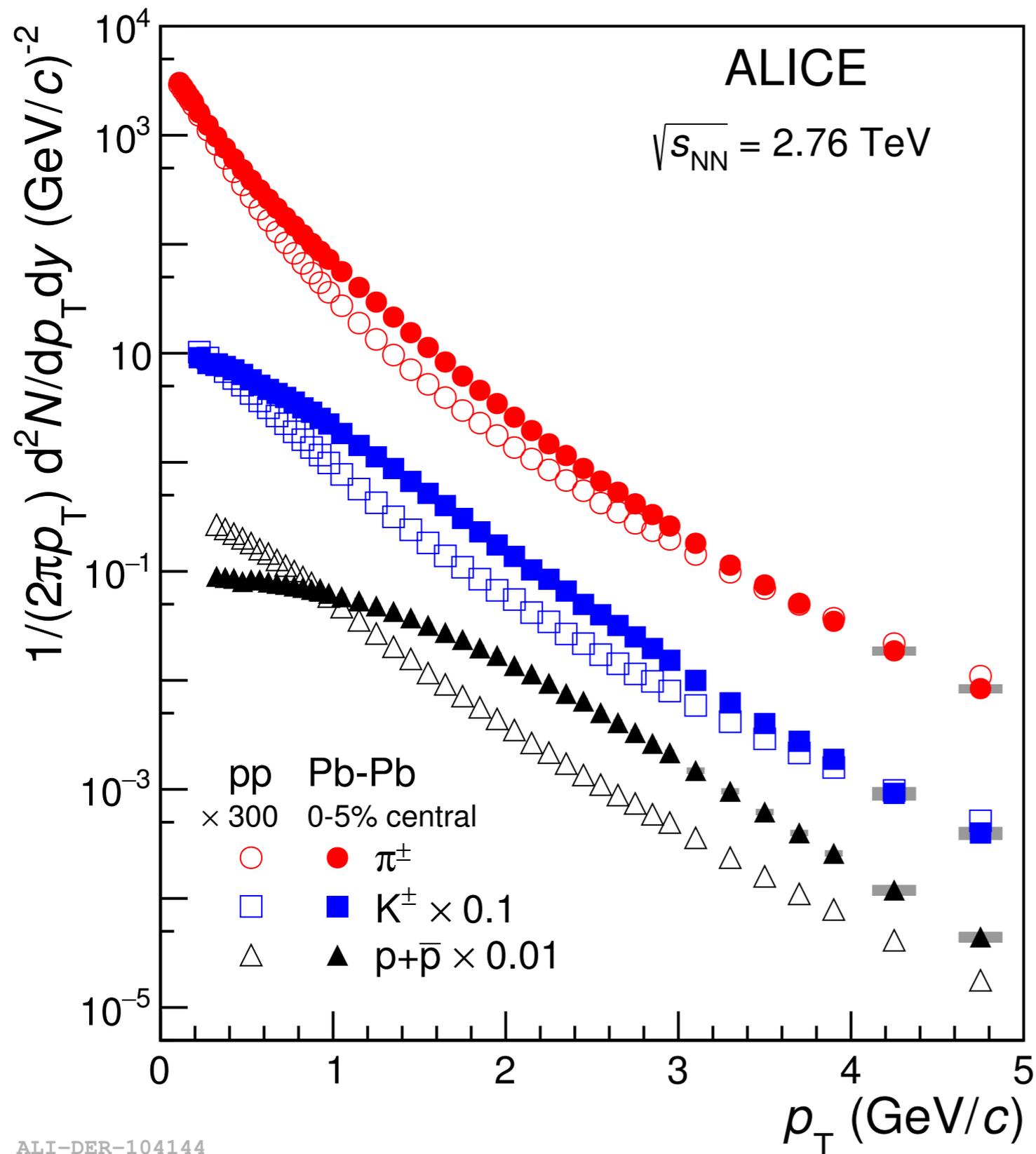
Prof. Dr. Peter Braun-Munzinger
Dr. Alexander Schmah
Heidelberg University
SS 2021

Evidence for collective behavior in heavy-ion collisions



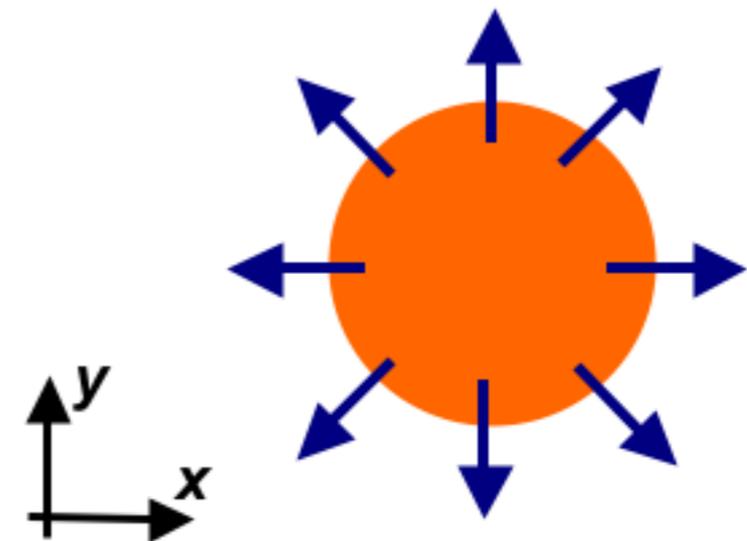
- Shape of low- p_T transverse momentum spectra for particles with different masses
- Azimuthal anisotropy of produced particles
- Source sizes from Hanbury Brown-Twiss correlations
- ...

Evidence for radial flow

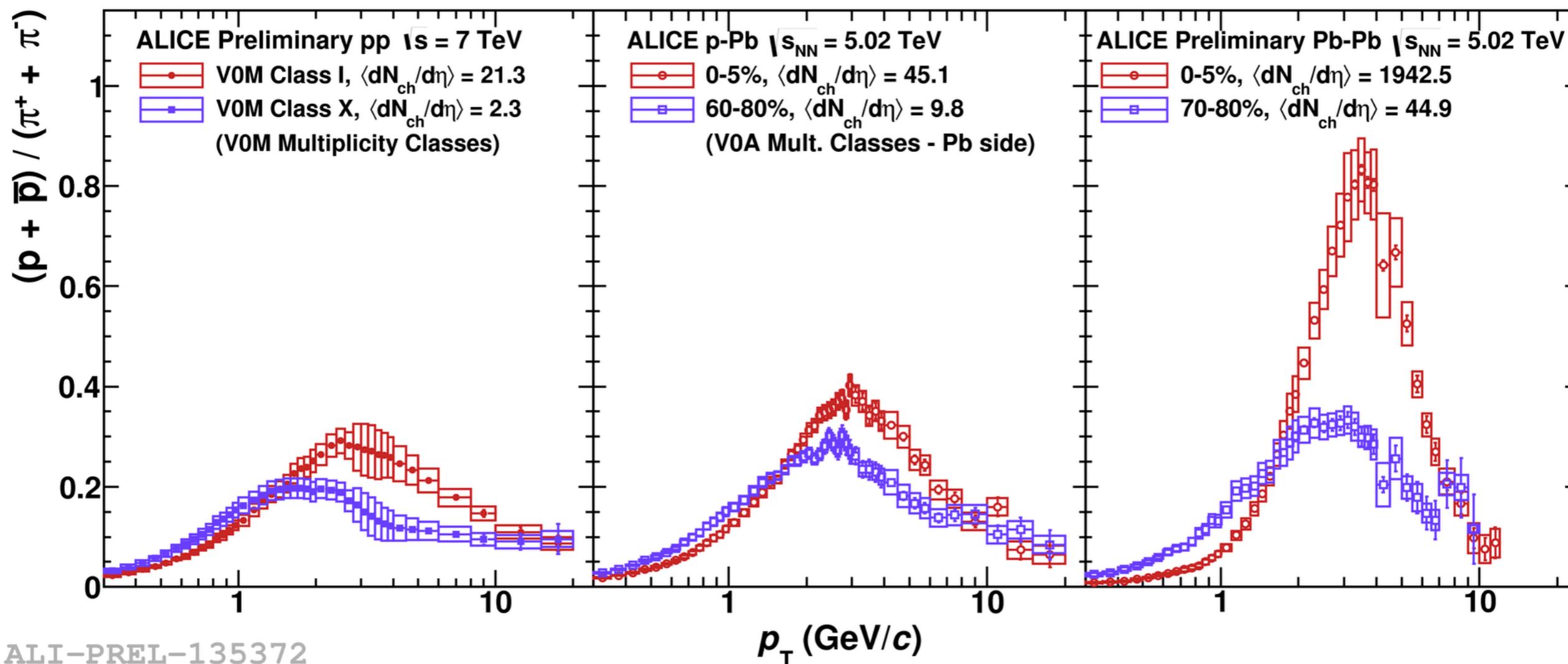


- Shape is different in pp and A-A
- Stronger effect for heavier particles

Radial flow



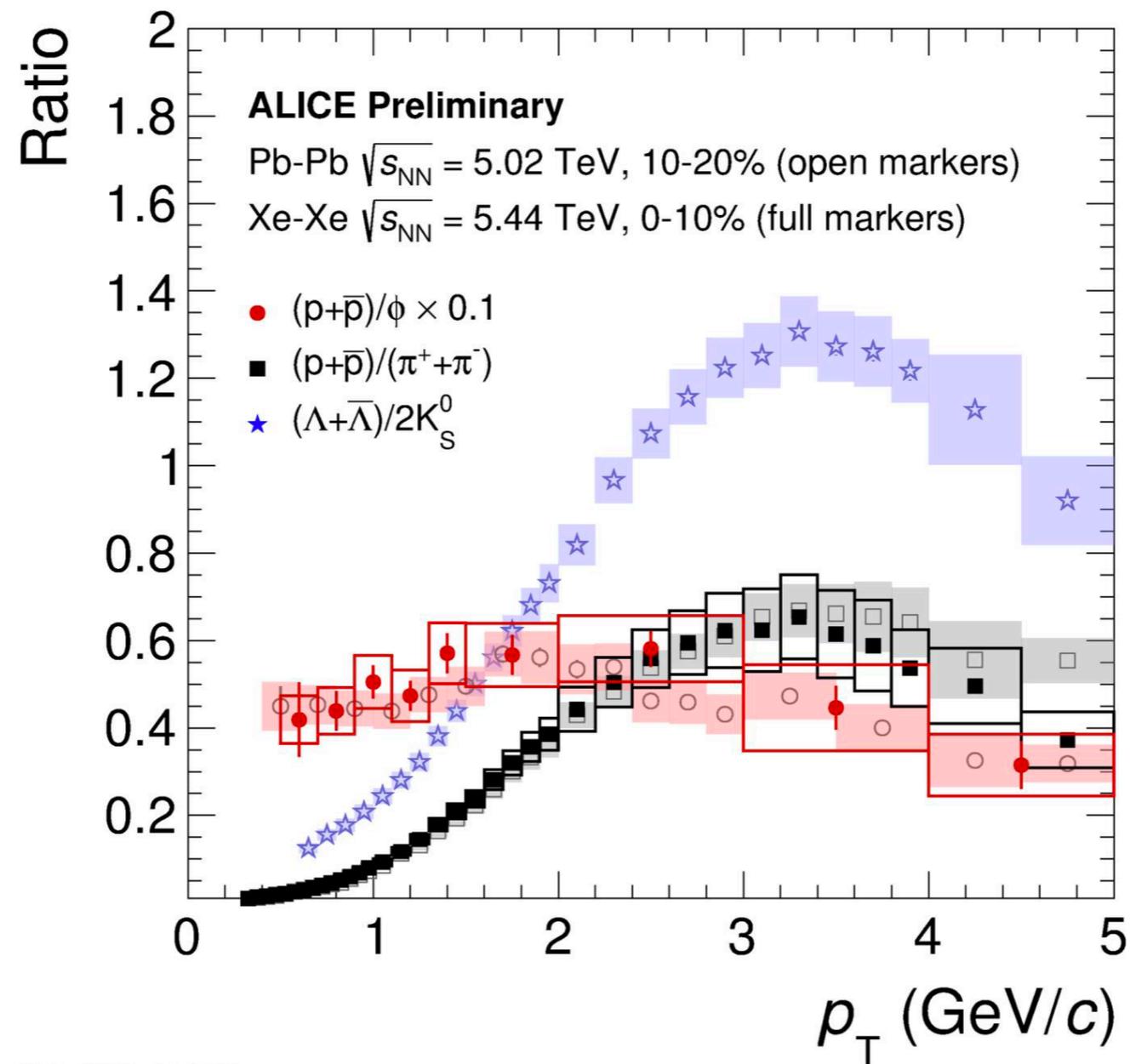
Ratio of particle spectra



$$p = \beta\gamma m$$

Centrality dependence of radial flow: effect increases with more central events

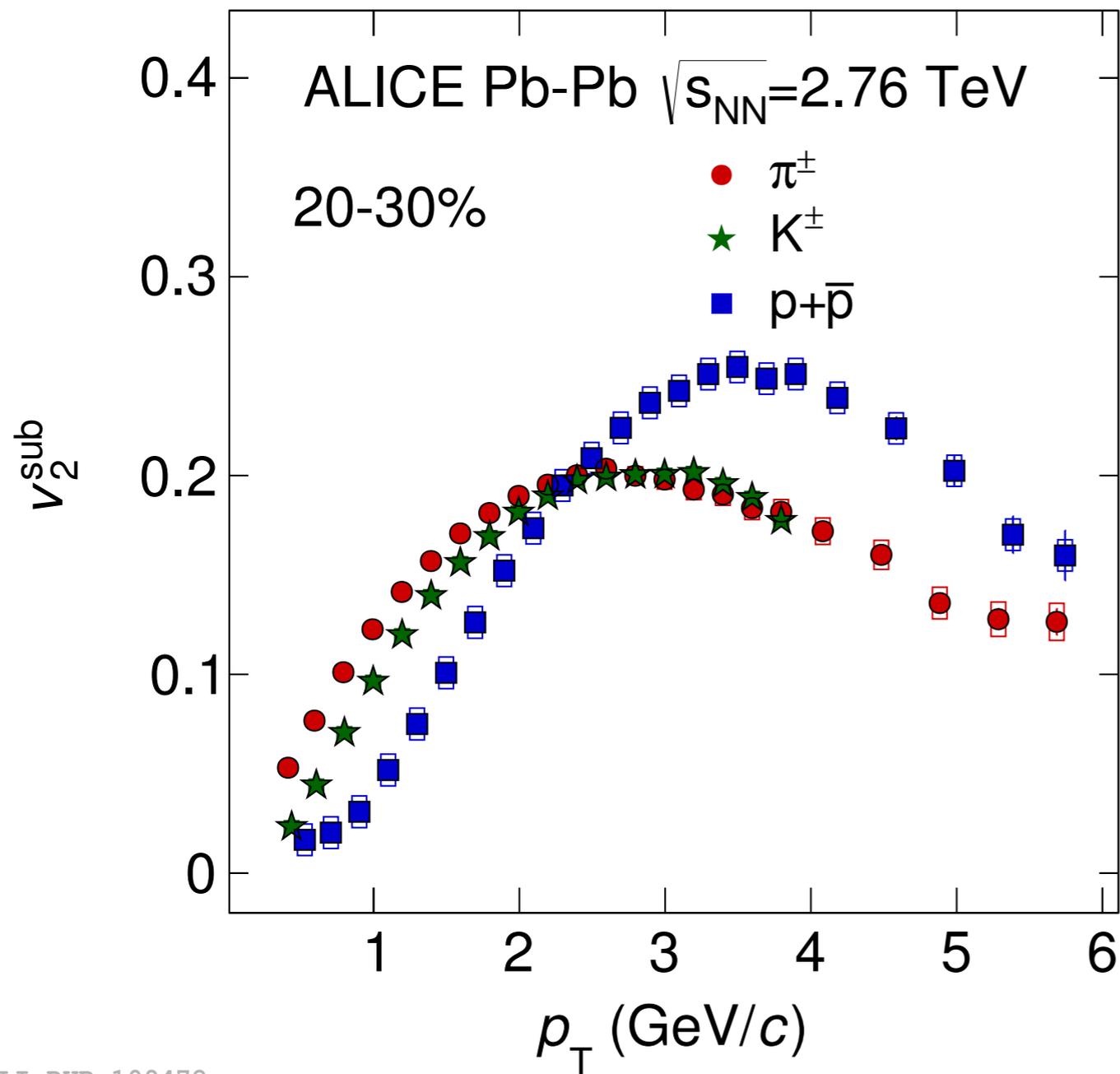
Baryon to Meson Ratios



ALI-PREL-156893

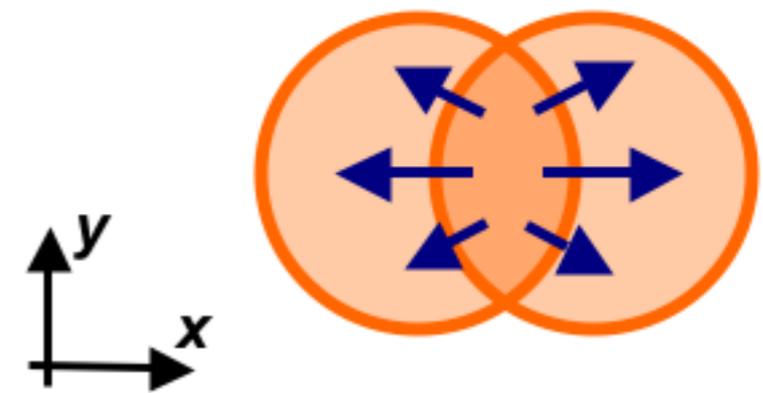
Ratio is almost flat for p/ϕ (similar mass)

Evidence for elliptic flow



Good explanation:
Azimuthal variation of the flow velocity

Elliptic flow



Basics of relativistic hydrodynamics

See e.g. Ollitrault,
arXiv:0708.2433

Standard thermodynamics: P , T , μ constant over the entire volume

Hydrodynamics assumes *local* thermodynamic equilibrium: $P(x^\mu)$, $T(x^\mu)$, $\mu(x^\mu)$

Local thermodynamic equilibrium only possible if mean free path between two collisions much shorter than all characteristic scales of the system:

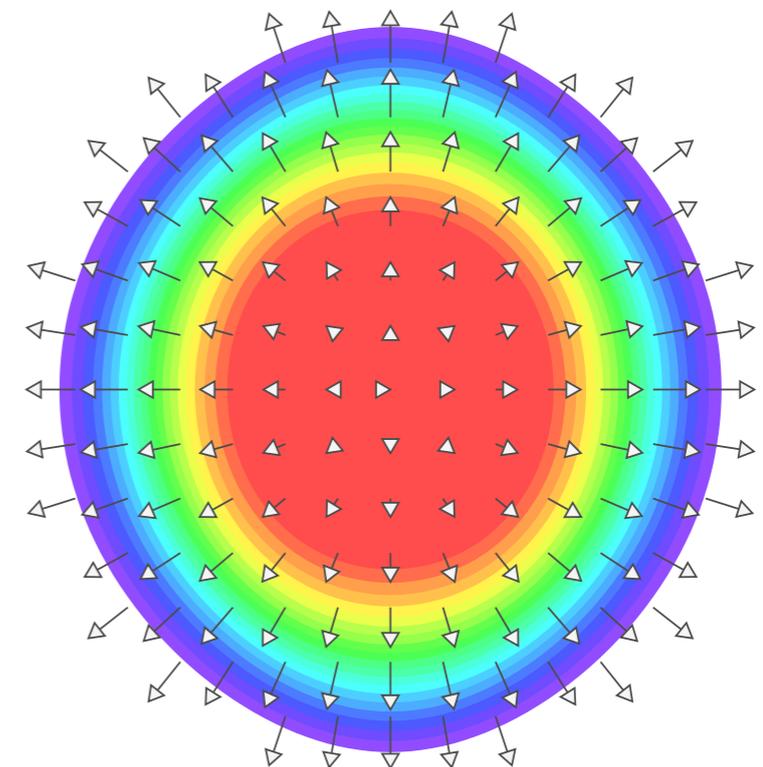
$$\lambda_{\text{mfp}} \ll L$$

This is the limit of non-viscous hydrodynamics.

4-velocity of a fluid element:

$$u = \gamma(1, \vec{\beta}), \quad u^\mu u_\mu = 1$$

$$\gamma = \frac{1}{\sqrt{1 - \vec{\beta}^2}}$$



Number conservation

Mass conservation in nonrelativistic hydrodynamics:

$$\frac{\partial \rho}{\partial t} + \vec{\nabla} \cdot (\rho \vec{v}) = 0 \quad [\text{continuity equation}]$$

Lorentz contraction in the relativistic case: $\rho \rightarrow n\gamma = nu^0$

conserved quantity,
e.g. baryon number

The continuity equation then reads: $\frac{\partial(nu^0)}{\partial t} + \vec{\nabla} \cdot (n\vec{u}) = 0$

nu^0 : baryon density
 $n\vec{u}$: baryon flux

The conservation of n can be written more elegantly as

$$\partial_\mu (nu^\mu) = 0$$

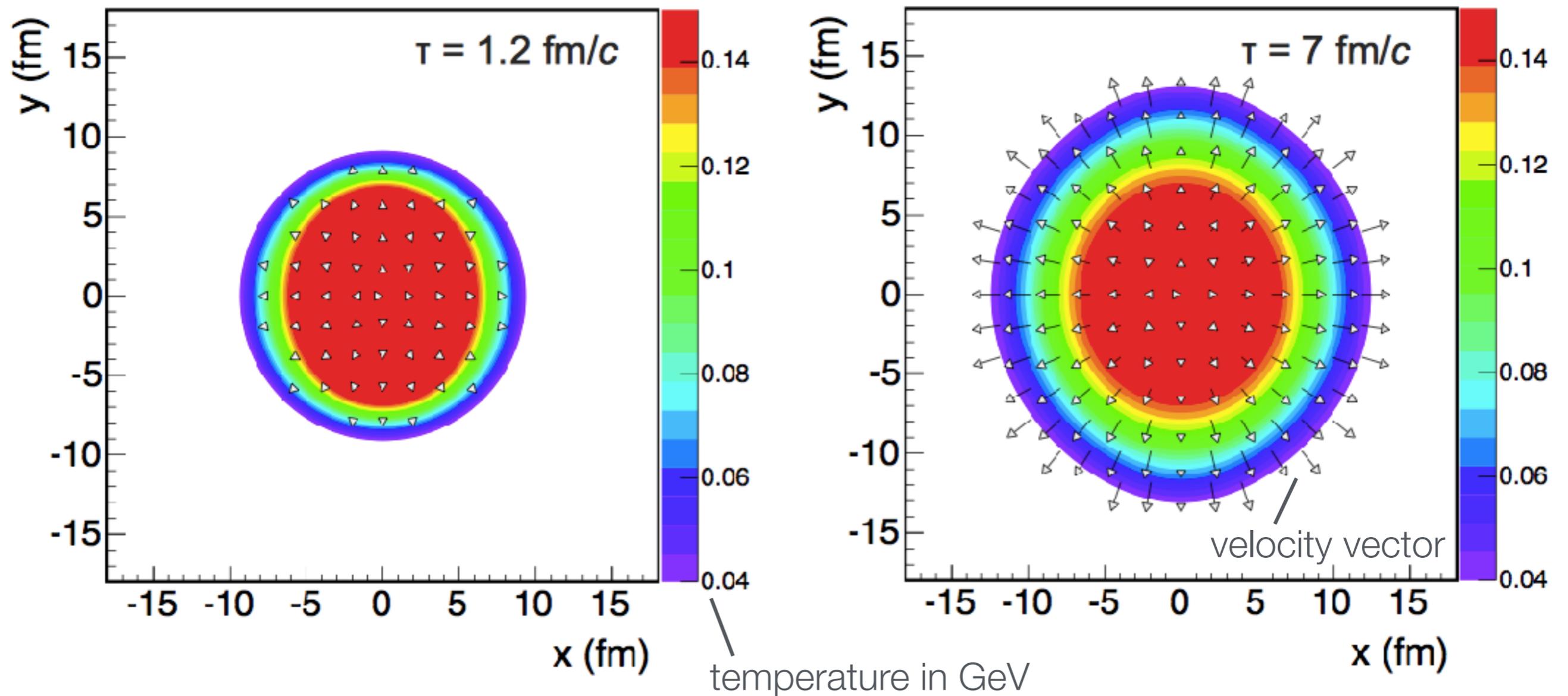
For a general 4-vector a we have:

$$\underbrace{\partial_\mu \equiv \frac{\partial}{\partial x^\mu}}_{\text{covariant derivative}} = \left(\frac{\partial}{\partial t}, \vec{\nabla} \right), \quad \underbrace{\partial^\mu \equiv \frac{\partial}{\partial x_\mu}}_{\text{contravariant derivative}} = \left(\frac{\partial}{\partial t}, -\vec{\nabla} \right), \quad \partial_\mu a^\mu = \left(\frac{\partial}{\partial t}, \vec{\nabla} \right) \cdot (a^0, \vec{a}) = \frac{\partial a^0}{\partial t} + \vec{\nabla} \cdot \vec{a}$$

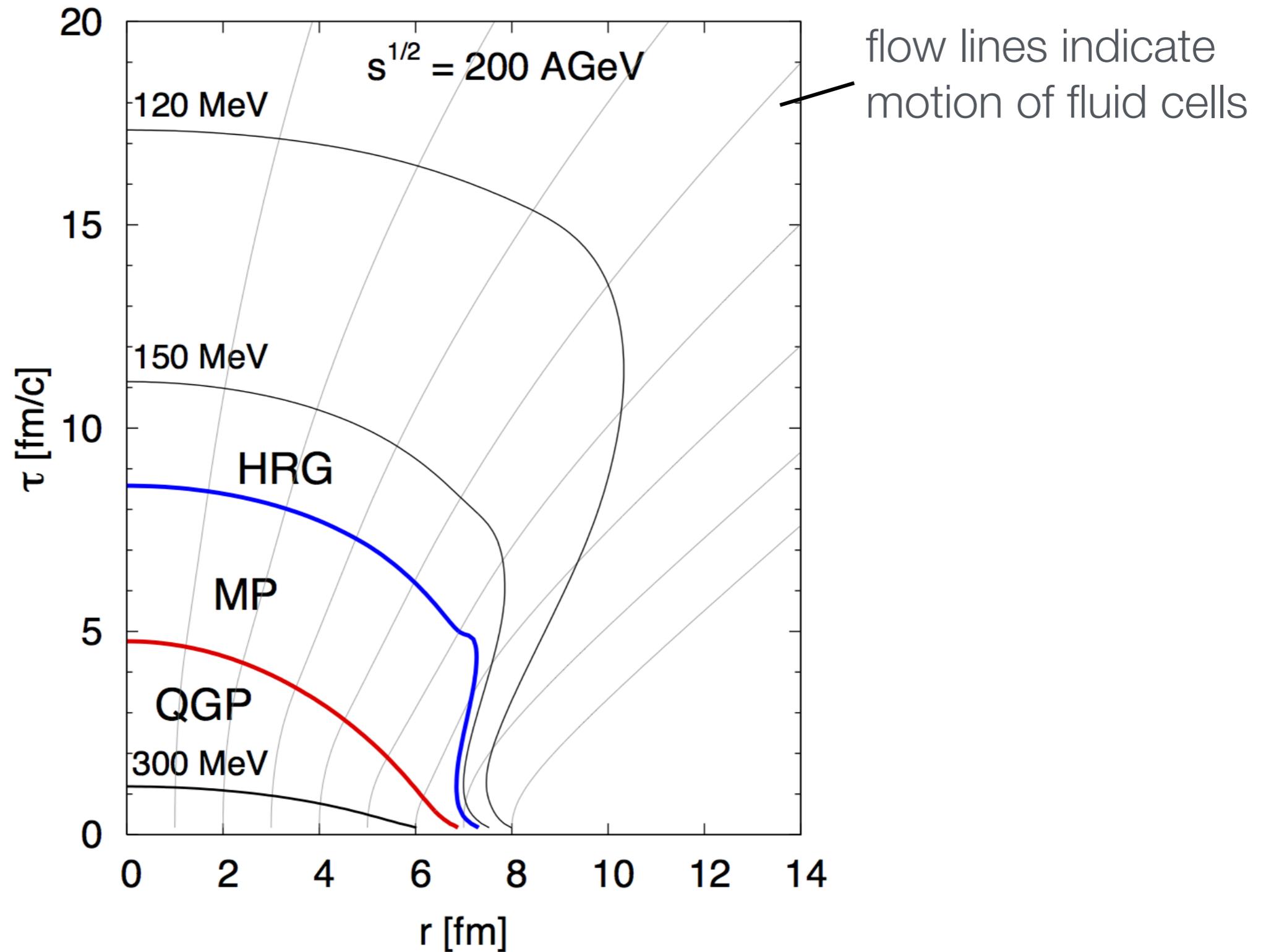
Transverse expansion

Transverse expansion of the fireball in a hydro model (temperature profile)

2+1 d hydro: Bjorken flow in longitudinal direction



Temperature Contours and Flow lines



Viscosity

Pitch drop experiment, started in Queensland, Australia in 1927

Date	Event	Duration		
		Years	Months	
1927	Hot pitch poured			
October 1930	Stem cut			
December 1938	1st drop fell	8.1	98	██████████
February 1947	2nd drop fell	8.2	99	██████████
April 1954	3rd drop fell	7.2	86	██████████
May 1962	4th drop fell	8.1	97	██████████
August 1970	5th drop fell	8.3	99	██████████
April 1979	6th drop fell	8.7	104	██████████
July 1988	7th drop fell	9.2	111	██████████
November 2000	8th drop fell ^[A]	12.3	148	██████████
April 2014	9th drop ^[B]	13.4	156	██████████

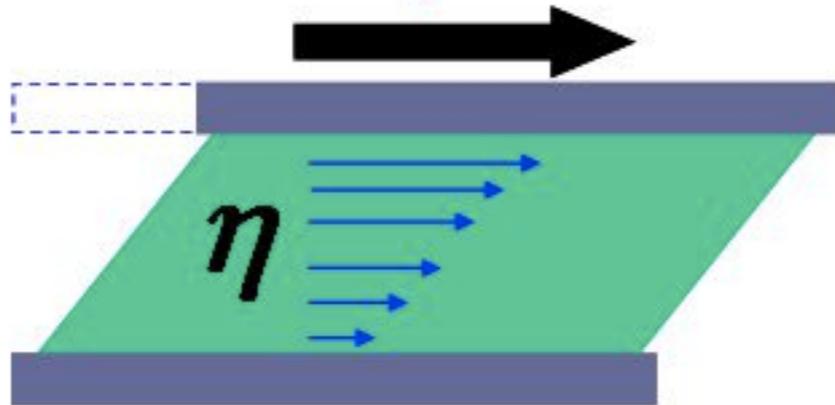
Meaningful comparison of different fluids: η/s



https://en.wikipedia.org/wiki/Pitch_drop_experiment

Shear and bulk viscosity

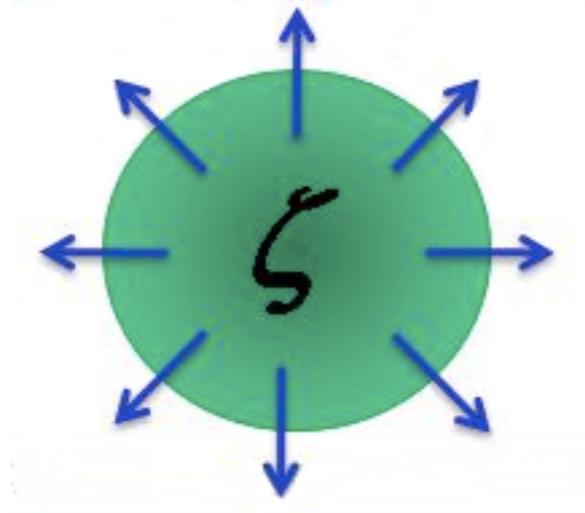
Shear viscosity



Acts against buildup of flow anisotropies ($v_2, v_3, v_4, v_5, \dots$)

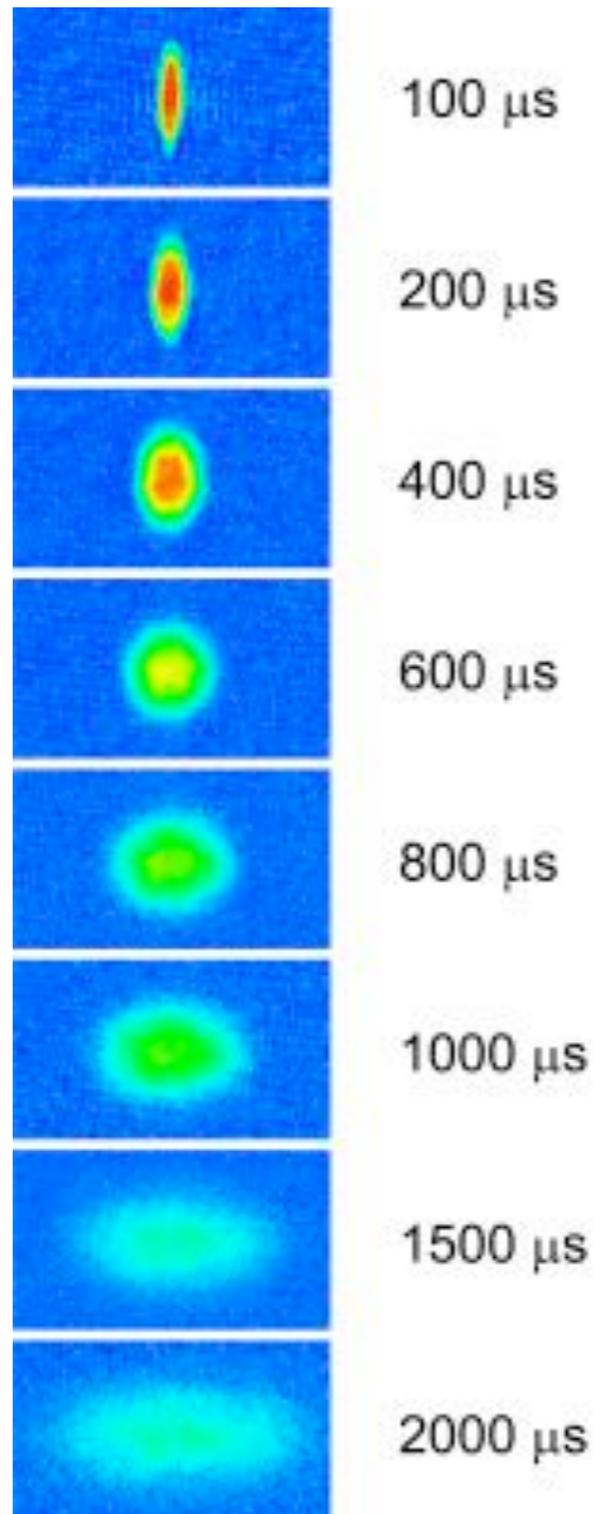
η/s : shear viscosity per entropy density ratio

Bulk viscosity



Acts against buildup of radial flow

Universal aspects of the underlying physics



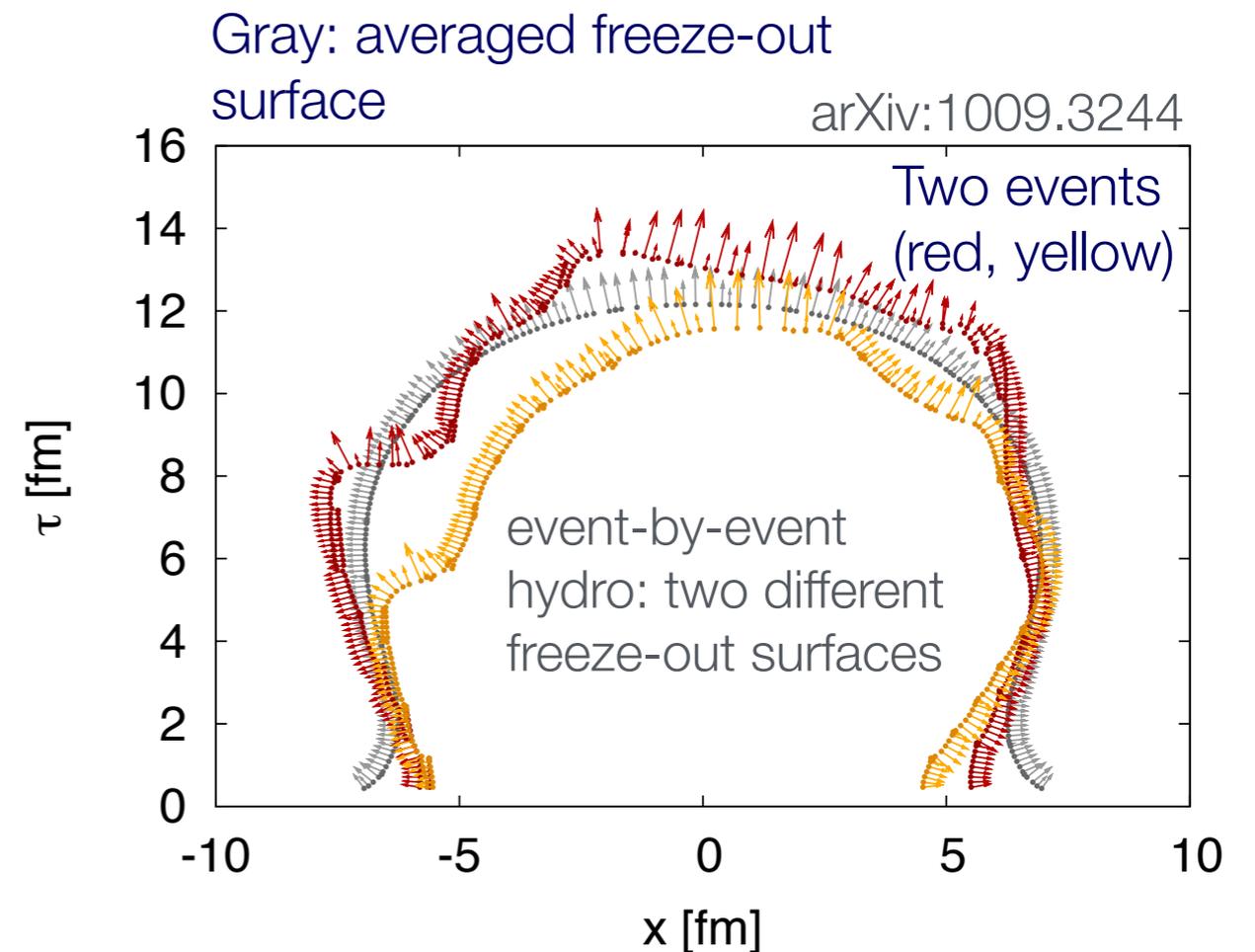
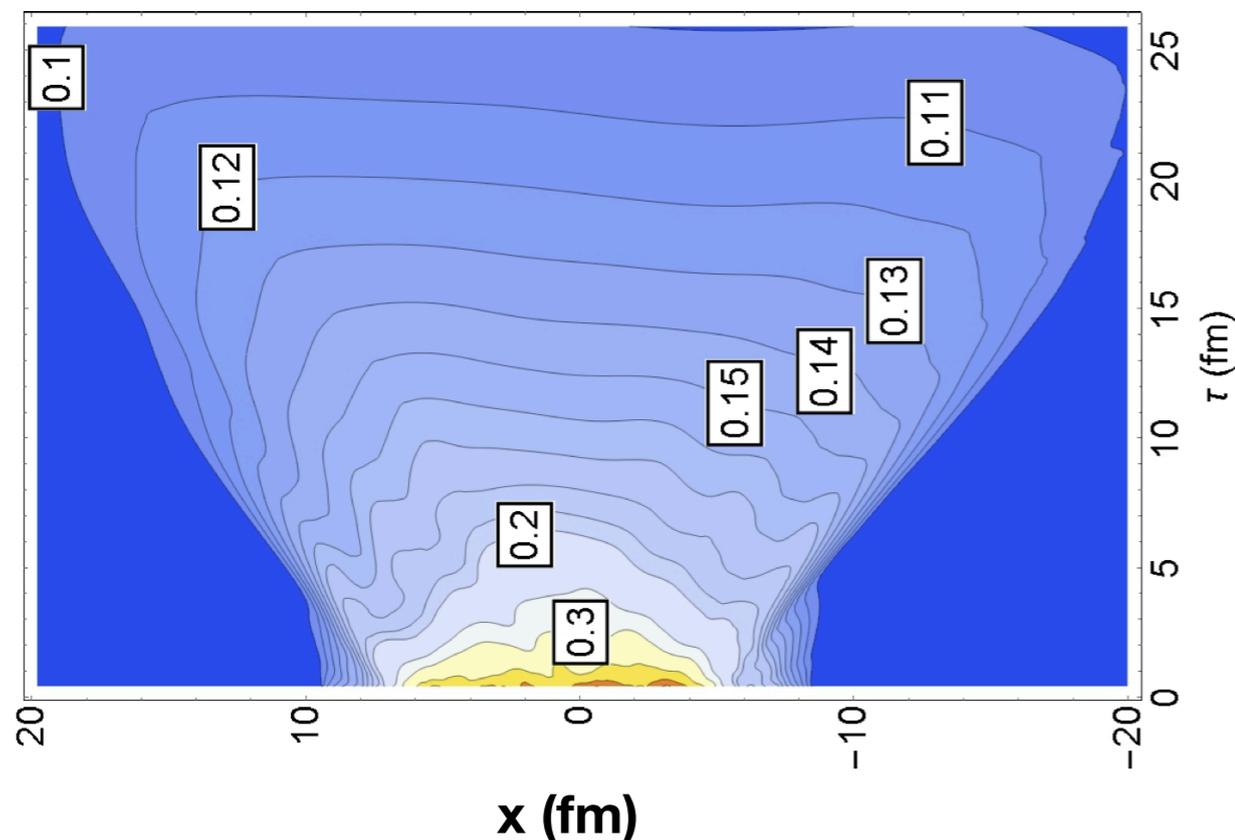
- Strongly-interacting degenerate gas of fermionic ${}^6\text{Li}$ atoms at $0.1 \mu\text{K}$
- Cigar-shaped cloud initially trapped by a laser field
- Anisotropic expansion upon abruptly turning off the trap: Elliptic flow!
- η/s can be extracted: [\[PhD thesis Chenglin Cao\]](#)

$$(\eta/s)_{{}^6\text{Li gas}} \approx 0.4 = 5 \times \frac{1}{4\pi}$$

The ultimate goal is to unveil the universal physical laws governing seemingly different physical systems (with temperature scales differing by 19 order of magnitude)

Hydrodynamic modeling of heavy-ion collisions: State of the art

- Equation of state from lattice QCD
- (2+1)D or (3+1)D viscous hydrodynamics
- Fluctuating initial conditions (event-by-event hydro)
- Hydrodynamic evolution followed by hadronic cascade



The blast-wave model: A Simple model to describe the effect of radial flow on particle spectra

Transverse velocity profile: $\beta_T(r) = \beta_s \left(\frac{r}{R}\right)^n$

Superposition of thermal sources with different radial velocities:

$$\frac{1}{m_T} \frac{dn}{dm_T} \propto \int_0^R r dr m_T l_0 \left(\frac{p_T \sinh \rho}{T} \right) K_1 \left(\frac{m_T \cosh \rho}{T} \right)$$

$\rho := \operatorname{arctanh}(\beta_T)$ "transverse rapidity"

l_0, K_1 : modified Bessel functions

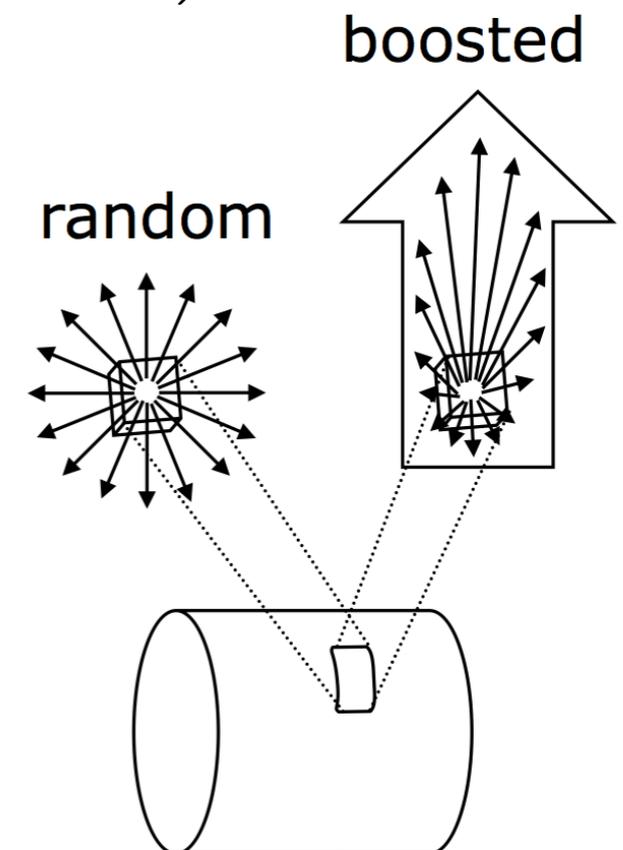
Schnedermann, Sollfrank, Heinz,
Phys.Rev.C48:2462-2475,1993

K. Reygers, A. Schmah, A. Berdnikova, X. Sun

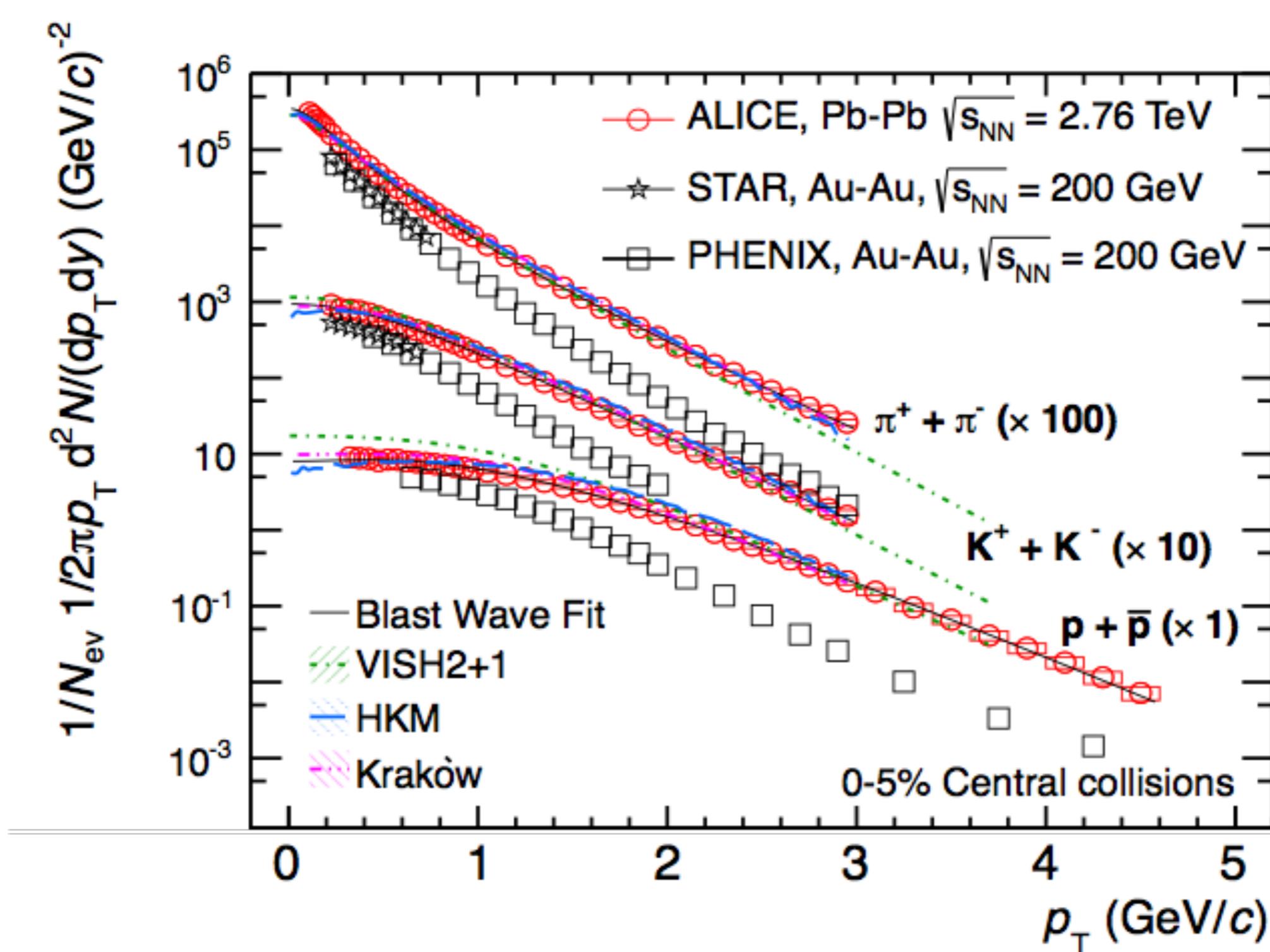
[10.1103/PhysRevC.101.064905](https://doi.org/10.1103/PhysRevC.101.064905)

Freeze-out at a 3d hyper-surface,
typically instantaneous, e.g.:

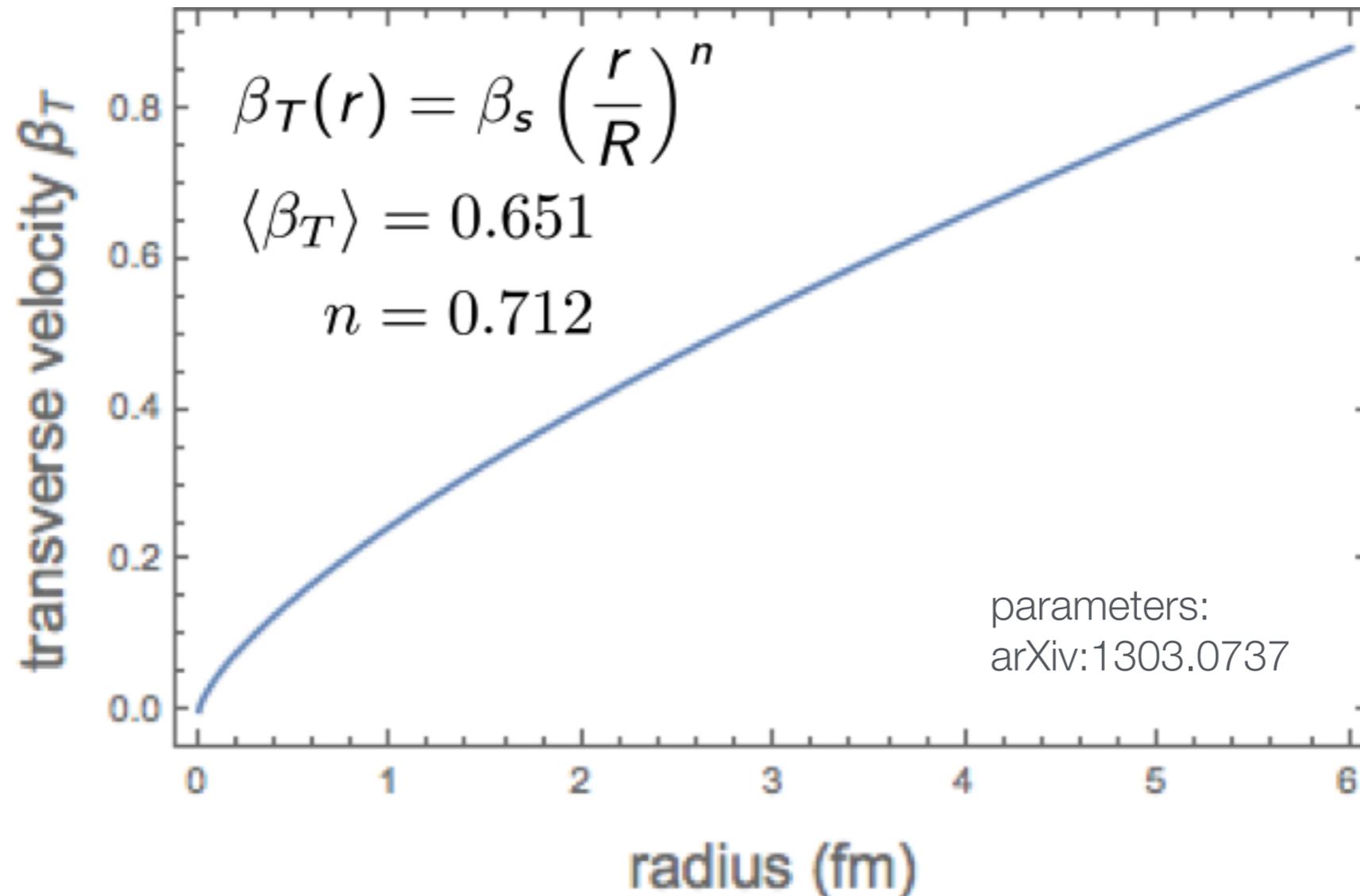
$$t_f(r, z) = \sqrt{\tau_f^2 + z^2}$$



Comparison of π , K , p spectra with hydro and blast-wave models



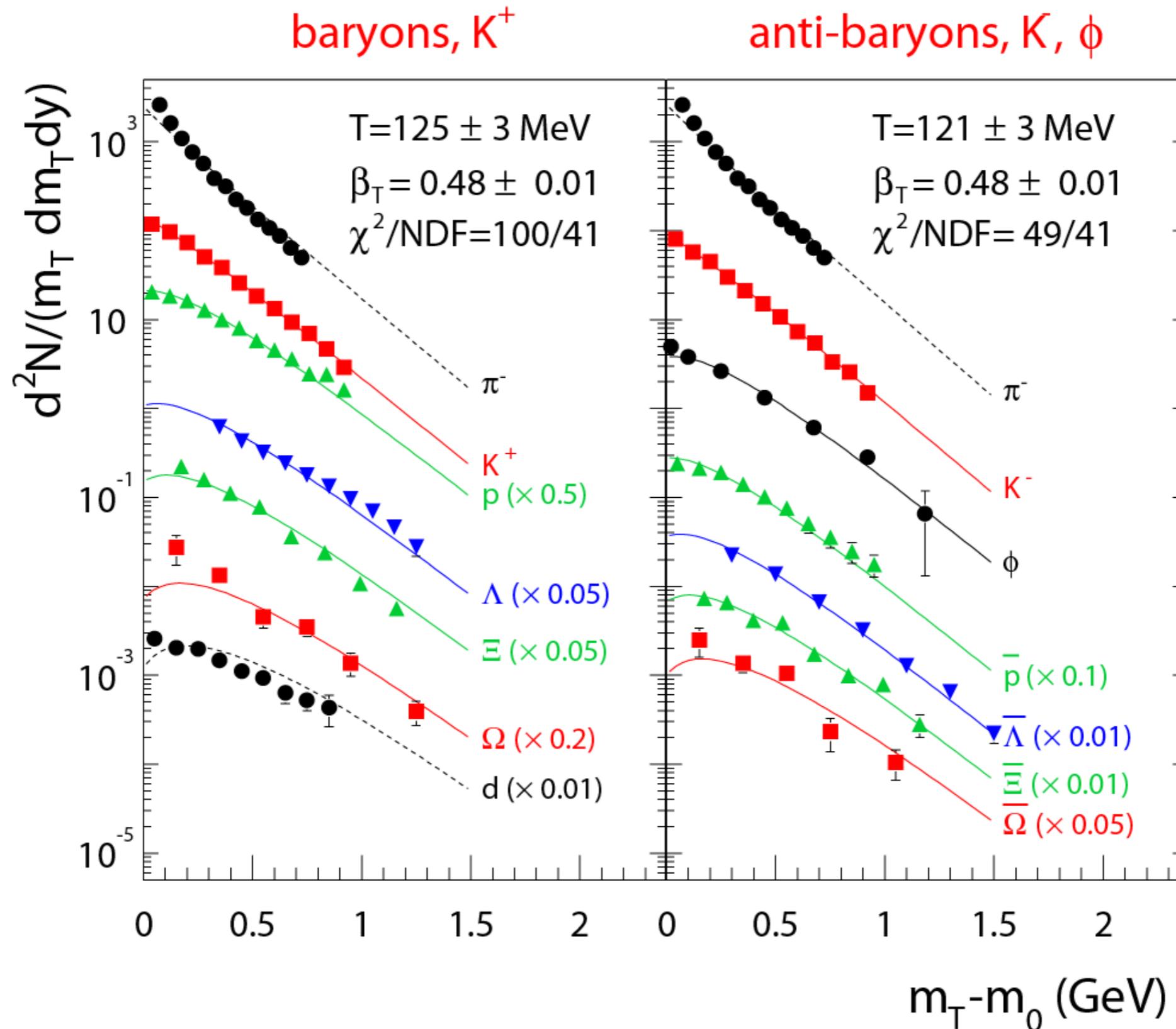
Example: Radial Flow Velocity Profile from Blast-wave Fit to 2.76 TeV Pb-Pb Spectra (0-5%)



$$\langle \beta_T \rangle = \frac{\int_0^R \int_0^{2\pi} r dr d\varphi \beta_T(r)}{\int_0^R \int_0^{2\pi} r dr d\varphi} = \frac{2}{n+2} \beta_s \quad \langle \beta_T \rangle = 0.651, n = 0.712$$

$$\rightarrow \beta_s = 0.8$$

Blast-wave fit for CERN SPS data (NA49)

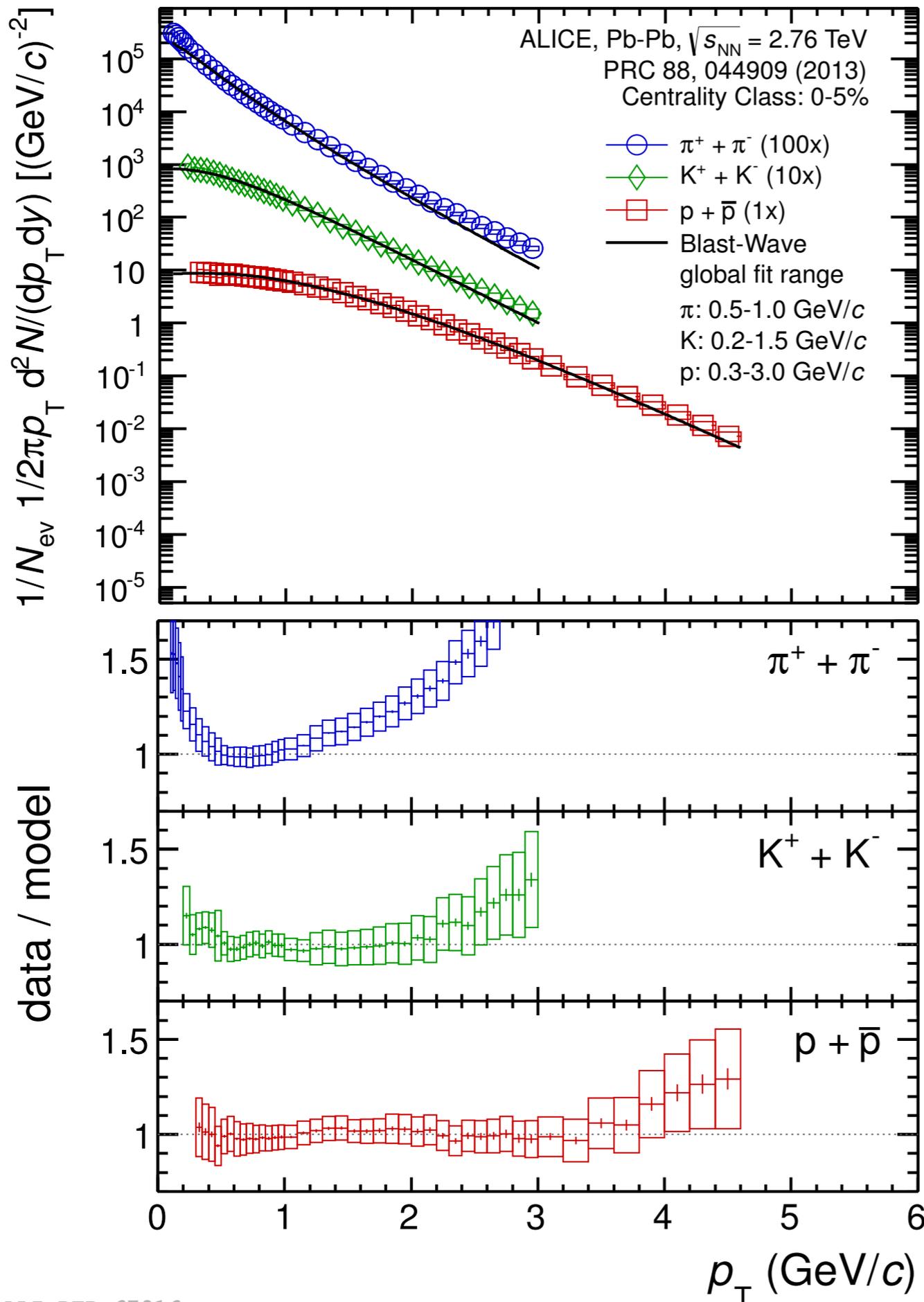


Blast-wave fit LHC

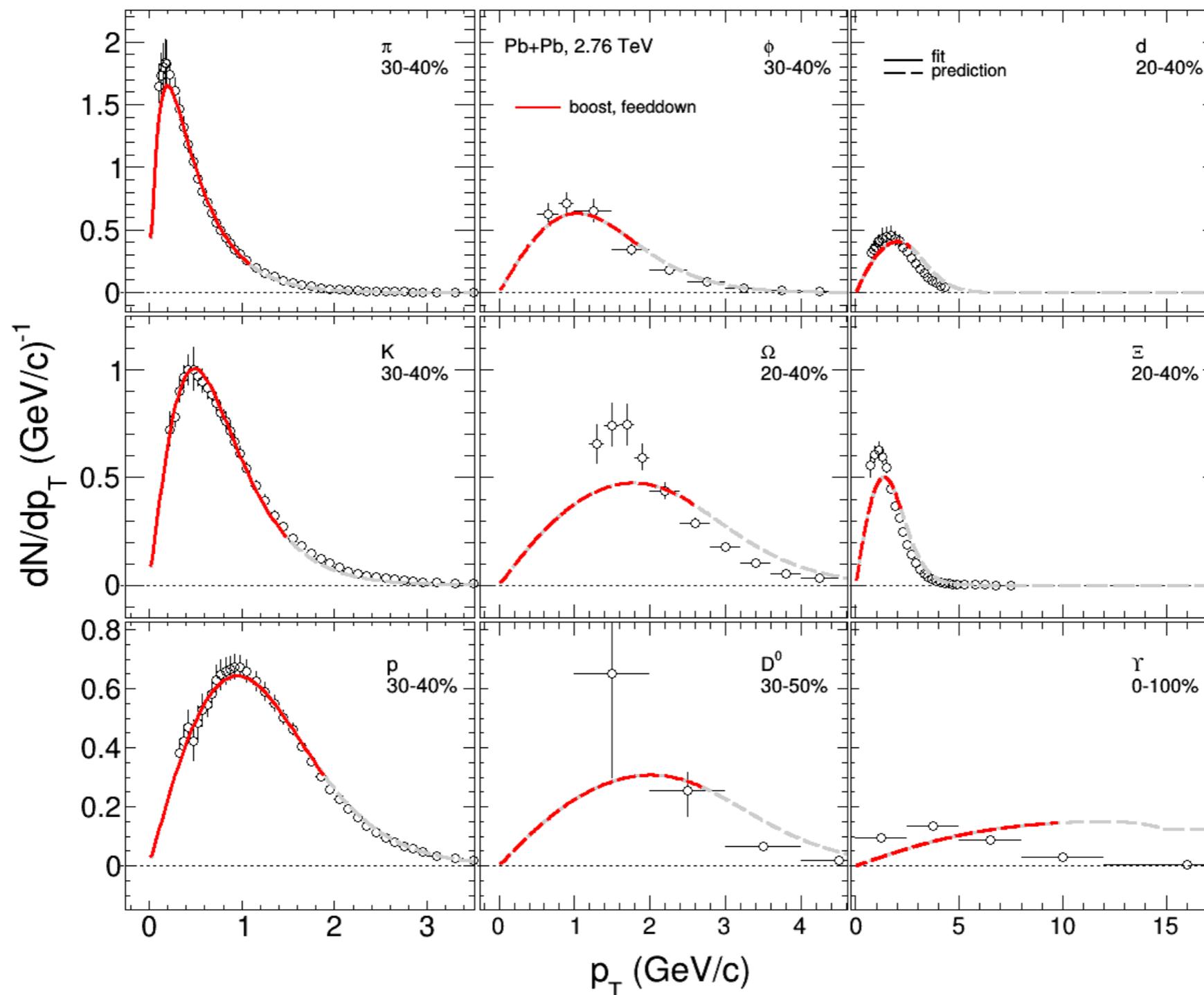
Works well for K and p

For pions, the deviation at low transverse momentum is currently the subject of intense investigations with no final conclusion yet.

At very high p_T the deviations are mainly due to jets (high p_T parton-parton scattering).

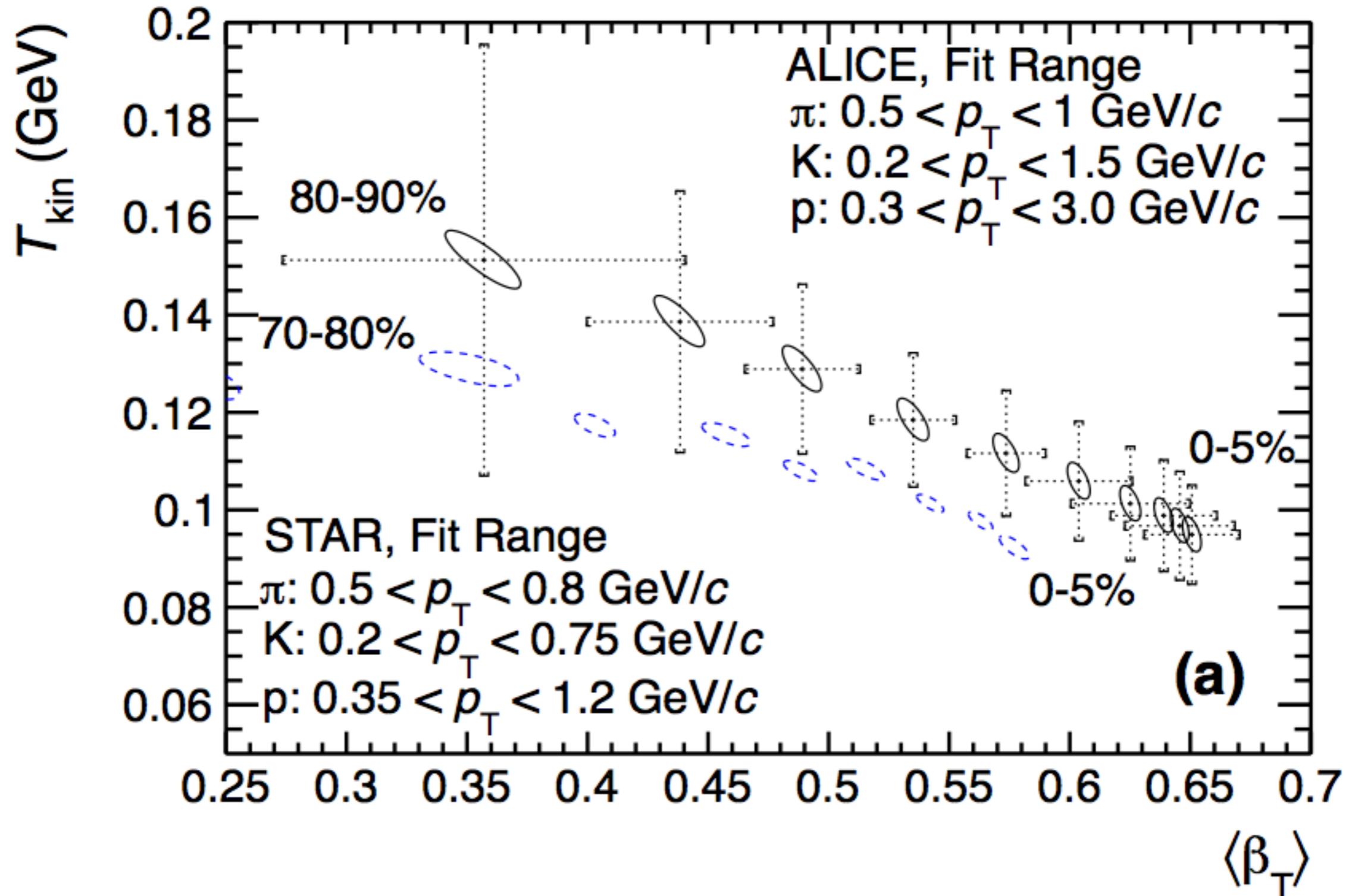


Most recent results on blast-wave



A. Schmah, K. Reygers, X. Sun, A. Berdnikova, N. Gruenwald
Publication in process

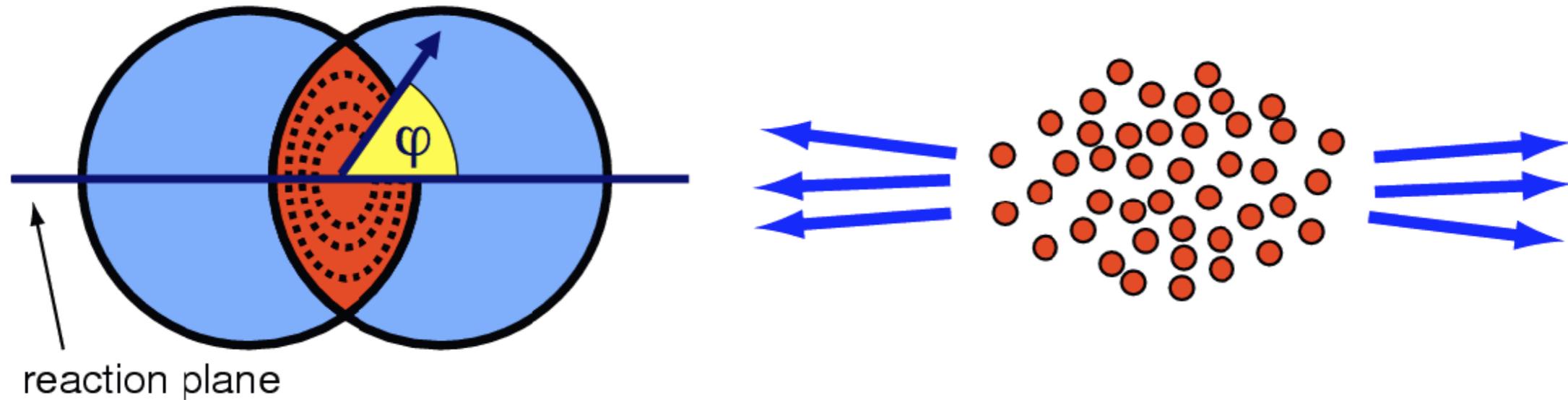
T und $\langle\beta\rangle$ for different centralities at RHIC and the LHC



10% larger flow velocities in central collisions at the LHC than at RHIC

Elliptic flow and higher flow harmonics

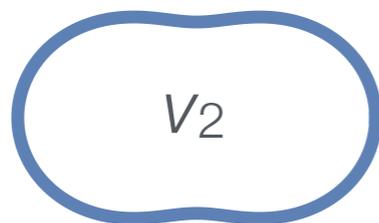
Azimuthal distribution of produced particles



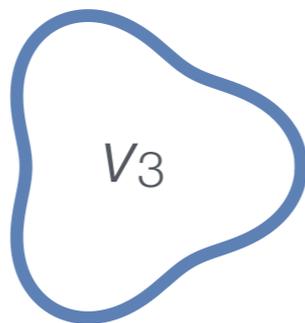
$$\frac{dN}{d\varphi} \propto 1 + 2 \sum_{n=1}^{\infty} v_n \cos[n(\varphi - \psi_n)]$$

Fourier coefficients:

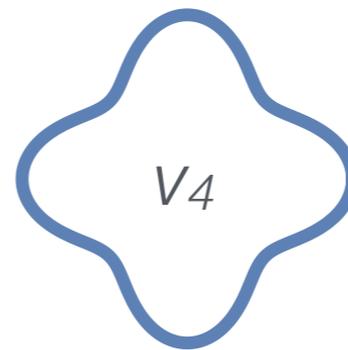
$$v_n(p_T, y) = \langle \cos[n(\varphi - \psi_n)] \rangle$$



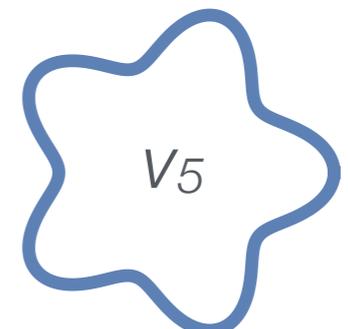
elliptic flow



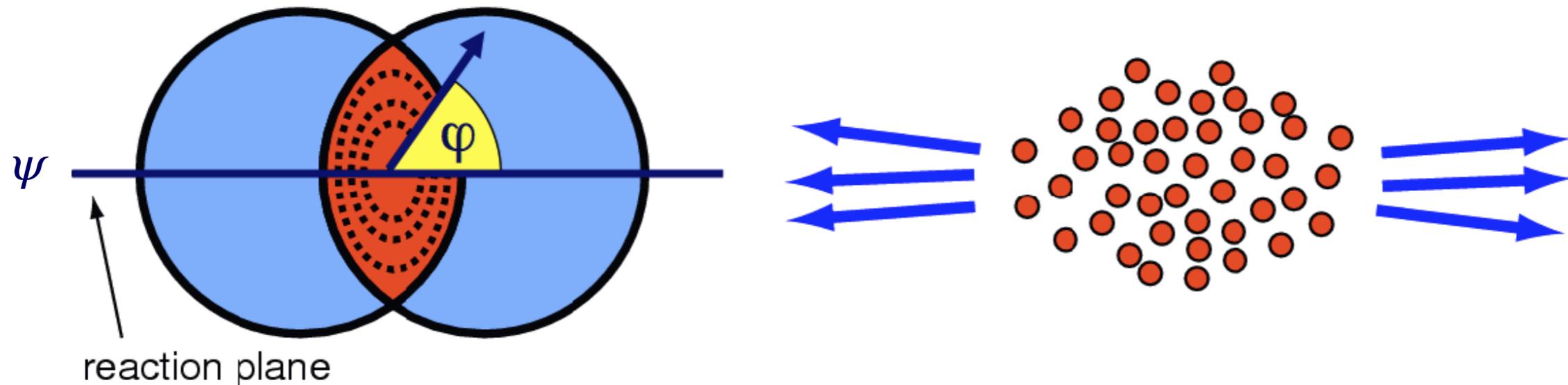
triangular flow



$$f(\varphi) = 1 + 2v_n \cos(n\varphi)$$



Event plane reconstruction



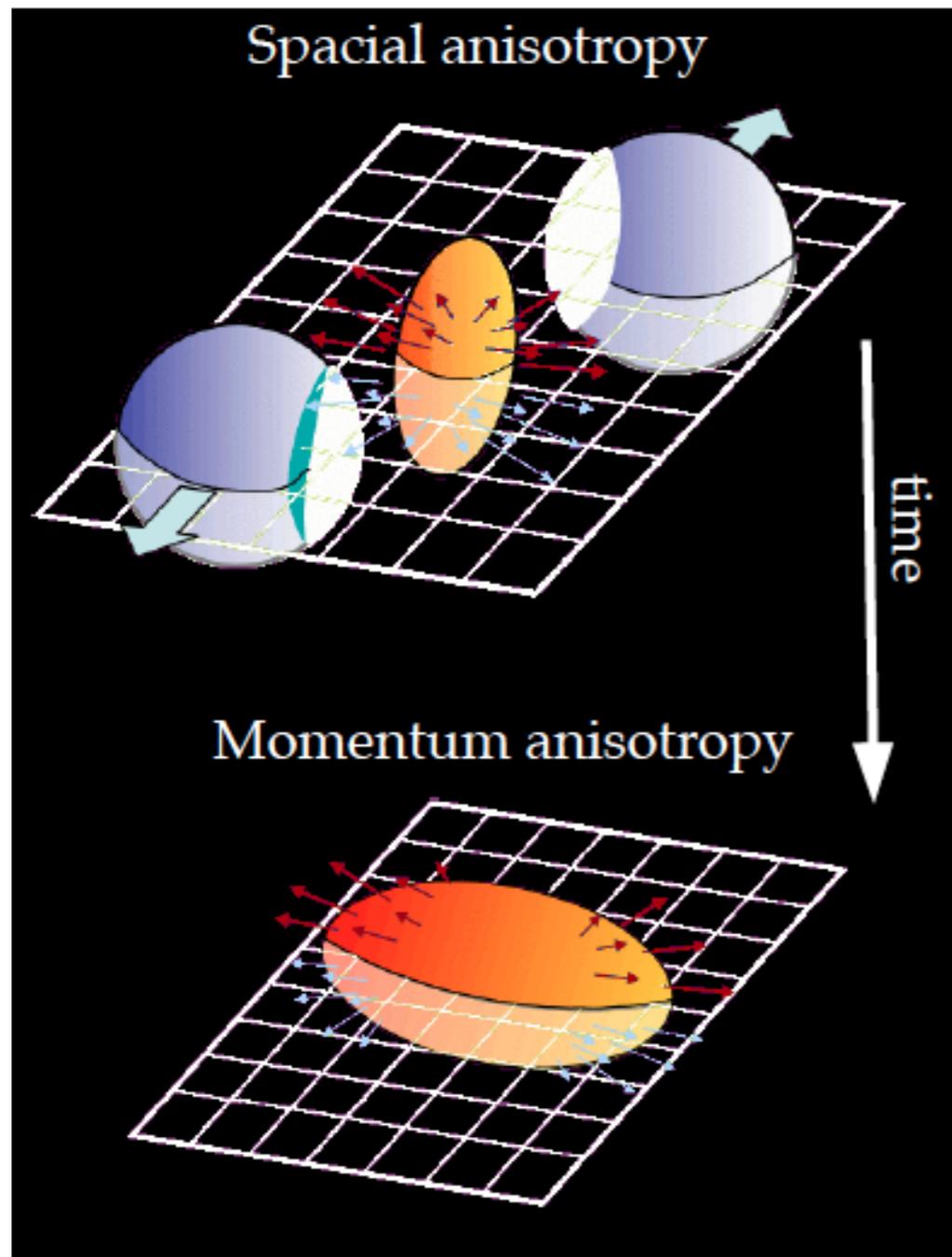
- Particles fly dominantly into the event (reaction) plane
- Use the produced particles to calculate event plane angle:

$$\Psi_2 = \frac{1}{2} \left(\tan^{-1} \frac{\sum_i w_i \sin(2\phi_i)}{\sum_i w_i \cos(2\phi_i)} \right)$$

- w_i : weights, usually the pT of the particle

Elliptic flow - Space-Time Evolution

Non-central heavy-ion collision



Hydrodynamic evolution

$$v_2 = \frac{\langle p_x^2 \rangle - \langle p_y^2 \rangle}{\langle p_x^2 \rangle + \langle p_y^2 \rangle}$$

$$v_2 = \langle \cos(2\phi - 2\Psi) \rangle$$

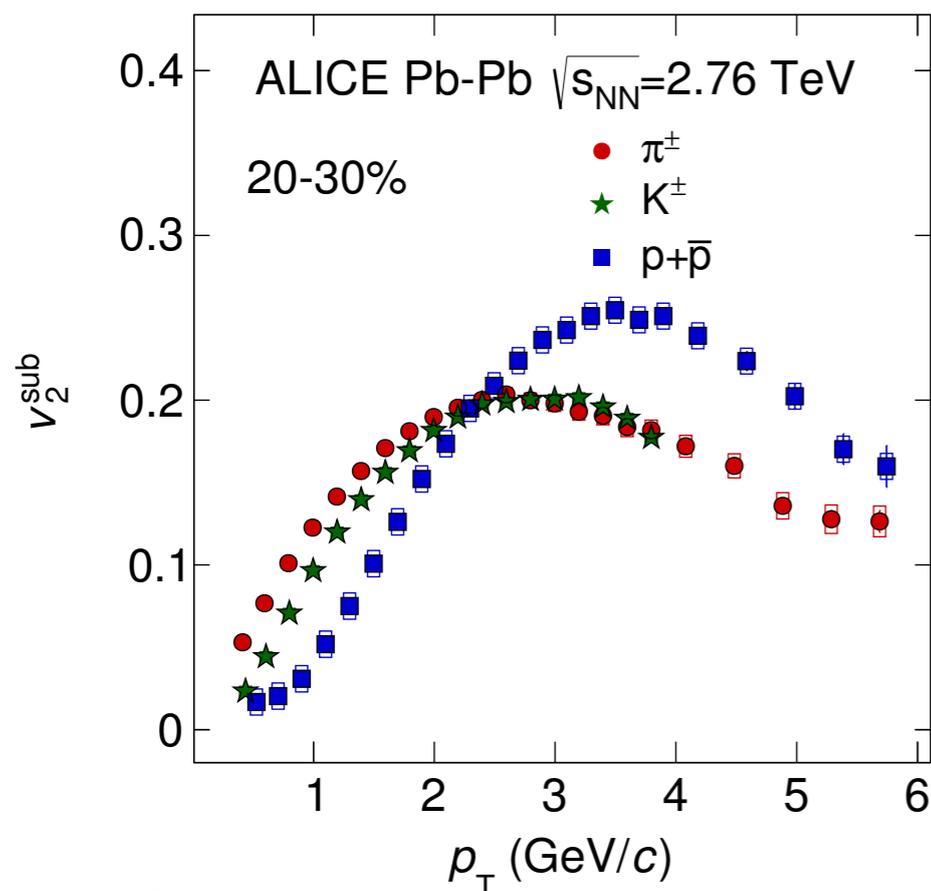
v_2 = elliptic flow

ϕ : particle azimuthal angle

ψ : event plane

- Anisotropic expansion can tell us a lot about the medium!
 - Kind of constituents
 - Interaction between the constituents
 - Viscosity

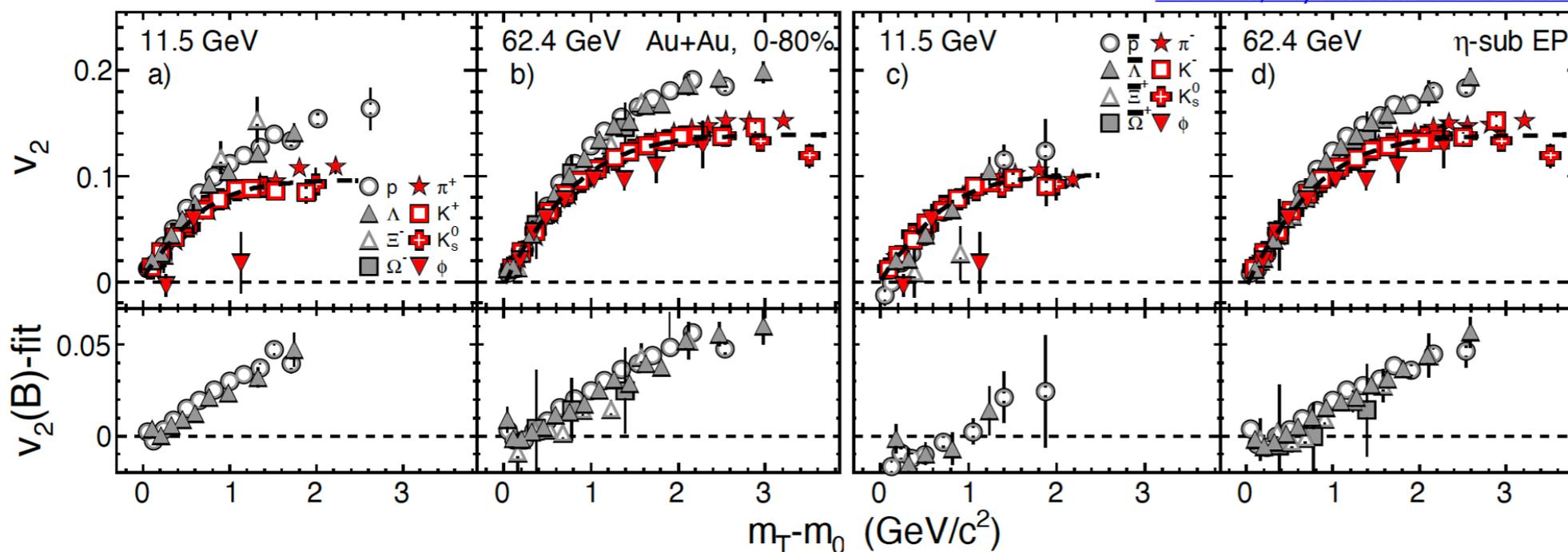
Elliptic flow - Data



- v_2 increases with p_T at low transverse momenta: direct correlation due to boosting
- At large p_T other effects like jets are getting important
- Difference at low p_T : mass dependence
- Difference at larger p_T : quark-number dependence
- Other way of plotting: using m_T instead of p_T , almost cancelation of radial flow effects

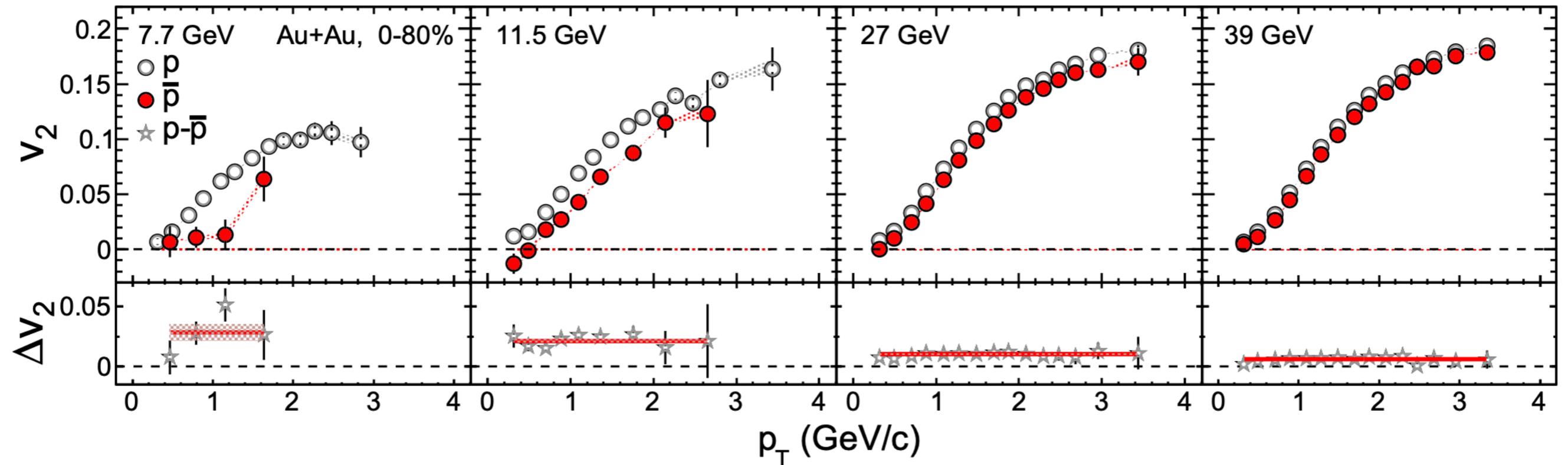
ALI-PUB-109472

[10.1103/PhysRevLett.110.142301](https://arxiv.org/abs/10.1103/PhysRevLett.110.142301)



Difference in elliptic flow between particles and anti-particles

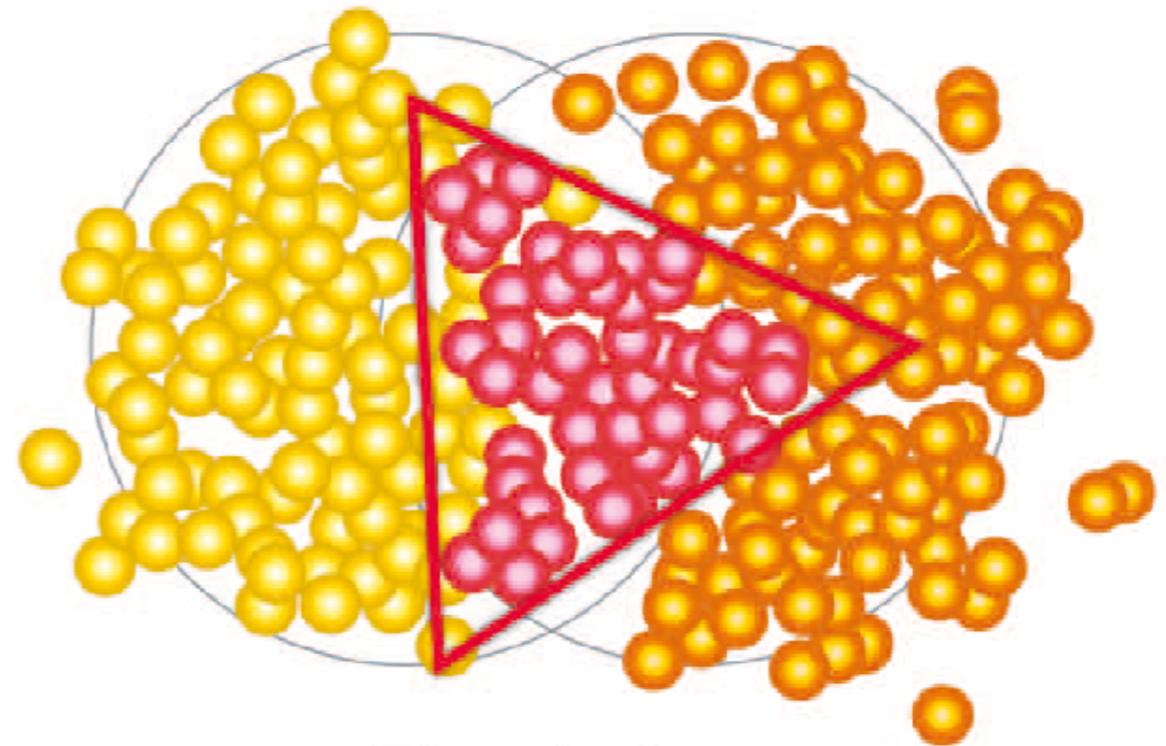
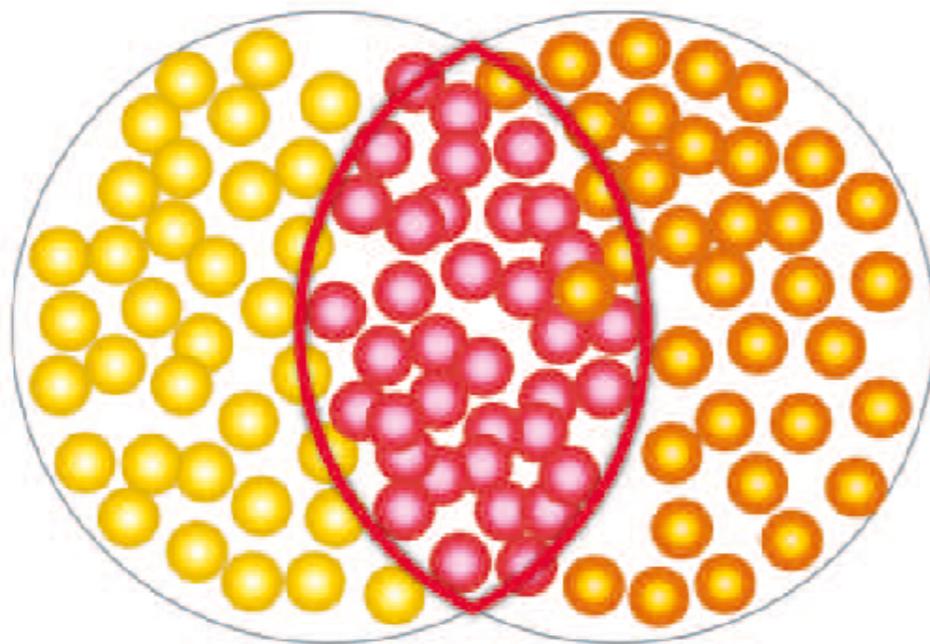
[10.1103/PhysRevLett.110.142301](https://arxiv.org/abs/10.1103/PhysRevLett.110.142301)



- Observation at lower collision energies: particle and anti-particle elliptic flow isn't the same anymore!
- Difference is most likely a result of baryon stopping. Global thermalization doesn't hold anymore on some scale.

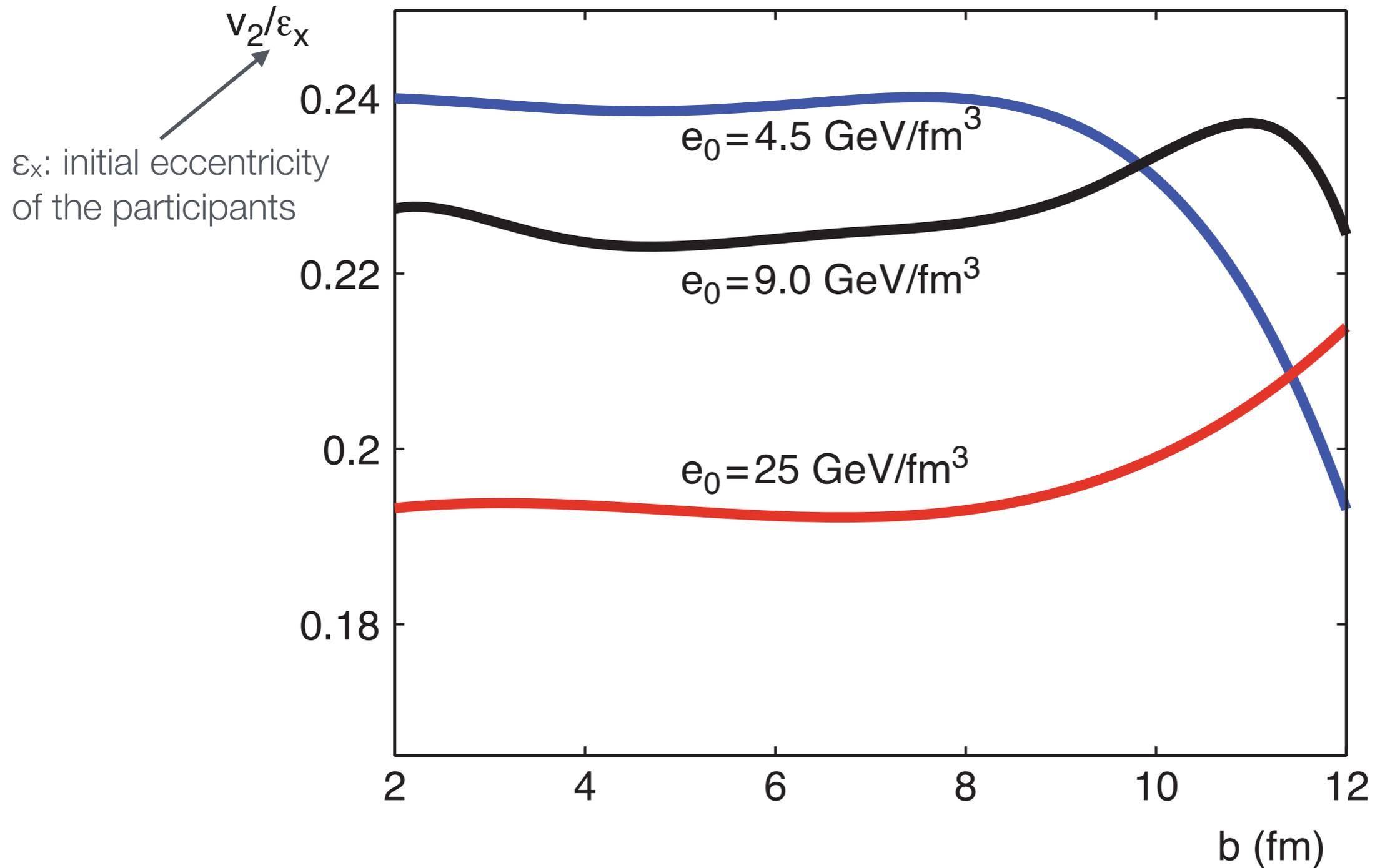
Origin of odd flow components (v_3, v_5, \dots)

- v_2 is related to the geometry of the overlap zone
- Higher moments result from fluctuations of the initial energy distribution



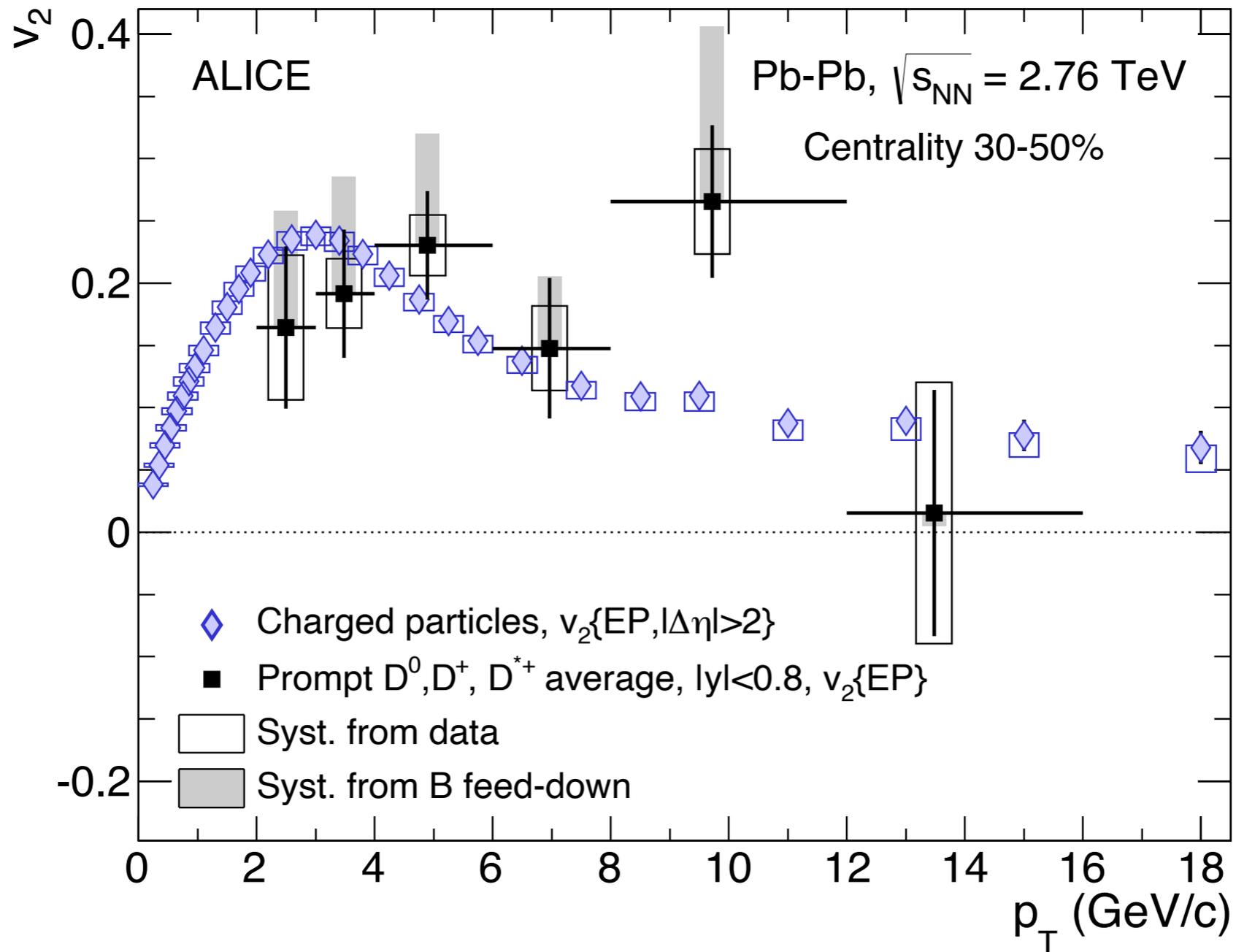
Müller, Jacak, <http://dx.doi.org/10.1126/science.1215901>

Hydrodynamic models: v_2/ε approx. constant



Ideal hydrodynamics gives $v_2 \approx 0.2 - 0.25 \varepsilon$

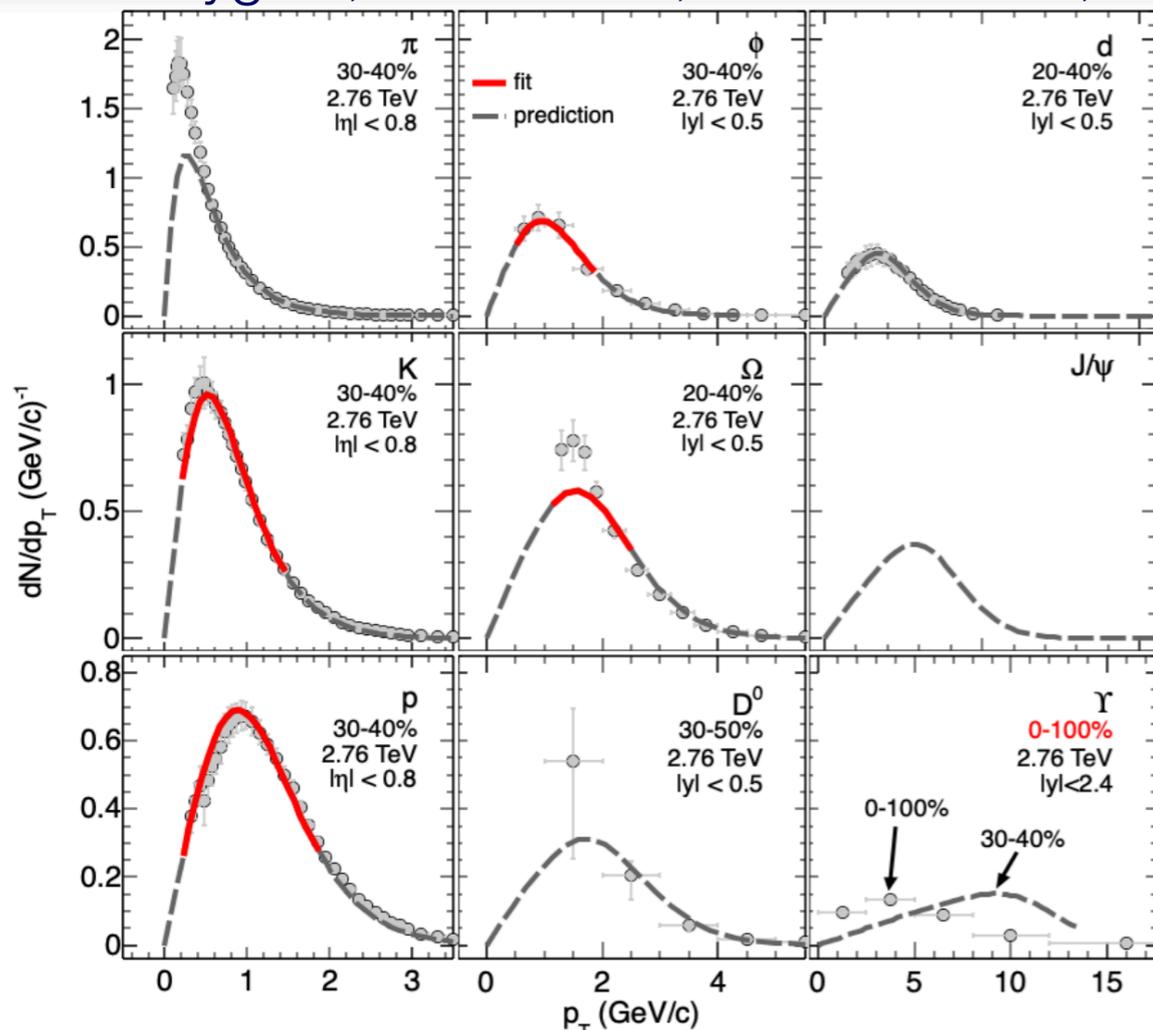
D meson v_2 in Pb-Pb: Heavy quarks seem to flow, too!



Given their large mass, it is not obvious that charm quarks take part in the collective expansion of the medium

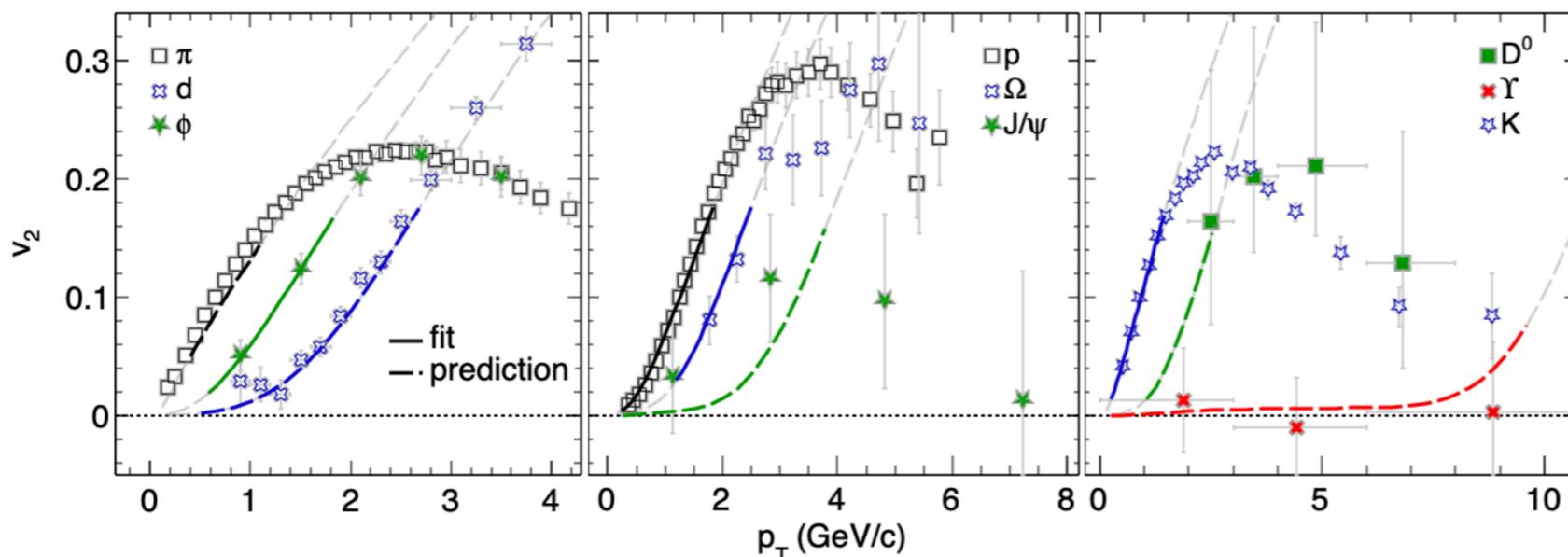
Flow of heavy quarks

K. Reygers, A. Schmah, A. Berdnikova, X. Sun

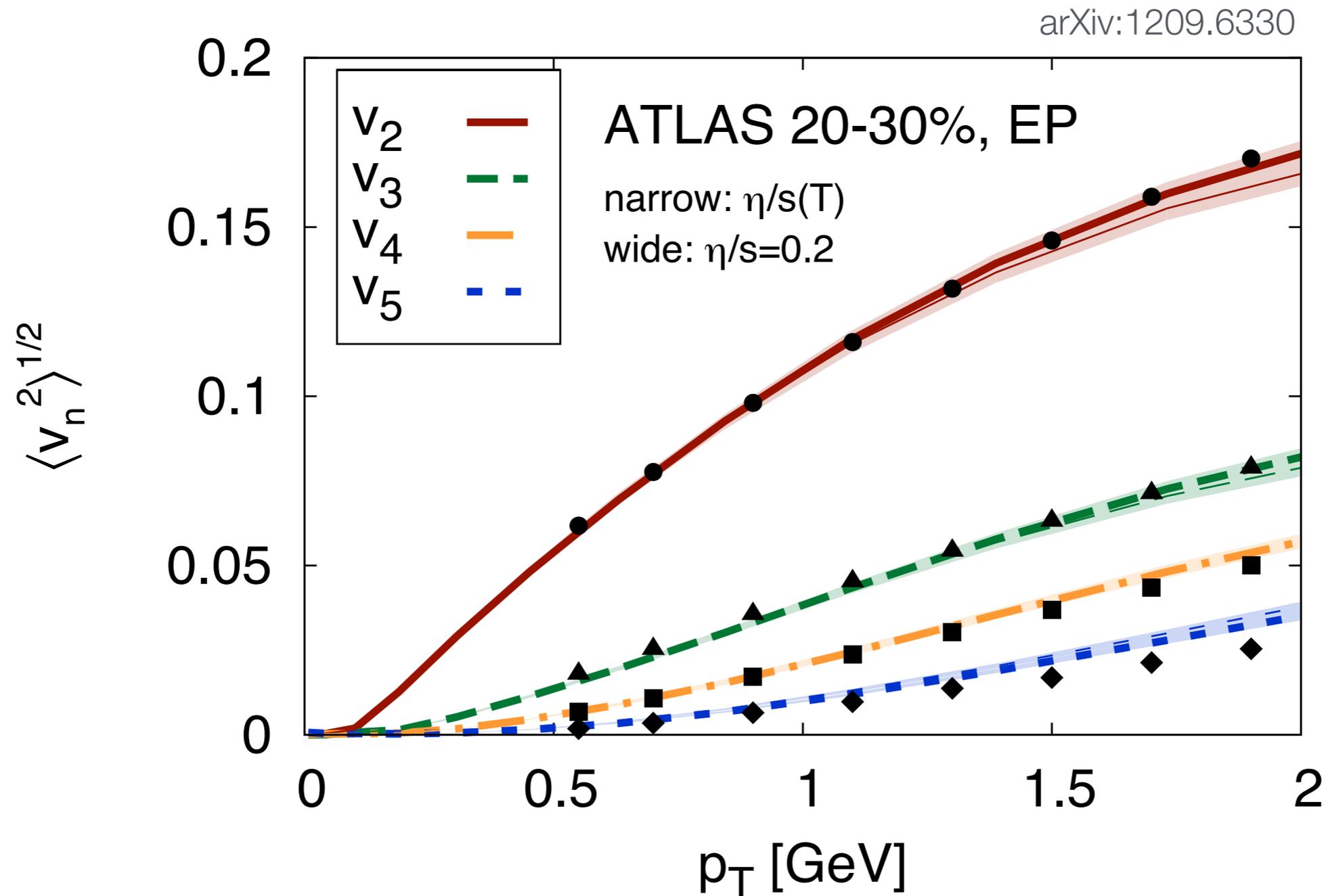


- State of the art blast-wave
- Simultaneous fit to spectra and elliptic flow data
- Good description of $Upsilon$

[10.1103/PhysRevC.101.064905](https://arxiv.org/abs/10.1103/PhysRevC.101.064905)



η/s from comparison to data



Current status (Pb-Pb at $\sqrt{s_{NN}} = 2.76$ TeV):

arXiv:1301.2826

$$(\eta/s)_{\text{QGP}} \approx 0.2 = 2.5 \times \frac{1}{4\pi} \quad (20\% \text{ stat. err.}, 50\% \text{ syst. err.})$$