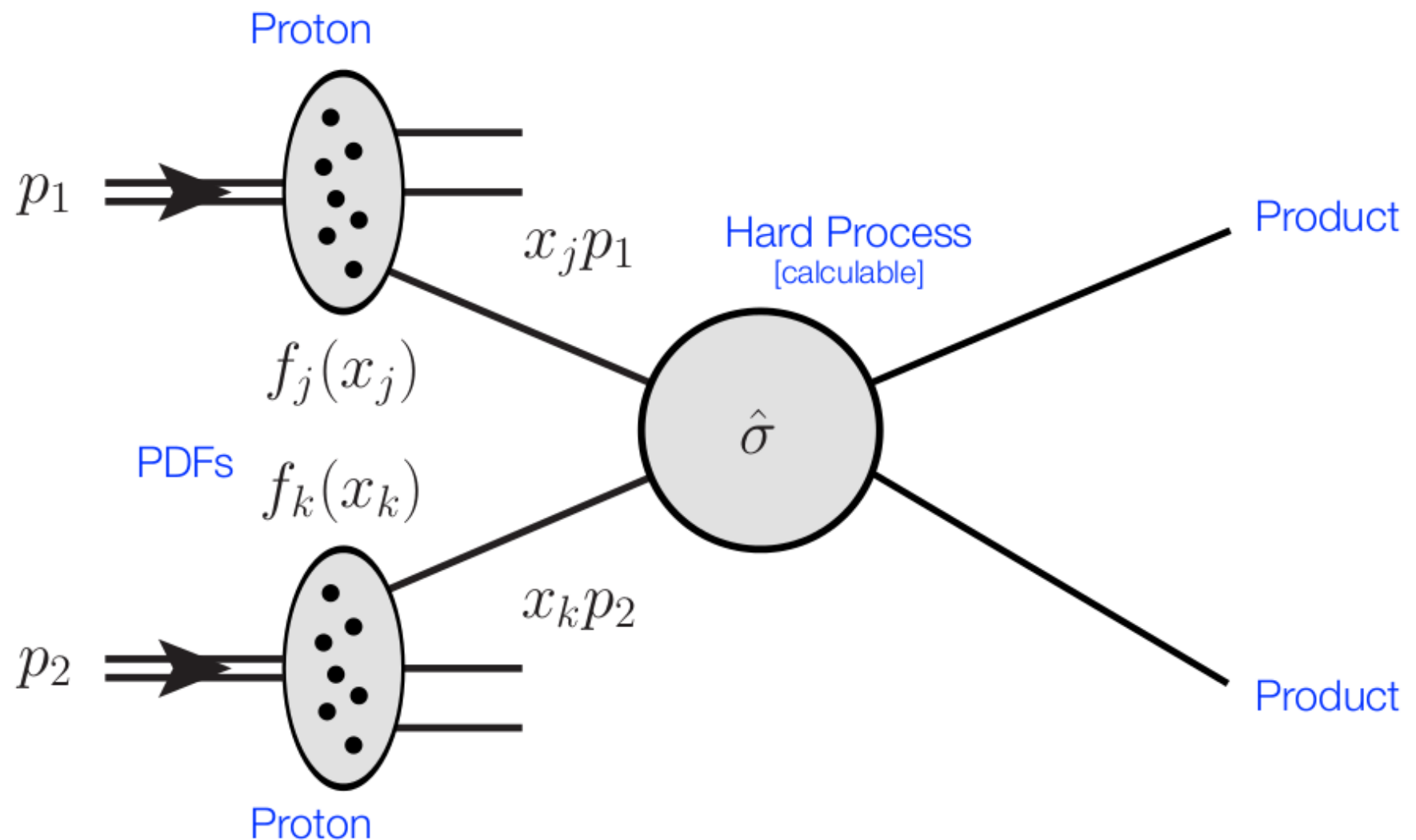
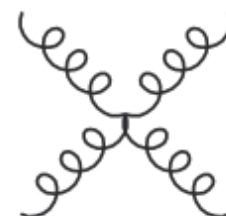
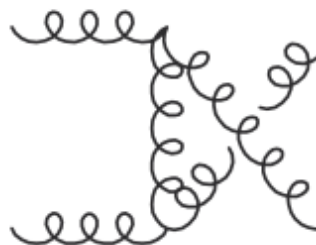
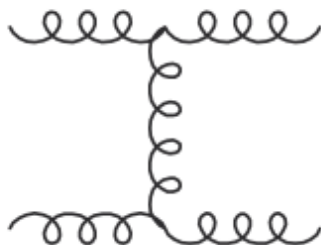
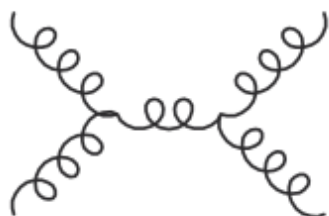
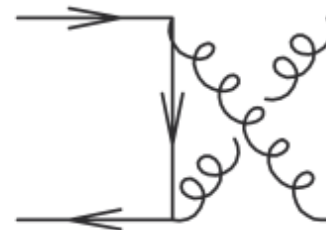
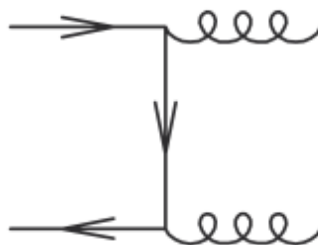
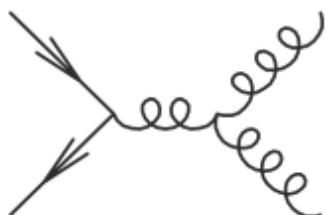
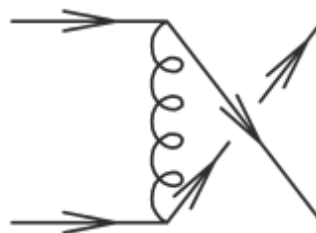
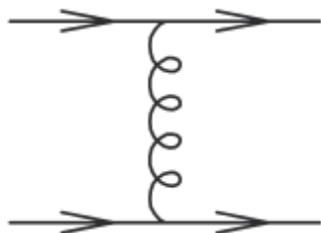


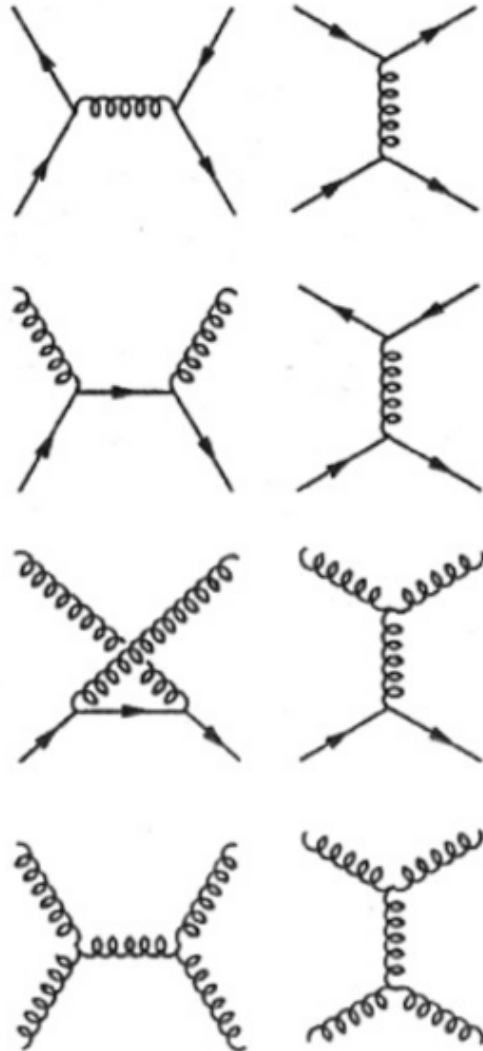
Proton-Proton Scattering @ LHC



Some Hard Processes ...

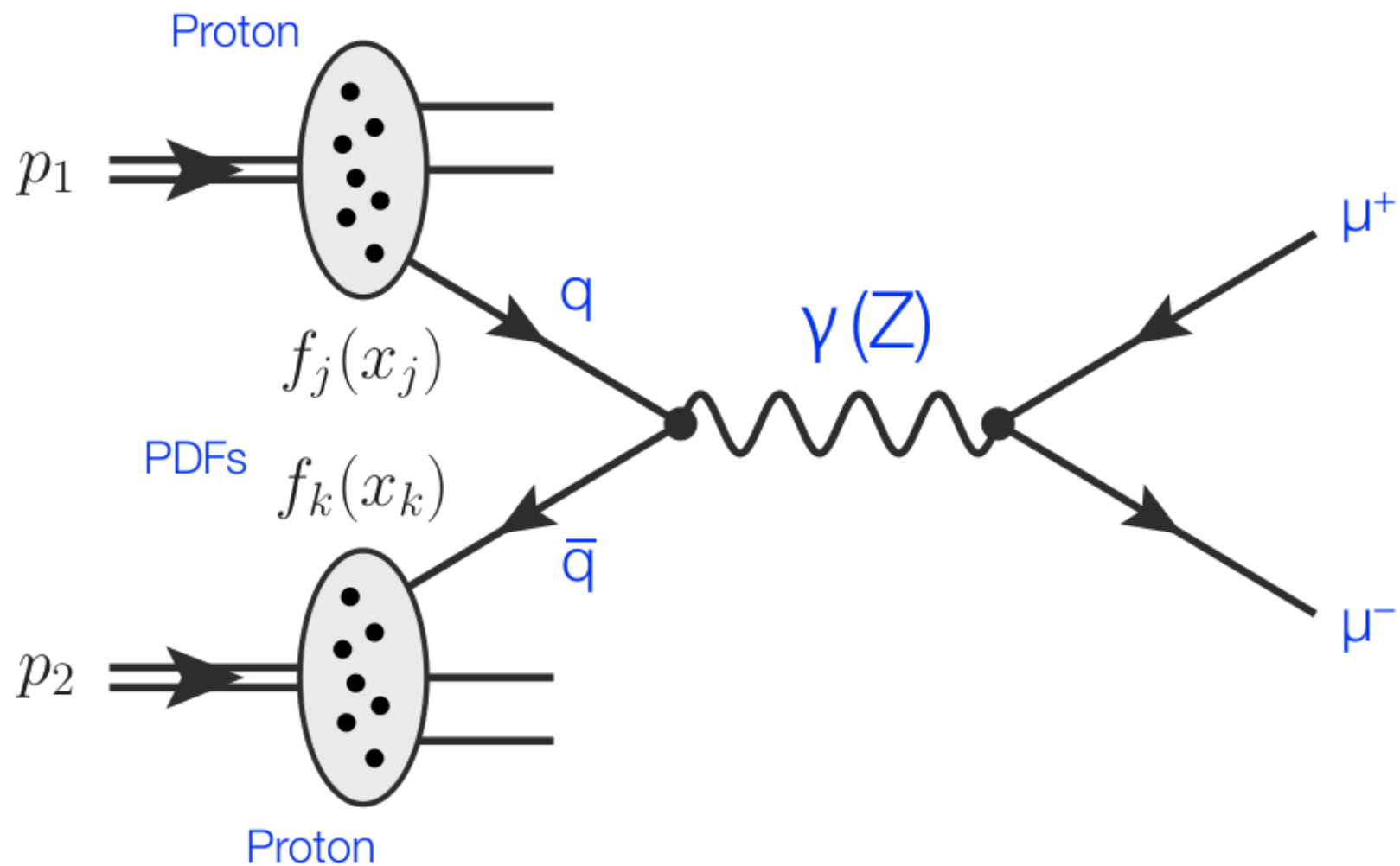


QCD Matrix Elements

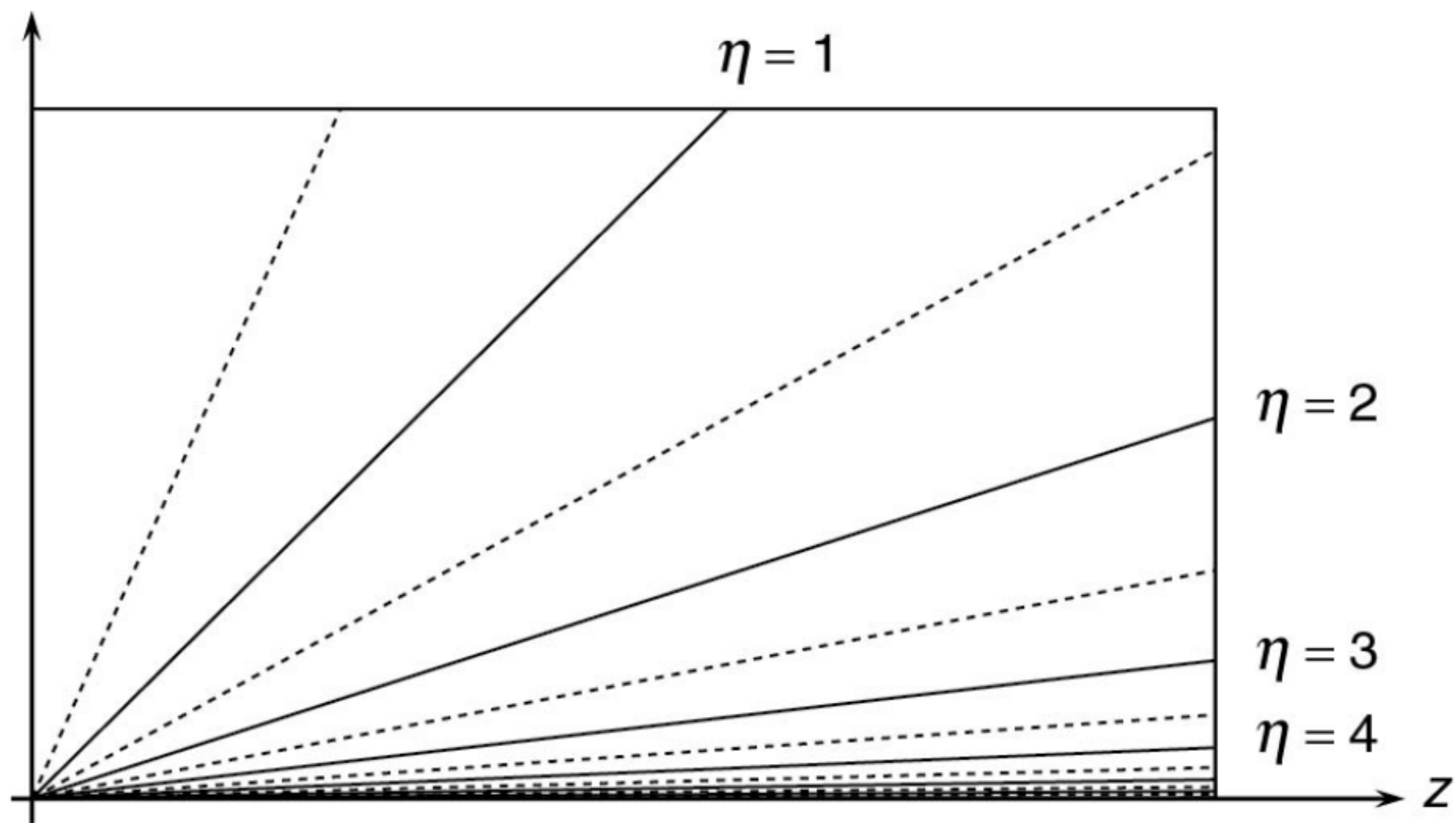


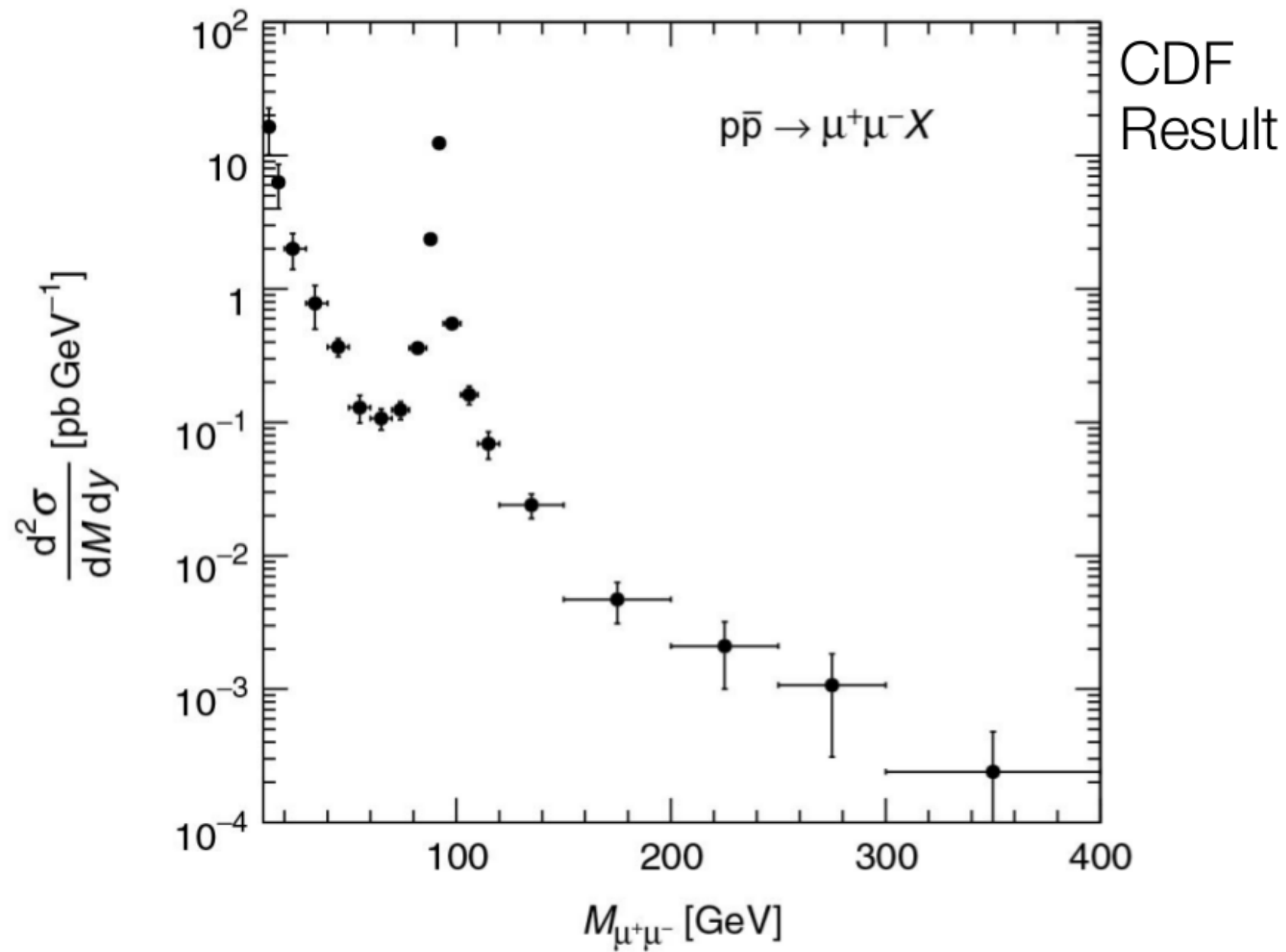
Subprocess	$ \mathcal{M} ^2/g_s^4$	$ \mathcal{M}(90^\circ) ^2/g_s^4$
$qq' \rightarrow qq'$ $q\bar{q}' \rightarrow q\bar{q}'$	$\frac{4}{9} \frac{\hat{s}^2 + \hat{u}^2}{\hat{t}^2}$	2.2
$qq \rightarrow qq$	$\frac{4}{9} \left(\frac{\hat{s}^2 + \hat{u}^2}{\hat{t}^2} + \frac{\hat{s}^2 + \hat{t}^2}{\hat{u}^2} \right) - \frac{8}{27} \frac{\hat{s}^2}{\hat{u}\hat{t}}$	3.3
$q\bar{q} \rightarrow q'\bar{q}'$	$\frac{4}{9} \frac{\hat{t}^2 + \hat{u}^2}{\hat{s}^2}$	0.2
$q\bar{q} \rightarrow q\bar{q}$	$\frac{4}{9} \left(\frac{\hat{s}^2 + \hat{u}^2}{\hat{t}^2} + \frac{\hat{t}^2 + \hat{u}^2}{\hat{s}^2} \right) - \frac{8}{27} \frac{\hat{u}^2}{\hat{s}\hat{t}}$	2.6
$q\bar{q} \rightarrow gg$	$\frac{32}{27} \frac{\hat{u}^2 + \hat{t}^2}{\hat{u}\hat{t}} - \frac{8}{3} \frac{\hat{u}^2 + \hat{t}^2}{\hat{s}^2}$	1.0
$gg \rightarrow q\bar{q}$	$\frac{1}{6} \frac{\hat{u}^2 + \hat{t}^2}{\hat{u}\hat{t}} - \frac{3}{8} \frac{\hat{u}^2 + \hat{t}^2}{\hat{s}^2}$	0.1
$qg \rightarrow qg$	$\frac{\hat{s}^2 + \hat{u}^2}{\hat{t}^2} - \frac{4}{9} \frac{\hat{s}^2 + \hat{u}^2}{\hat{u}\hat{s}}$	6.1
$gg \rightarrow gg$	$\frac{9}{4} \left(\frac{\hat{s}^2 + \hat{u}^2}{\hat{t}^2} + \frac{\hat{s}^2 + \hat{t}^2}{\hat{u}^2} + \frac{\hat{u}^2 + \hat{t}^2}{\hat{s}^2} + 3 \right)$	30.4

Example: Drell-Yan Process

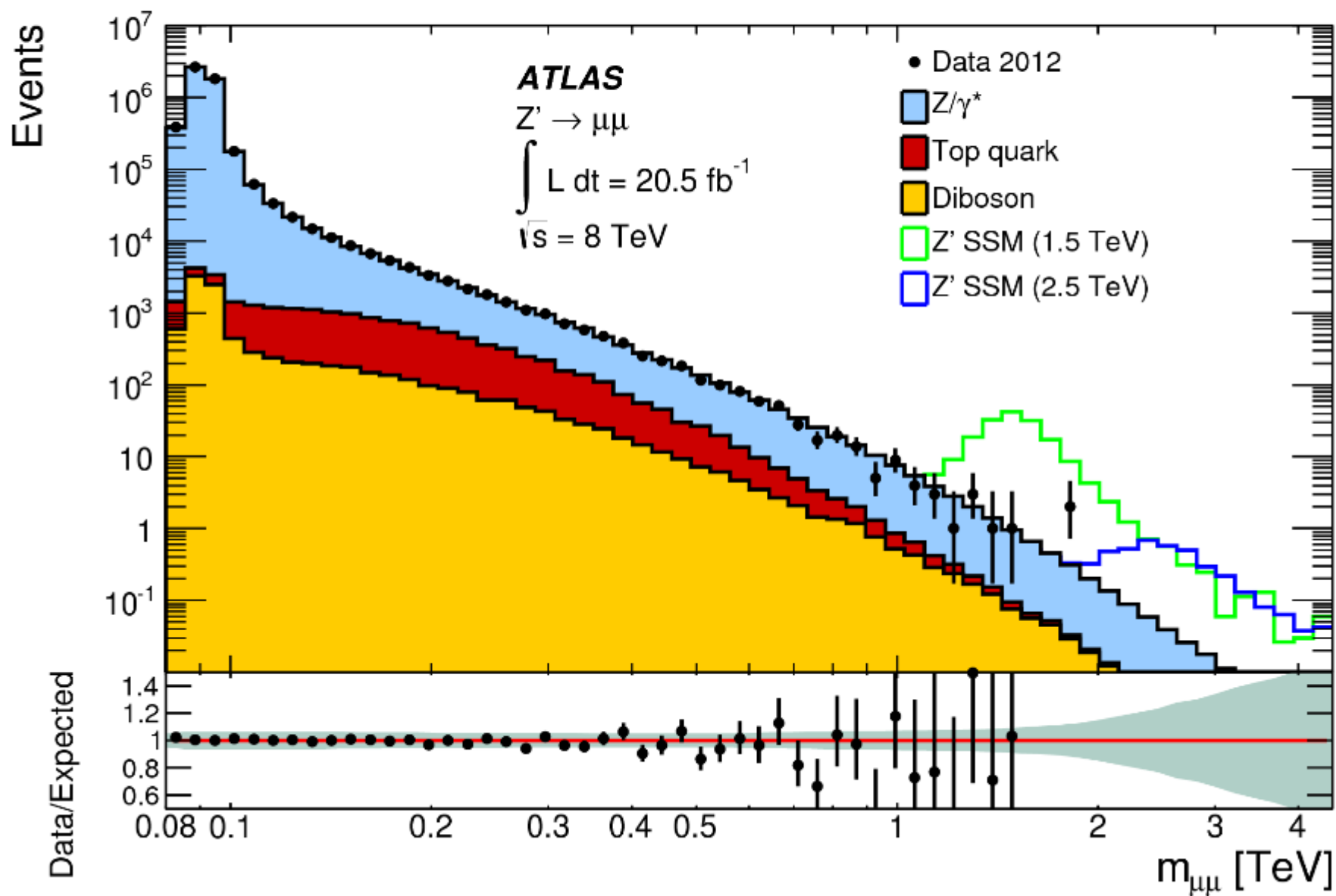


Pseudorapidity Regions

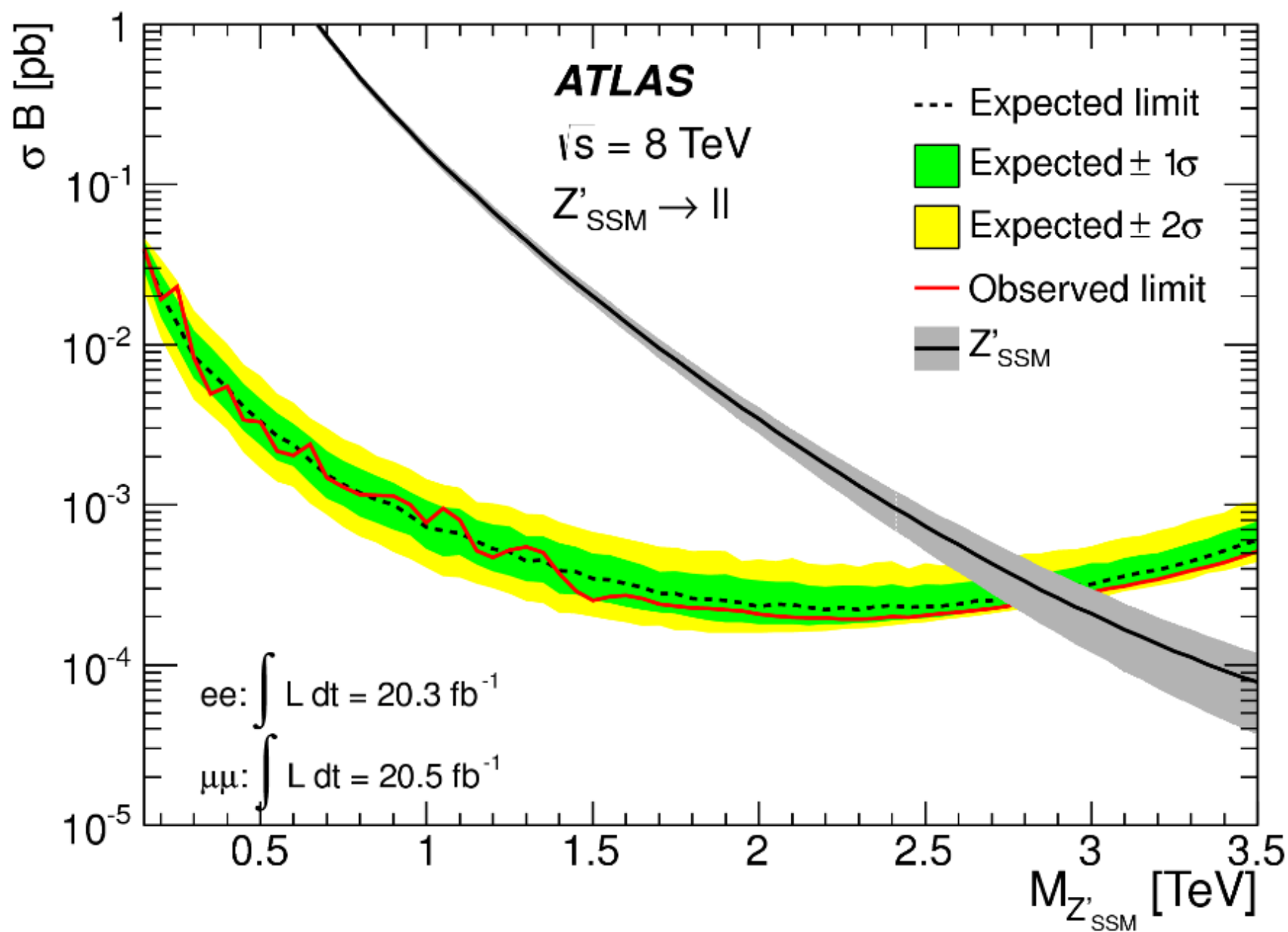




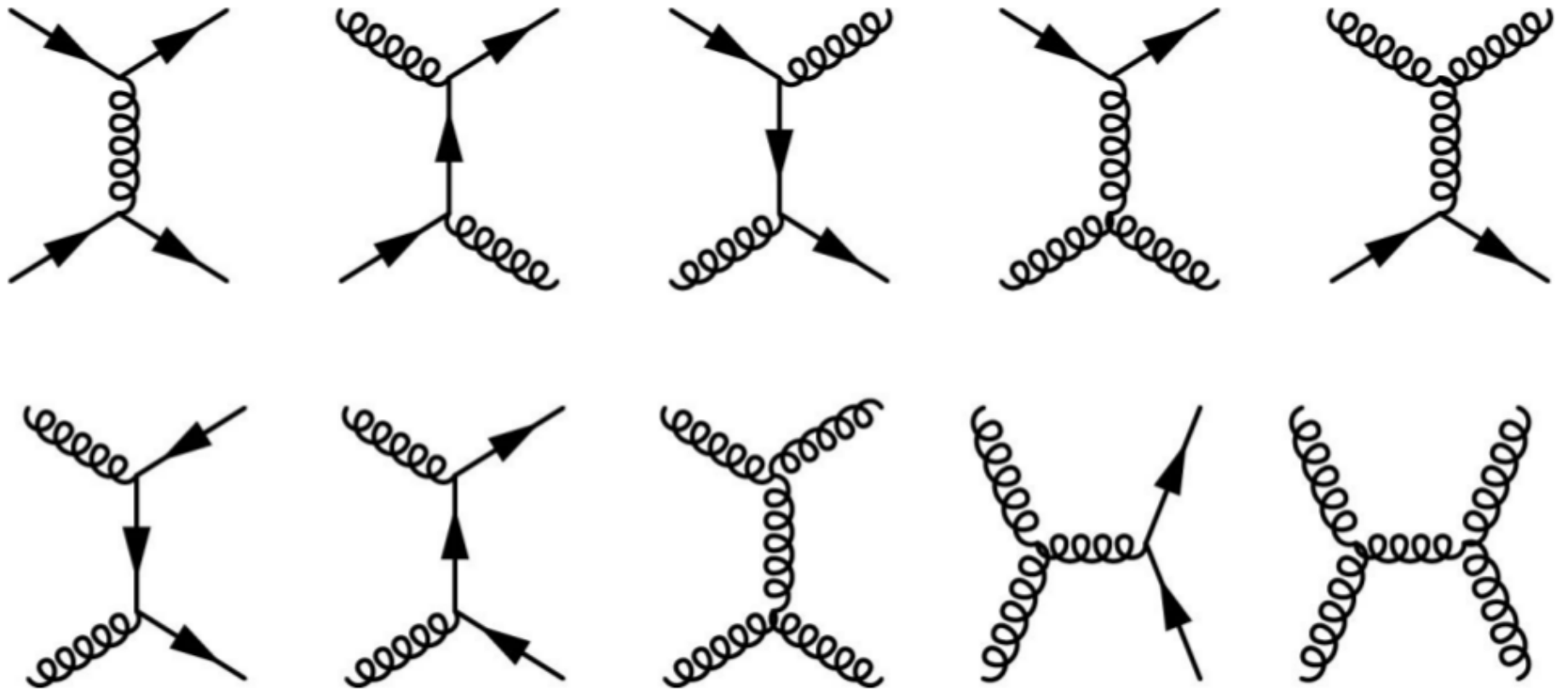
ATLAS Z' Search Result



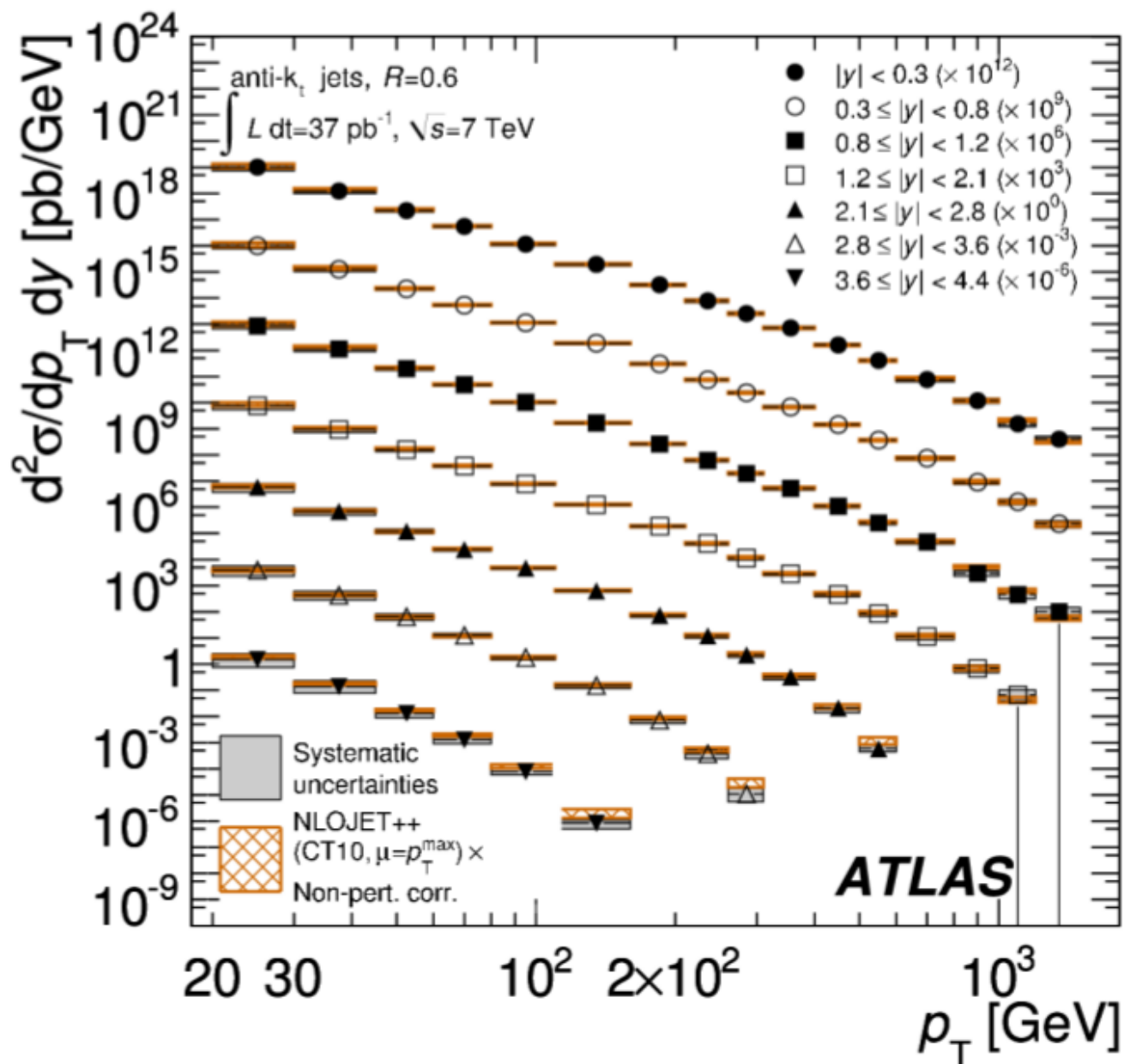
ATLAS Z' Search Result



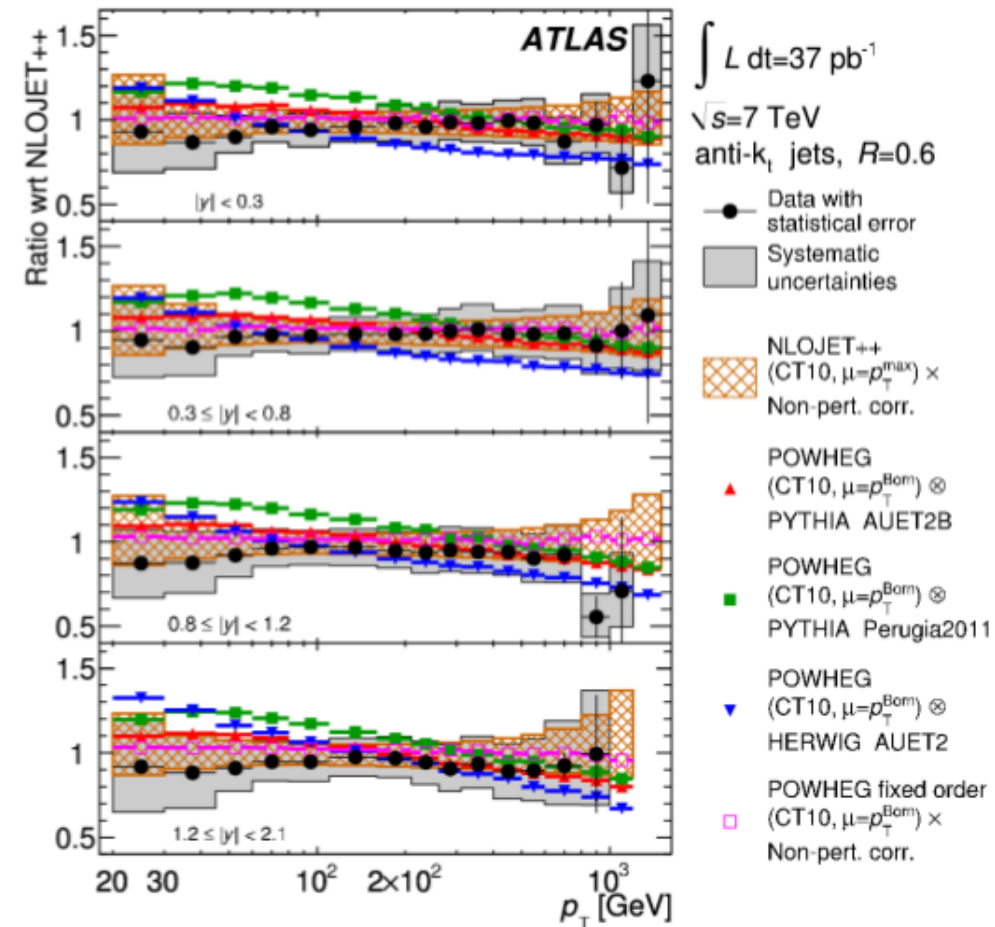
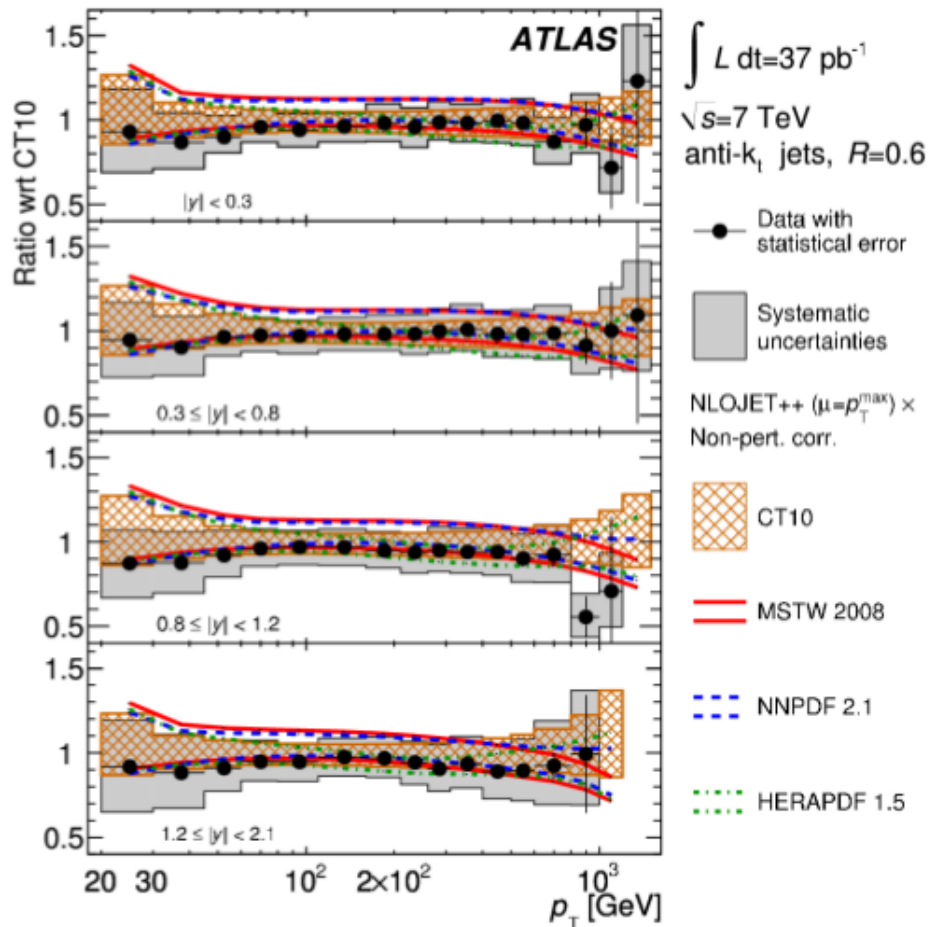
QCD Dijet-Production



Inclusive Jet Cross Section



Inclusive Jet Cross Section



Reminder of Parity

particle solution $\Psi_i = u_i(E, \vec{p}) e^{i(\vec{p}\vec{x} - Et)}$

$$u_1 = \sqrt{E + m} \begin{pmatrix} 1 \\ 0 \\ \frac{p_z}{E+m} \\ \frac{p_x + ip_y}{E+m} \end{pmatrix} \quad u_2 = \sqrt{E + m} \begin{pmatrix} 0 \\ 1 \\ \frac{p_x - ip_y}{E+m} \\ \frac{-p_z}{E+m} \end{pmatrix}$$

antiparticle solution $\Psi_i = v_i(E, \vec{p}) e^{-i(\vec{p}\vec{x} - Et)}$

$$v_1 = \sqrt{E + m} \begin{pmatrix} \frac{p_x - ip_y}{E+m} \\ \frac{-p_z}{E+m} \\ 1 \\ 0 \end{pmatrix} \quad v_2 = \sqrt{E + m} \begin{pmatrix} \frac{p_z}{E+m} \\ \frac{p_x + ip_y}{E+m} \\ 0 \\ 1 \end{pmatrix}$$

These solutions have positive energies.

Parity operator P:

$$\gamma^0 = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix}$$

Spin ½ particles AT REST have intrinsic parity P = +1

Spin ½ particles AT REST have intrinsic parity P = -1

Historical θ/τ

In 1956, parity conservation as well as T and C symmetry was a “dogma”

→ very little experimental tests done

θ/τ puzzle:

$$\theta \rightarrow \pi^+ \pi^0; \quad P(\pi^+ \pi^0) = +1$$

$$\tau \rightarrow \pi^+ \pi^+ \pi^-; \quad P(\pi^+ \pi^+ \pi^-) = -1$$

$$P(q) = 1; P(\bar{q}) = -1;$$

$$P(\text{meson}) = P_q P_{\bar{q}} (-1)^L;$$

$$\text{lowest energy, } S = 0$$

$$P = -1$$

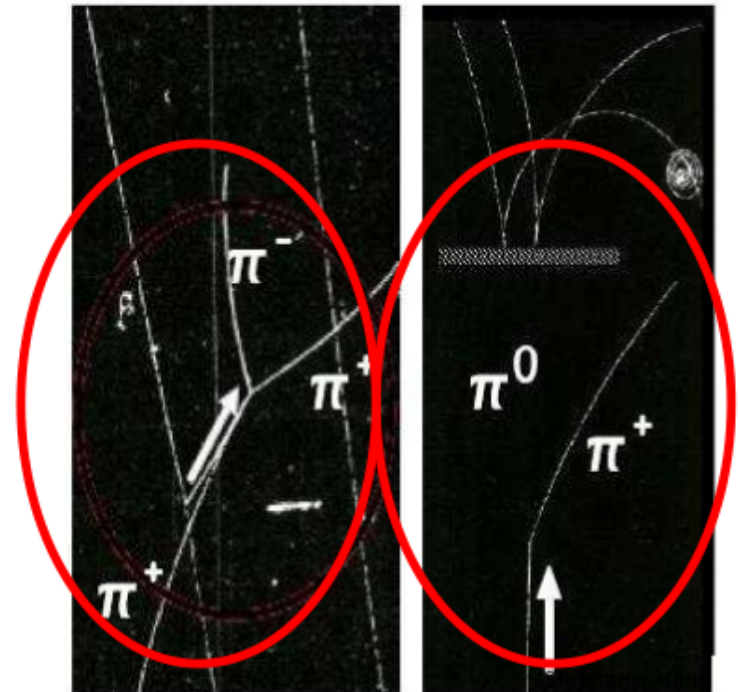
θ, τ have same mass, same lifetime, however different parity ...

Yang, Lee:

$$\rightarrow \theta = \tau = K^+$$

weak interaction violates parity

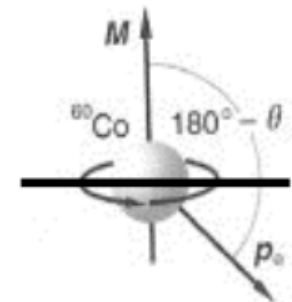
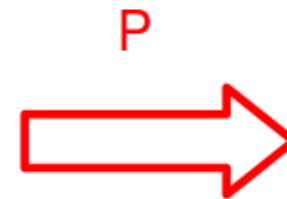
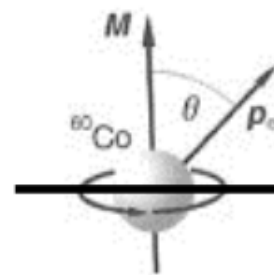
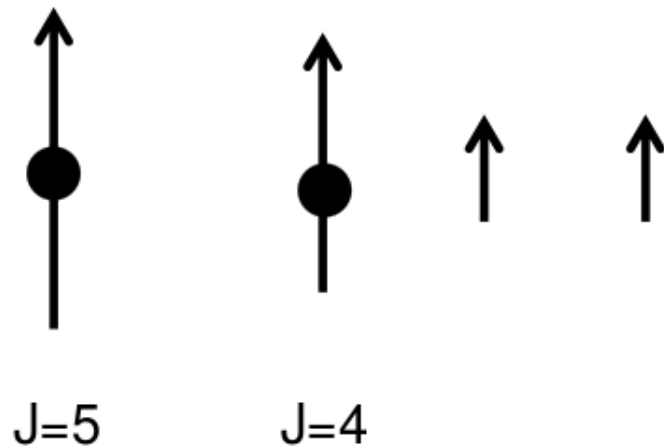
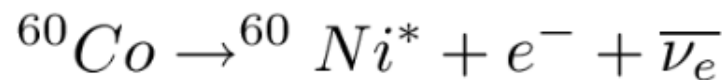
proposed a set of measurements which
test parity



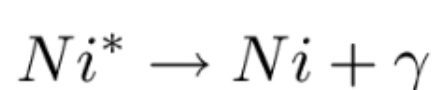
Wu-Experiment

Partly conservation: physics stays invariant under parity conservation

Idea: Check that number of electrons emitted in direction of spin (\vec{J}) of ^{60}Co and in opposite direction ($-\vec{J}$) are the same.



Experiment: Invert polarization of ^{60}Co and compare electron rate in same angle θ



photons are preferentially emitted in direction of spin. Use photon distribution to test polarization of ^{60}Co . (elm IA conserves parity)

MAIN CHALLENGE: Polarization of ^{60}Co

Spin of ^{60}Co : $J=5 \rightarrow M = -5, -4, \dots, 4, 5$

$$M=5 \begin{array}{|c|} \hline \hline \hline \hline \hline \hline \end{array} \} \Delta E = g \mu_K B$$

$$\mu_K \sim 5.05 \times 10^{-27} \text{ J/T}$$

Population of energy levels follows Boltzmann distribution:

$$e^{-\frac{E}{k_B T}}$$

$$M=-5 \begin{array}{|c|} \hline \hline \hline \hline \hline \hline \end{array}$$

for $\Delta E \gg k_B T$ only lowest energy level is populated, however for given B field in experiment (2.3 T) **very low temperatures** needed

g factor depends on gitter structure

Example: $g = 7.5$ (^{60}Co), $B = 2.3 \text{ T}$, $T = 0.003 \text{ K}$

$$\frac{P(m=-4)}{P(m=-5)} = e^{-\frac{\Delta E}{k_B T}} = 0.074 \rightarrow 92\% \text{ polarized } ^{60}\text{Co}$$

Solution Part-I: embedding ^{60}Co in a paramagnetic material ($B \sim \mu_r$; $\mu_r \sim 3-4$)
still temperatures of $T=0.01\text{K}$ needed

Adiabatic Colling

1926 von Debye proposed method to create low temperature

Fundamental relation of thermodynamics:

$$dU = T dS - p dV$$

1. Step: **isotherm** magnetisation

- paramagnetic material **in helium gas** is put into magnetic field
- energy levels are split up, only lower ones are populated
- entropy gets smaller: $dS < 0 \rightarrow dU < 0$, **helium gas absorbs heat**

2. Step: helium gas removed \rightarrow **thermal isolation** of nitrit

3. Step: adiabatic cooling

- **magnetic field is slowly switched off**
- split off of energy levels get smaller
- system likes to polpulate higher states, however $dU = \text{const}$ due to **isolation**
- dS gets larger thus **T gets smaller**



Caveat: need magnetic field to get polarized ^{60}Co

How to combine Cooling and Polarization?

Two competing effects needed in the nitrit-crystal to get high degree of polarization

- 1) Need high B field and low temperature to get polarization
- 2) Switch off B field to lower temperature via adiabatical cooling
B field on \rightarrow warm up, B field off \rightarrow cool down

How does this work?

Solution: Some paramagnetic material have large anisotropic distribution of g-factors
(artefact of crystal structure, different binding mechanisms)

B field for adiabatic cooling in direction with high g-factor

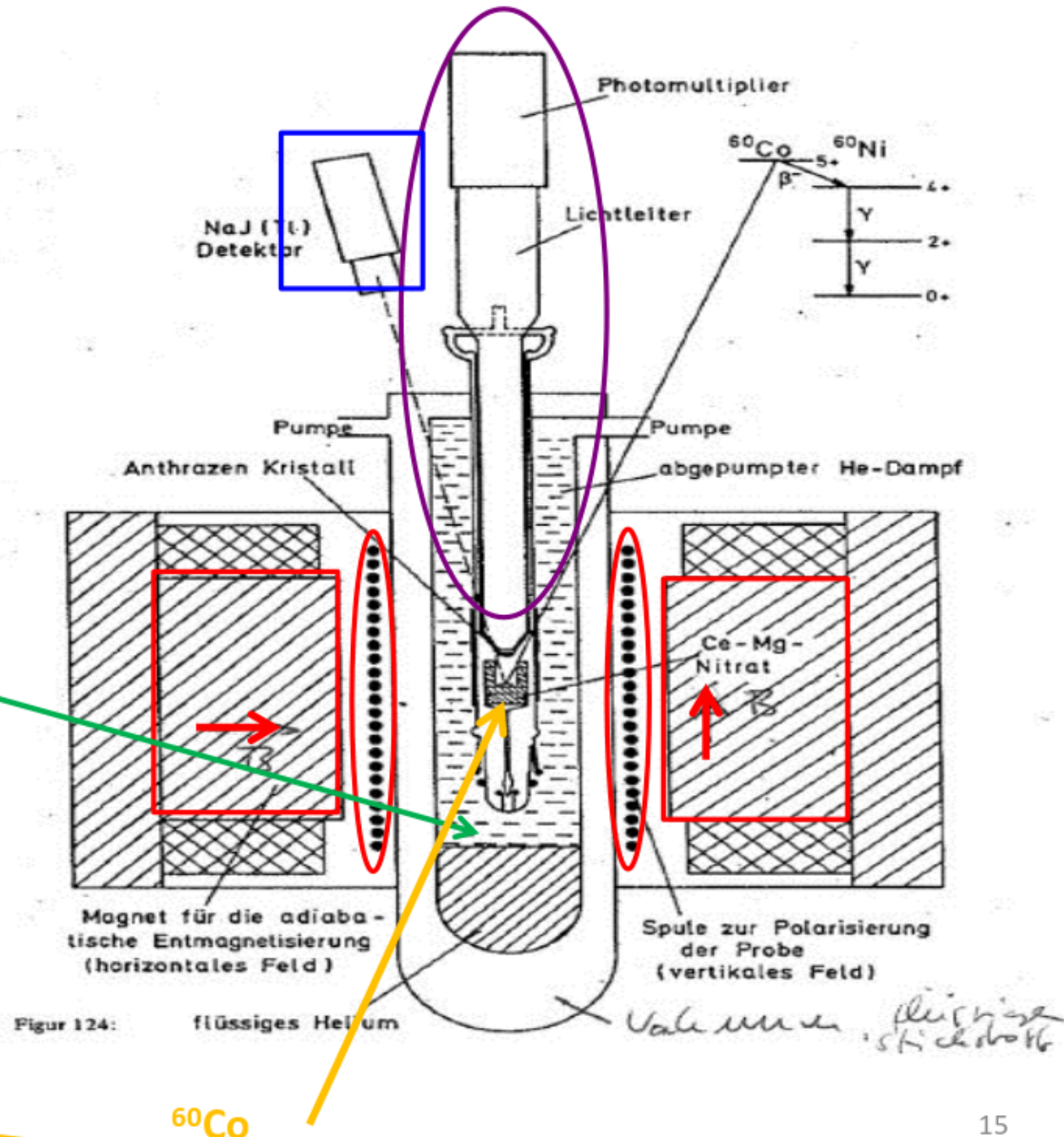
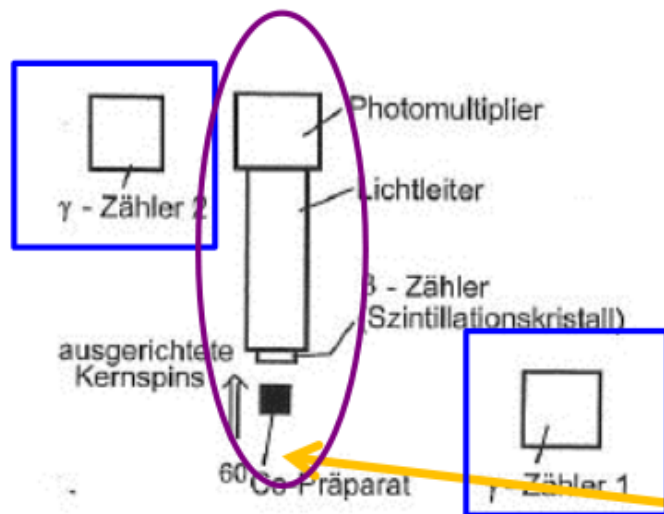
Thus large split up of energy levels, thus large cooling effect

B field for polarization in direction of low g-factor, thus only little warm up

Wu-Experiment

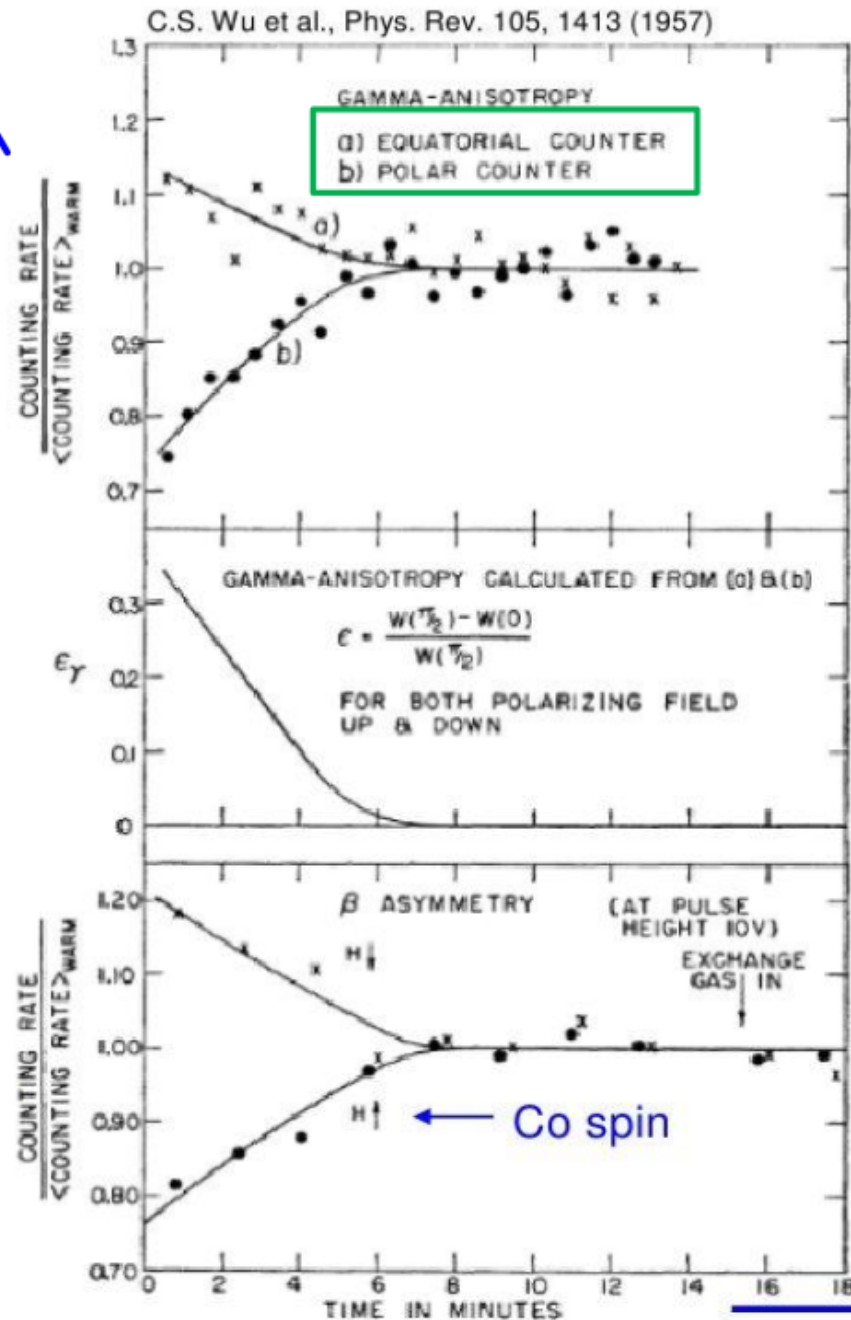
Requirements:

- 2 B fields in orthogonal directions
- detection of emitted electron
(cover a small opening angle Θ)
- detection of emitted gamma
(to test polarization of ^{60}Co)
- crystal needs to be located
in helium bath first than in vacuum



Wu-Experiment: Results

counting rate relative to warm (unpolarized) rate



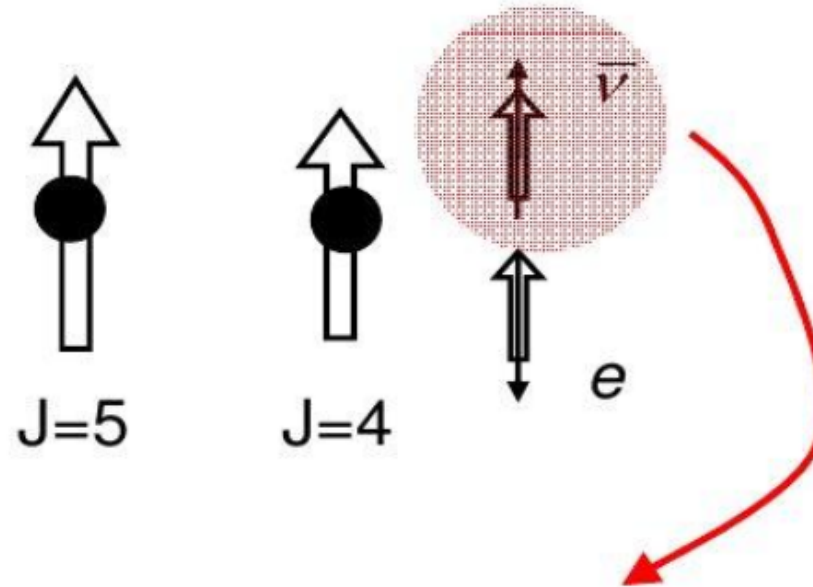
measure photon anisotropy, to determine degree of polarization

electron rates are different depending on the polarization!

warm up with time

→ parity violation

Qualitative Explanation



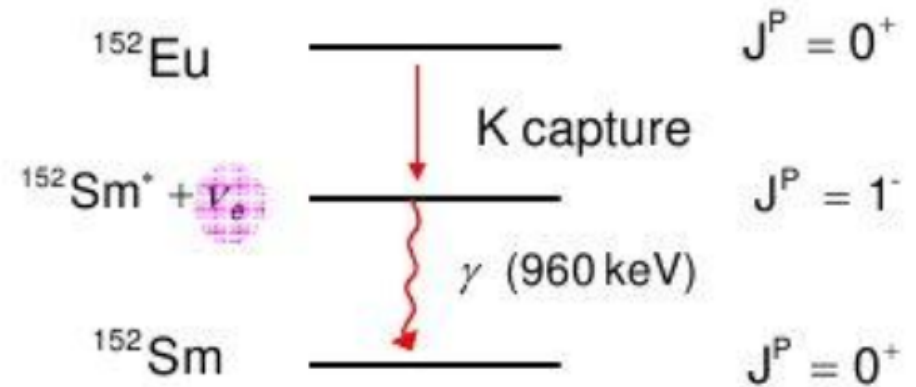
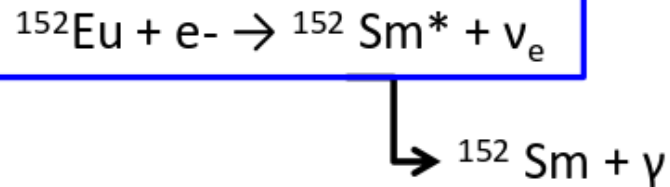
Consequence of existence of
only left-handed (LH) neutrinos
(RH anti-neutrinos)

Wu experiment established P violation!

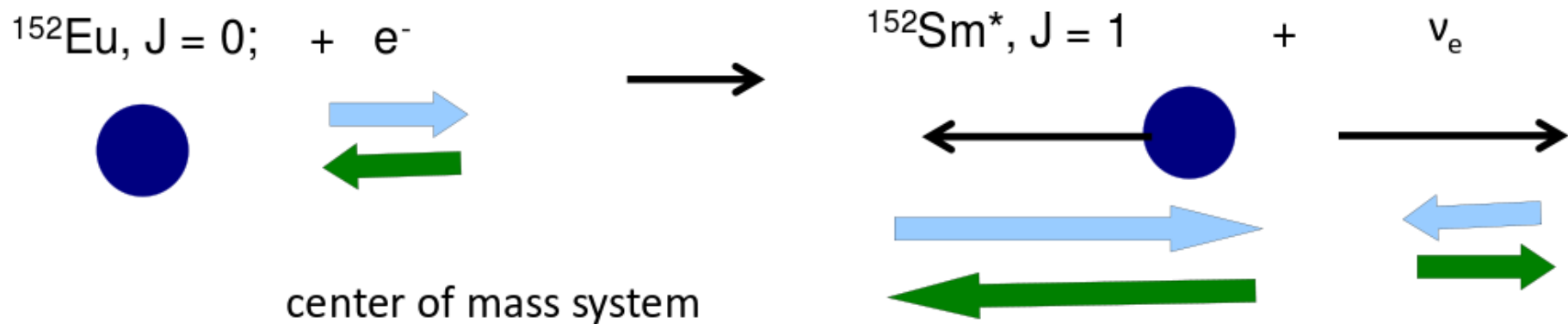
It was however not precise enough to measure helicity of neutrino
 $H \sim 0.7 \pm \text{large uncertainties}$

 Goldhaber experiment

Goldhaber Experiment

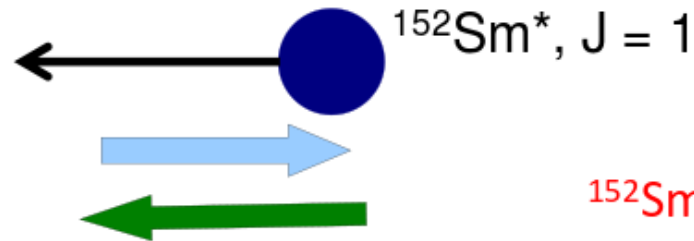
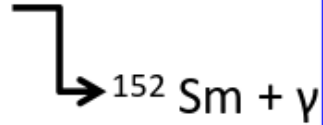
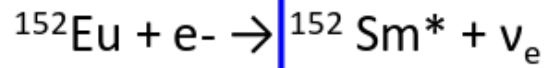


Light blue and green arrows indicate possible spin configurations

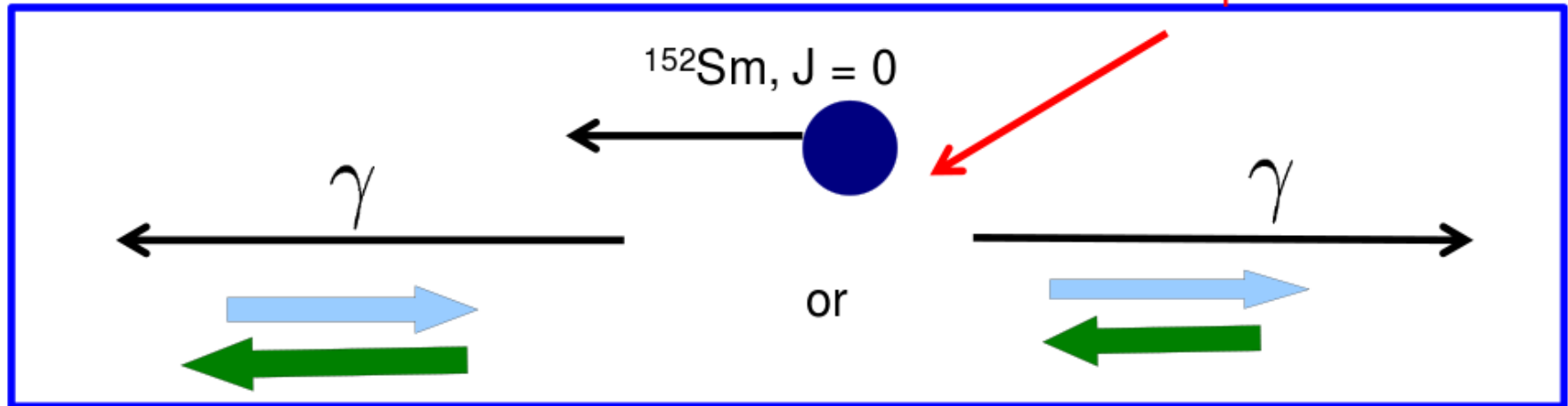


spin of neutrino is in **opposite direction** than the one of $^{152}\text{Sm}^*$,
momentum of is **in opposite direction** than the one of $^{152}\text{Sm}^*$

Goldhaber Experiment



^{152}Sm get's small recoil from photon!



direction of spin of photon is opposite of neutrino

emitted in direction of Sm^*

$$h(\gamma) = h(\nu_e)$$

emitted in opposite direction of Sm^*

$$h(\gamma) = -h(\nu_e)$$

Two open question: 1) What is the direction of emission of the photon?

2) What is the polarization of the photon?

Resonant Scattering

To compensate the nuclear recoil, the photon energy must be slightly larger than 960 keV.

This is the case for photons which have been emitted in the direction of the $\text{Eu} \rightarrow \text{Sm}$ recoil (Doppler-effect).



Resonant scattering only possible for "forward" emitted photons, which carry the polarization of the Sm^* and thus the polarization of the neutrinos.

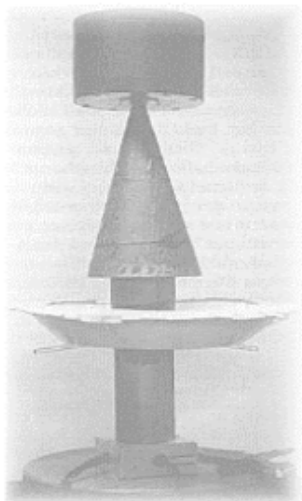


foto of exeperiment

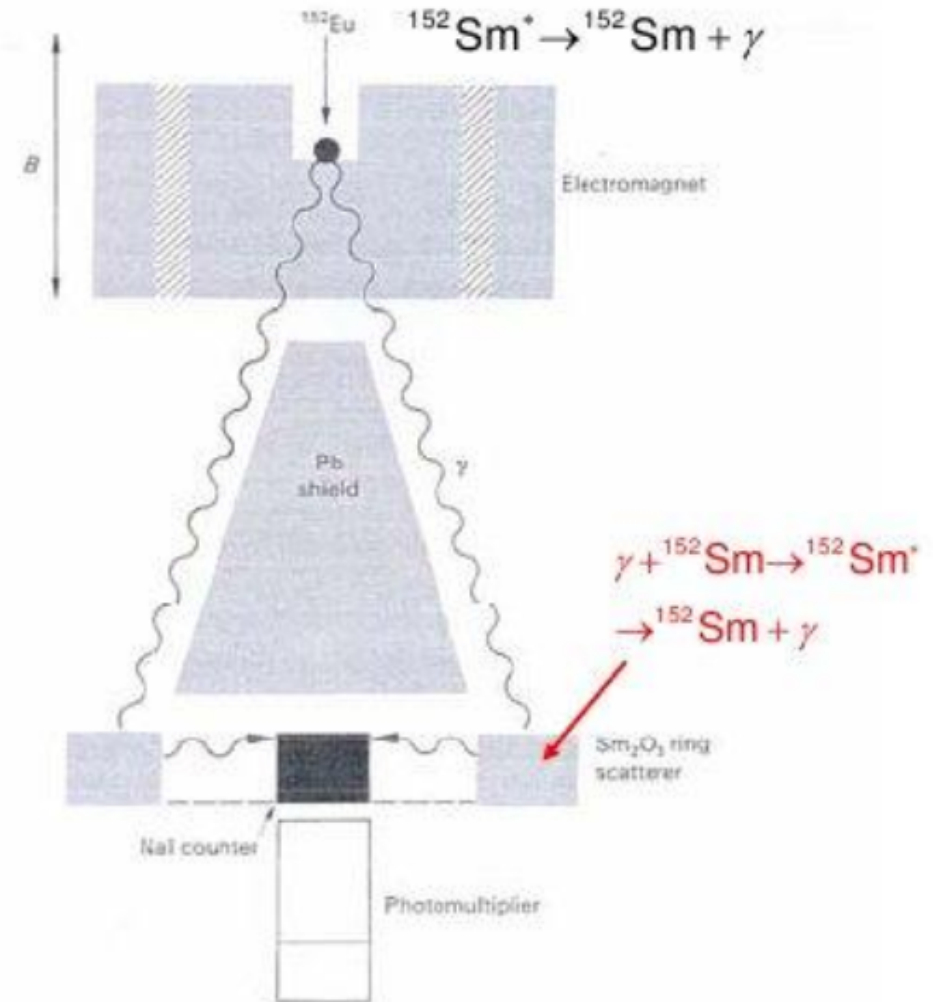
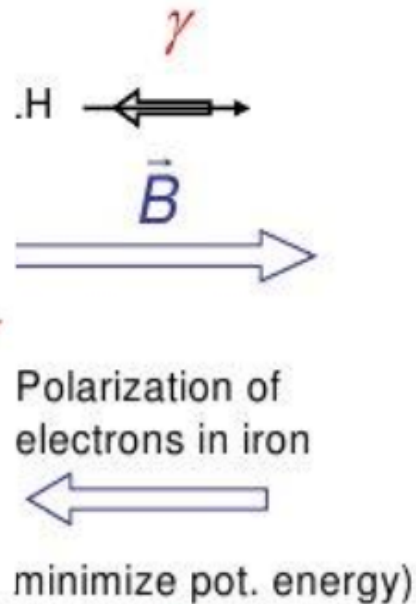
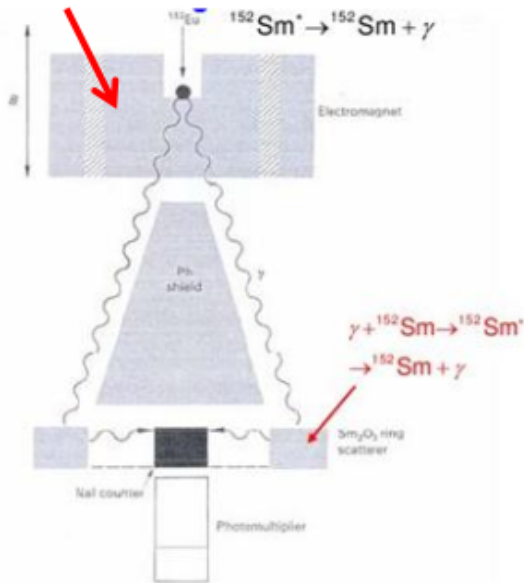


Fig. 7.8. Schematic diagram of the apparatus used by Goldhaber *et al.*, in which γ -rays from the decay of $^{152}\text{Sm}^*$, produced following K-capture in ^{152}Eu , undergo resonance scattering in Sm_2O_3 and are recorded by a sodium iodide scintillator and photomultiplier. The transmission of photons through the iron surrounding the source depends on their helicity and the direction of the magnetic field B .

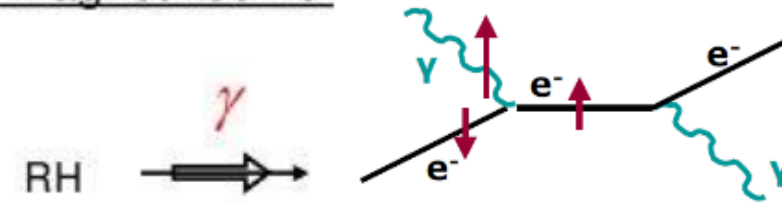
Measurement of Polarization of Photon

Exploit that the transmission index through magnetized iron is polarization dependent: Compton scattering in magnetized iron

iron in B field



LH photons cannot be absorbed:
Good transmission



Absorption leads to spin flip

RH photons undergo Compton scattering: Bad transmission

Photons w/ polarization anti-parallel to magnetization undergo less absorption

Goldhaber Experiment: Result

- Due to geometry of experiment, only resonant scattered photons are detected
Helicity of detected photons identical to helicity of neutrino.
- Detect photons which pass through magnetized iron.

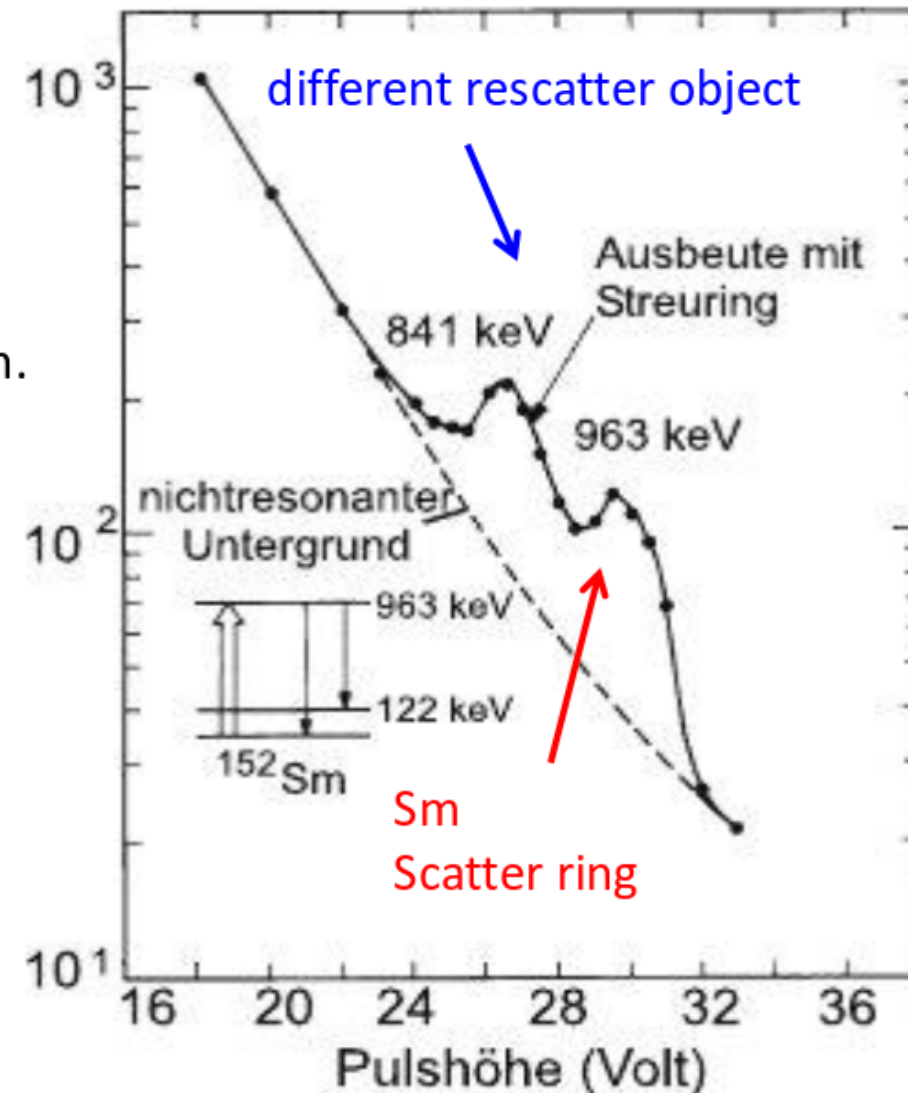
B field points in flight direction of photons
→ measure fraction of (mainly) LH photons

B field points in opposite direction
→ measure fraction of (mainly) RH photons

$$\delta = \frac{N_- - N_+}{0.5(N_- + N_+)} = 0.017 \pm 0.003$$

N_- : counting rate with magnetic field down

N_+ : counting rate with magnetic field up



Goldhaber-Experiment: Result

Result: $\delta = +0.017 \pm 0.003$

Theoretical expectation (for 100% polarized photons)

$$\delta = \pm 0.025 \quad \left\{ \begin{array}{ll} \text{„-“} & \text{for } h > 0 \\ \text{„+“} & \text{for } h < 0 \end{array} \right.$$

Only 5-8% of electrons in iron are polarized, thus asymmetry of scattering for LH and RH photons is very small, thus heavily wash out the asymmetry.

Due to background effects (thermal movements inside the source, polarisation can depend on angle, ...) expect for pure LH neutrinos 75% polarized photons

Measured photo polarisation: $66 \pm 15\%$, consistent with 75%

 **Neutrino are left handed particles**